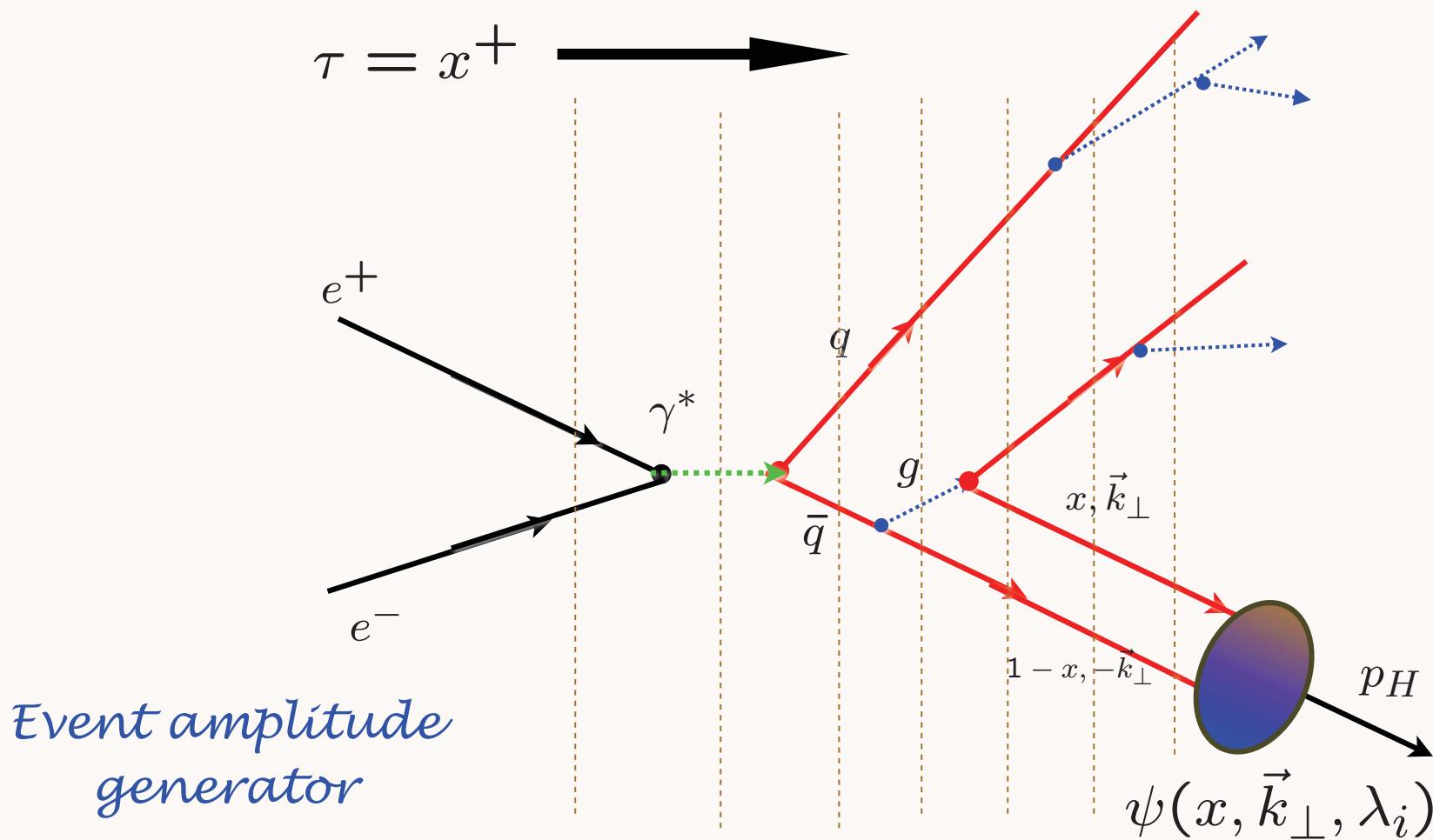
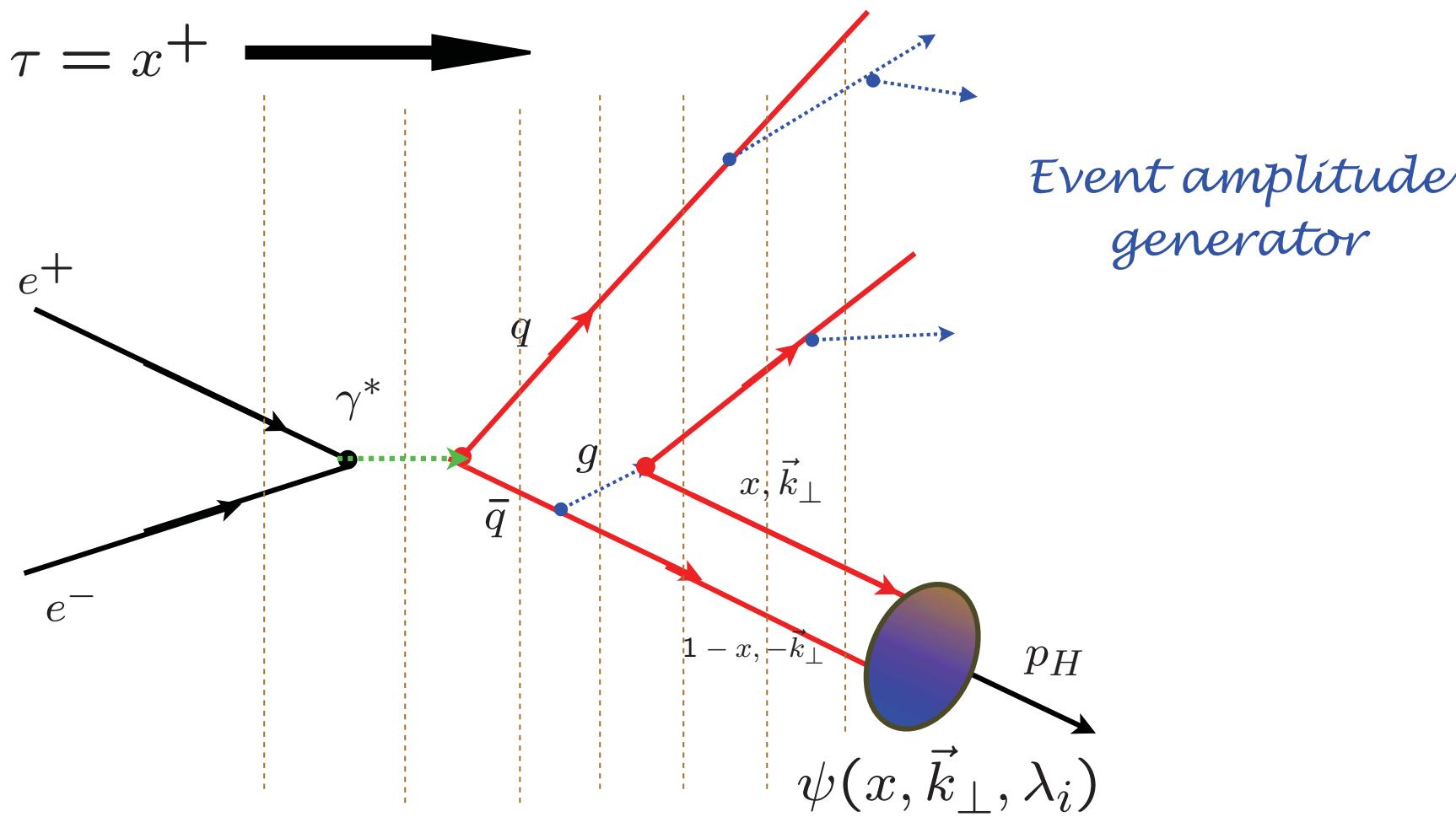


Hadronization at the Amplitude Level



Construct helicity amplitude using Light-Front Perturbation theory; coalesce quarks via LFWFs

Hadronization at the Amplitude Level



AdS/QCD

Hard Wall

Confinement: $\mathcal{M}^2 = \frac{k_\perp^2}{x(1-x)} < \Lambda_{QCD}^2$

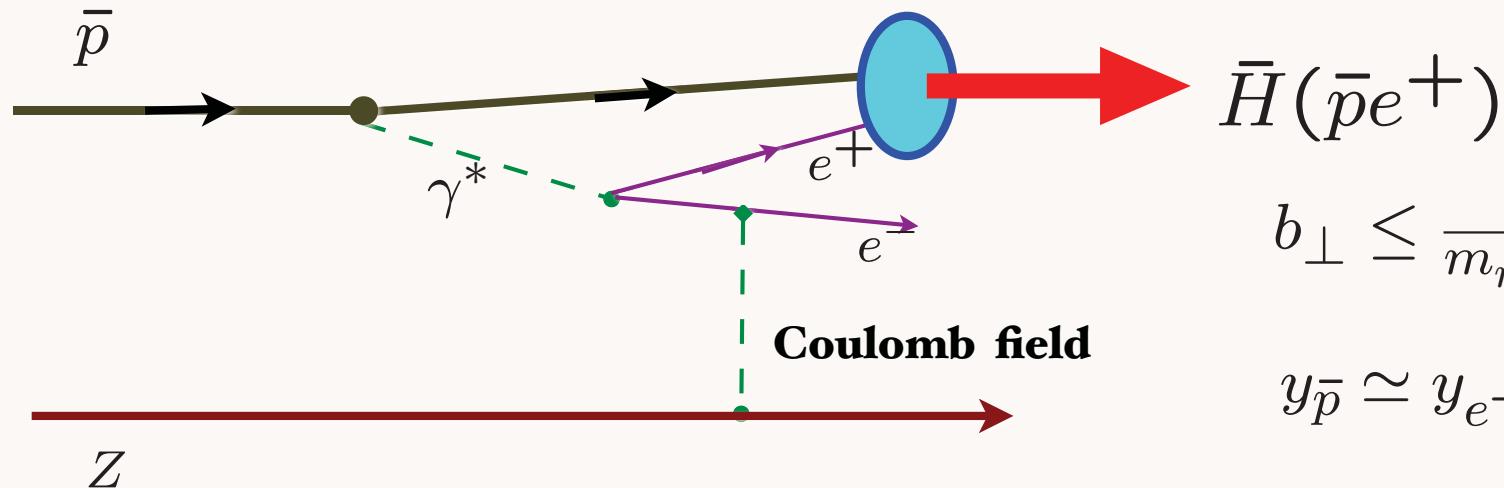
Capture if $\zeta^2 = x(1-x)b_\perp^2 > \frac{1}{\Lambda_{QCD}^2}$

i.e.,

Formation of Relativistic Anti-Hydrogen

Measured at CERN-LEAR and FermiLab

Munger, Schmidt, sjb



$$b_{\perp} \leq \frac{1}{m_{red}\alpha}$$

$$y_{\bar{p}} \simeq y_{e^+}$$

Coalescence of off-shell co-moving positron and antiproton.

Wavefunction maximal at small impact separation and equal rapidity

“Hadronization” at the Amplitude Level

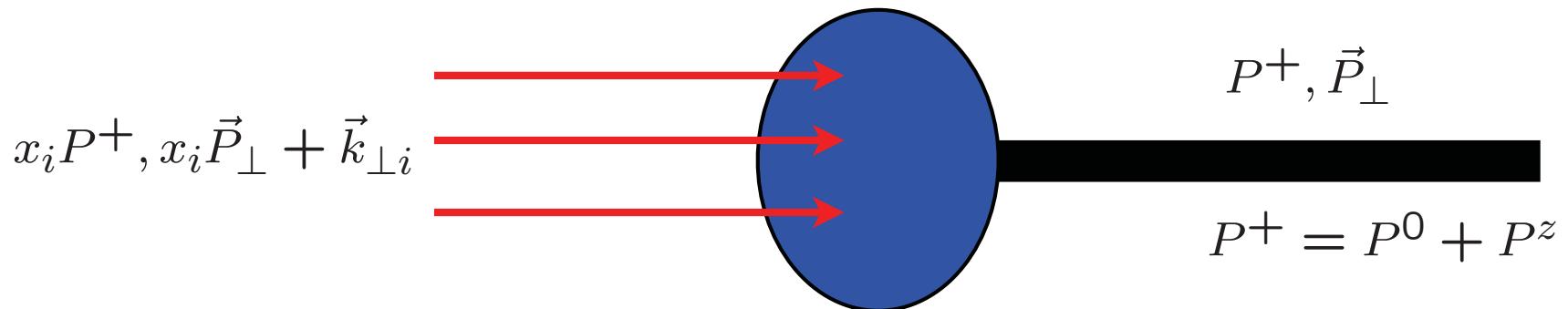
Features of LF T-Matrix Formalism

“Event Amplitude Generator”

- Coalesce color-singlet cluster to hadronic state if

$$\mathcal{M}_n^2 = \sum_{i=1}^n \frac{k_{\perp i}^2 + m_i^2}{x_i} < \Lambda_{QCD}^2$$

- The coalescence probability amplitude is the LF wavefunction $\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$
- No IR divergences: Maximal gluon and quark wavelength from confinement

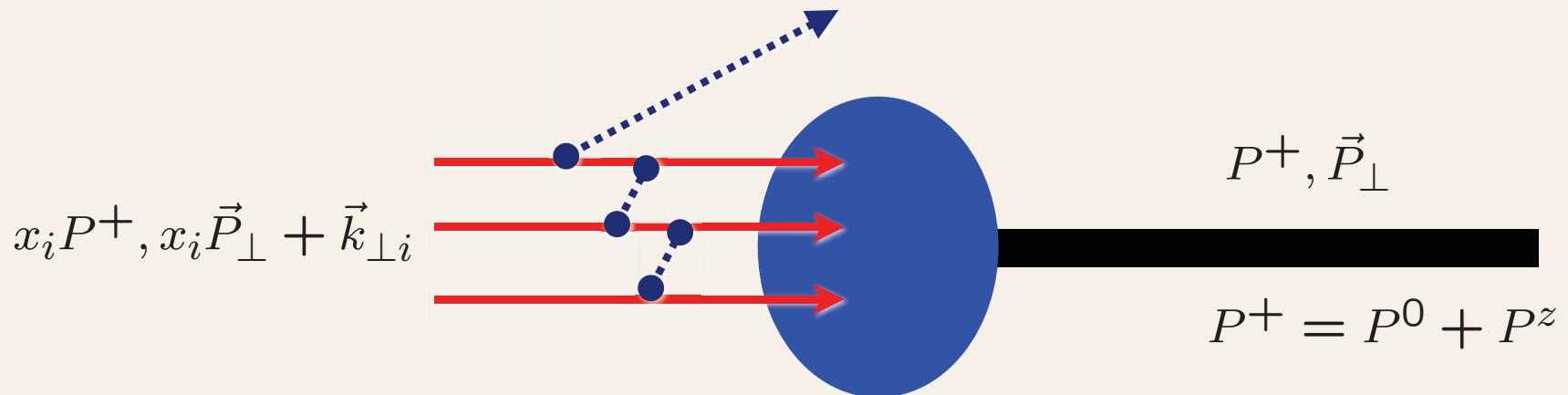


Features of LF T-Matrix Formalism

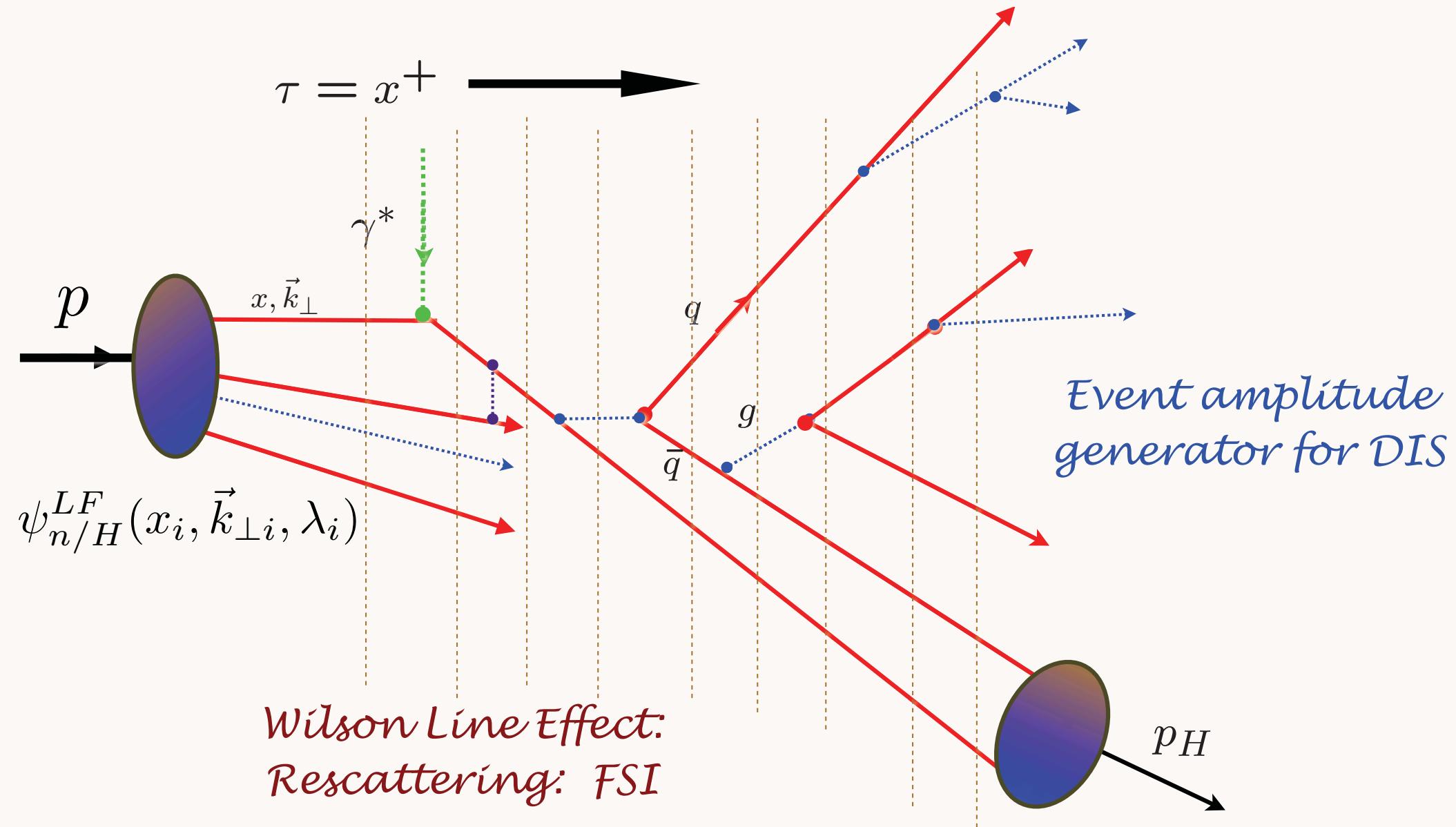
“Event Amplitude Generator”

If $\mathcal{M}_n^2 \geq \Lambda_{QCD}^2$ use PQCD hard gluon exchange

- DGLAP and ERBL Evolution from gluon emission and exchange
- Factorization Scale for structure functions and fragmentation functions set: $\mu_{fact} = \Lambda_{QCD}$



Hadronization at the Amplitude Level



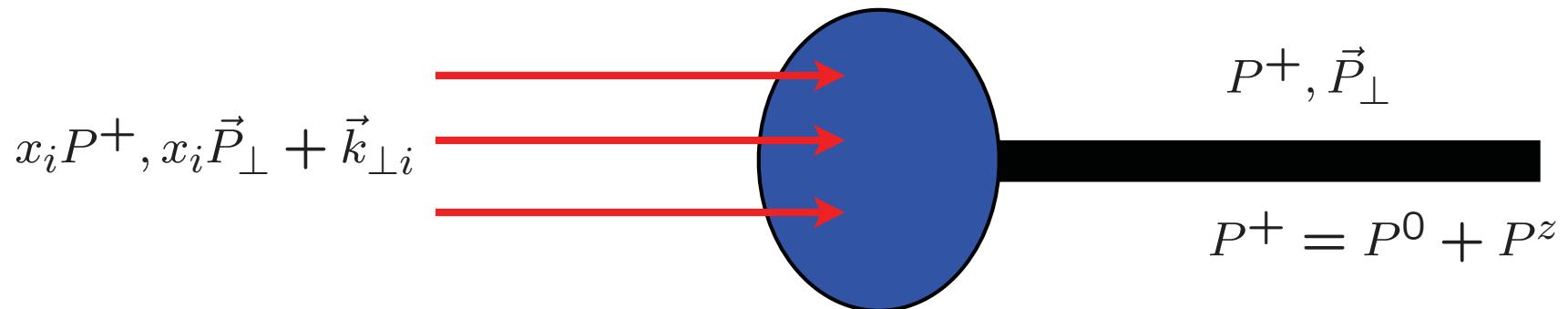
Features of LF T-Matrix Formalism

- Only positive + momenta; no backward time-ordered diagrams
- Frame-independent! Independent of P^+ and P^z
- LC gauge: No ghosts; physical helicity
- $J^z = L^z + S^z$ conservation at every vertex
- Sum all amplitudes with same initial-and final-state helicity, then square to get rate
- Renormalize each UV-divergent amplitude using “alternating denominator” method
- Multiple renormalization scales (BLM)

Features of LF T-Matrix Formalism

“Event Amplitude Generator”

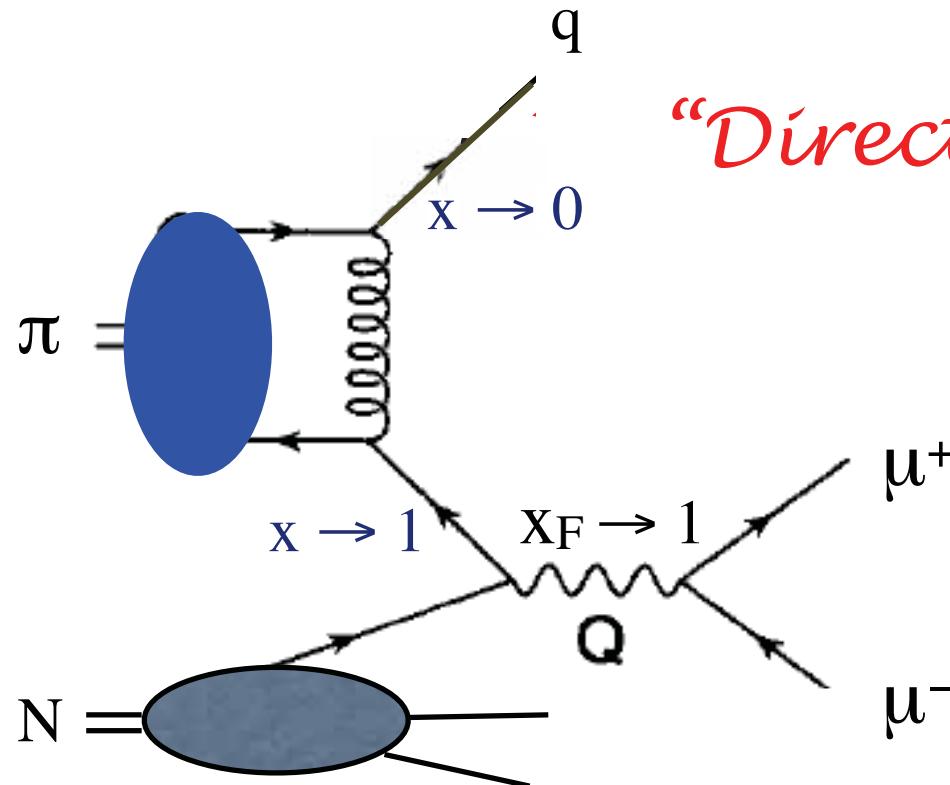
- Same principle as antihydrogen production: off-shell coalescence
- coalescence to hadron favored at equal rapidity, small transverse momenta
- leading heavy hadron production: D and B mesons produced at large z
- hadron helicity conservation if hadron LFWF has $L^z = 0$
- Baryon AdS/QCD LFWF has aligned and anti-aligned quark spin



$$\pi N \rightarrow \mu^+ \mu^- X \text{ at high } x_F$$

In the limit where $(1-x_F)Q^2$ is fixed as $Q^2 \rightarrow \infty$

Entire pion wf contributes to hard process



Virtual photon is longitudinally polarized

Berger, sjb
Khoze, Brandenburg, Muller, sjb
Hoyer Vanttinen

$\pi^- N \rightarrow \mu^+ \mu^- X$ at 80 GeV/c

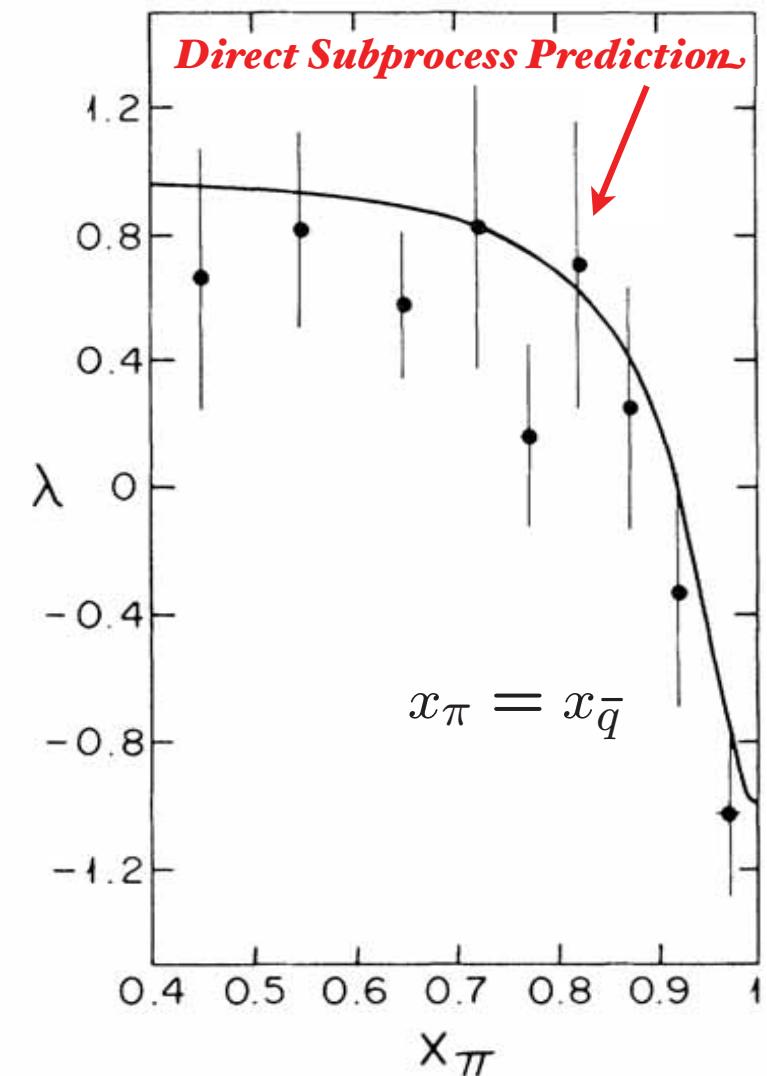
$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2\theta + \rho \sin 2\theta \cos \phi + \omega \sin^2\theta \cos 2\phi.$$

$$\frac{d^2\sigma}{dx_\pi d\cos\theta} \propto x_\pi \left((1-x_\pi)^2 (1+\cos^2\theta) + \frac{4}{9} \frac{\langle k_T^2 \rangle}{M^2} \sin^2\theta \right)$$

$$\langle k_T^2 \rangle = 0.62 \pm 0.16 \text{ GeV}^2/c^2$$

Dramatic change in angular distribution at large x_F

Example of a higher-twist direct subprocess



Chicago-Princeton
Collaboration

Phys.Rev.Lett.55:2649,1985

Crucial Test of Leading -Twist QCD: Scaling at fixed x_T

$$x_T = \frac{2p_T}{\sqrt{s}}$$

$$E \frac{d\sigma}{d^3 p}(pN \rightarrow \pi X) = \frac{F(x_T, \theta_{CM})}{p_T^{n_{eff}}}$$

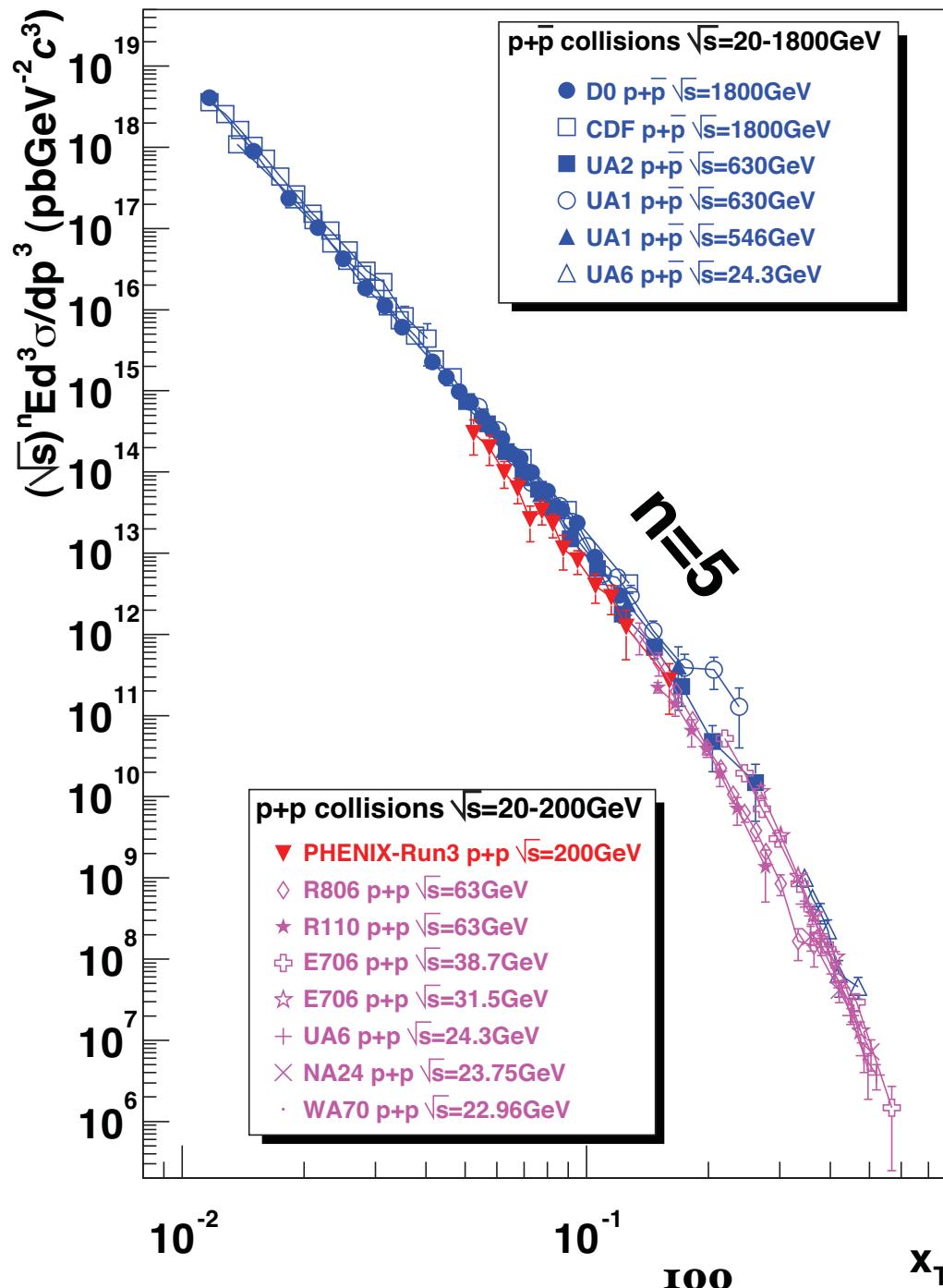
Parton model: $n_{eff} = 4$

As fundamental as Bjorken scaling in DIS

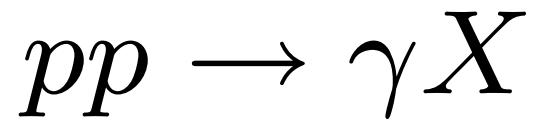
Conformal scaling: $n_{eff} = 2 n_{active} - 4$

$\sqrt{s}^n E \frac{d\sigma}{d^3 p}(pp \rightarrow \gamma X)$ at fixed x_T

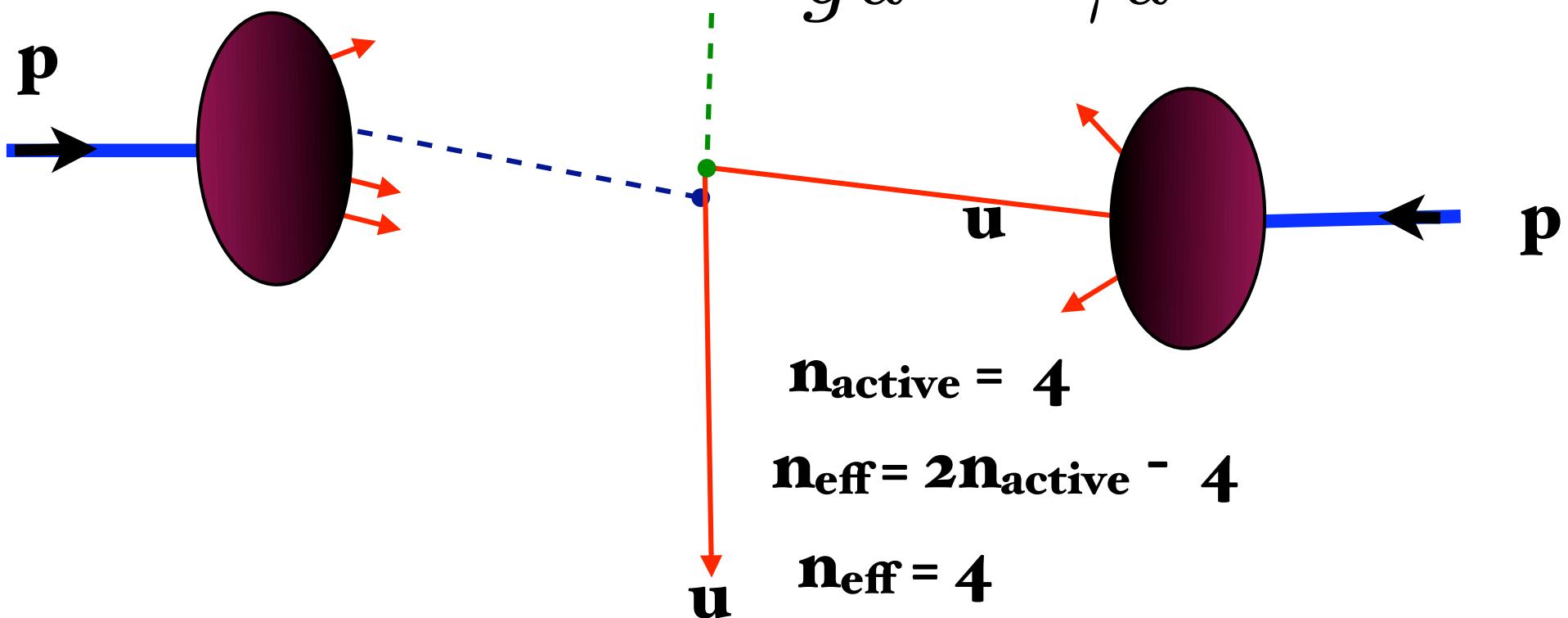
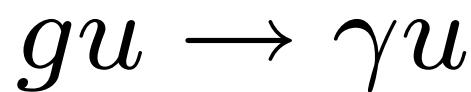
Tannenbaum



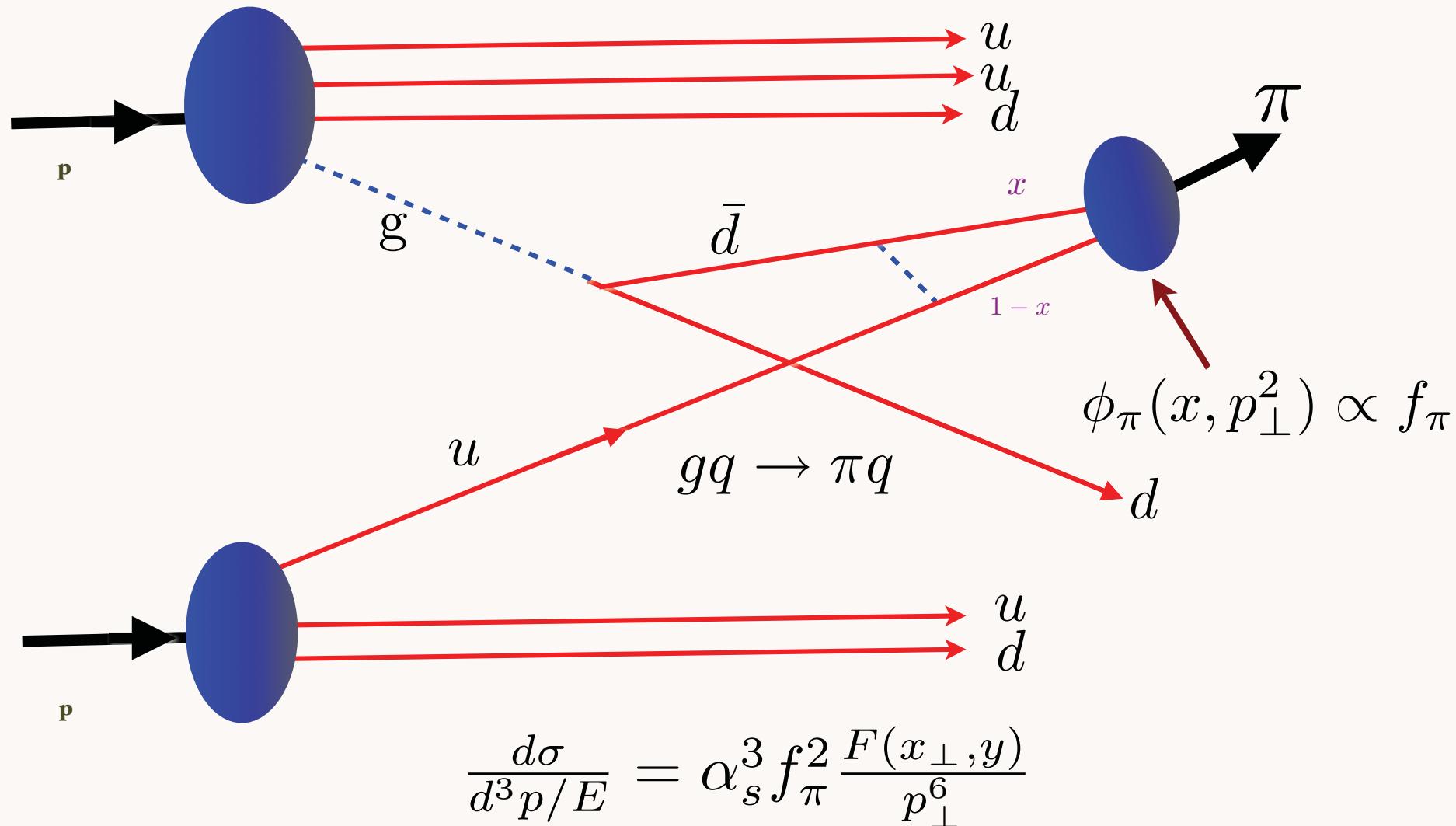
Scaling of direct
photon production
consistent with
PQCD



$$E \frac{d\sigma}{d^3 p}(pp \rightarrow \gamma X) = \frac{F(\theta_{cm}, x_T)}{p_T^4}$$

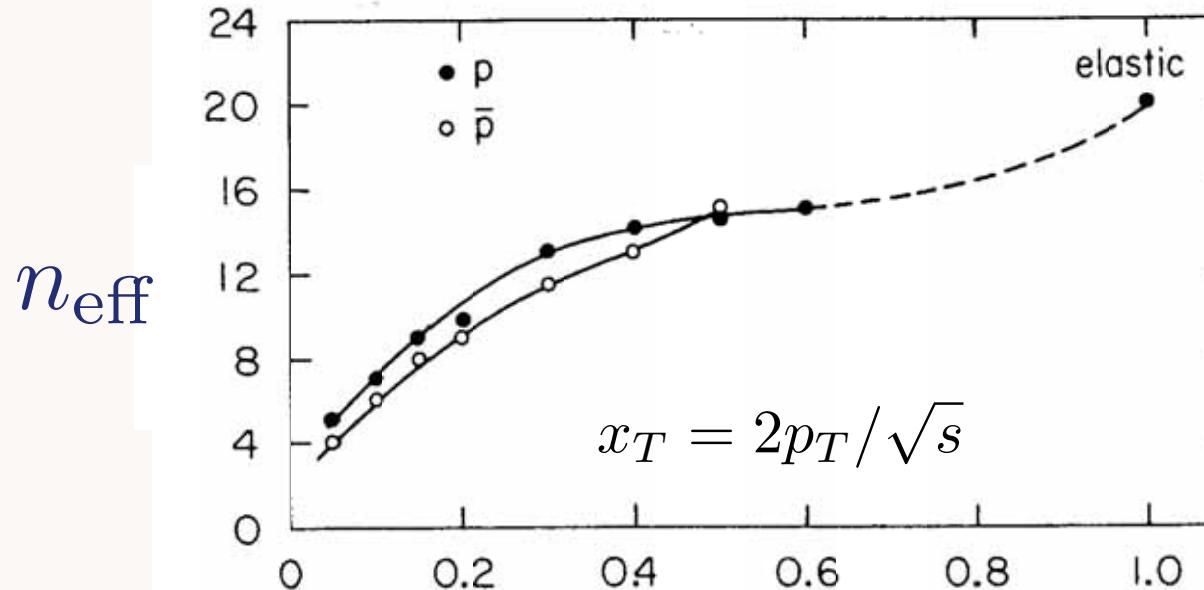


Higher-Twist Contribution to Hadron Production



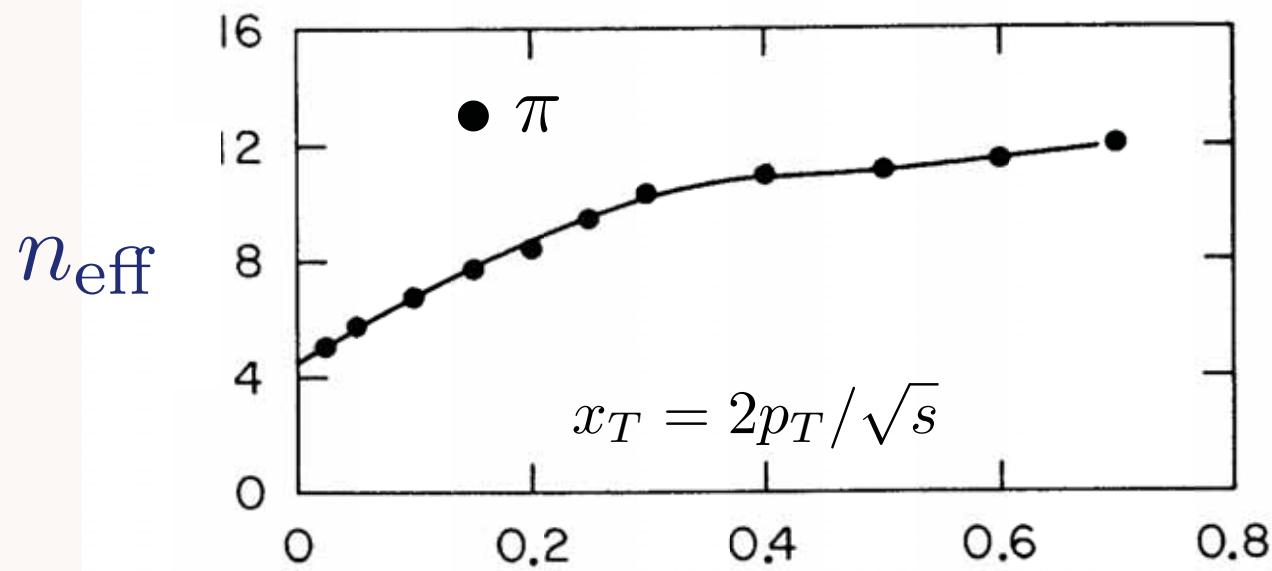
No Fragmentation Function

$$E \frac{d\sigma}{d^3 p} (pp \rightarrow HX) = \frac{F(x_T, \theta_{cm} = \pi/2)}{p_T^{n_{eff}}}$$



Clear evidence
for higher-twist
contributions

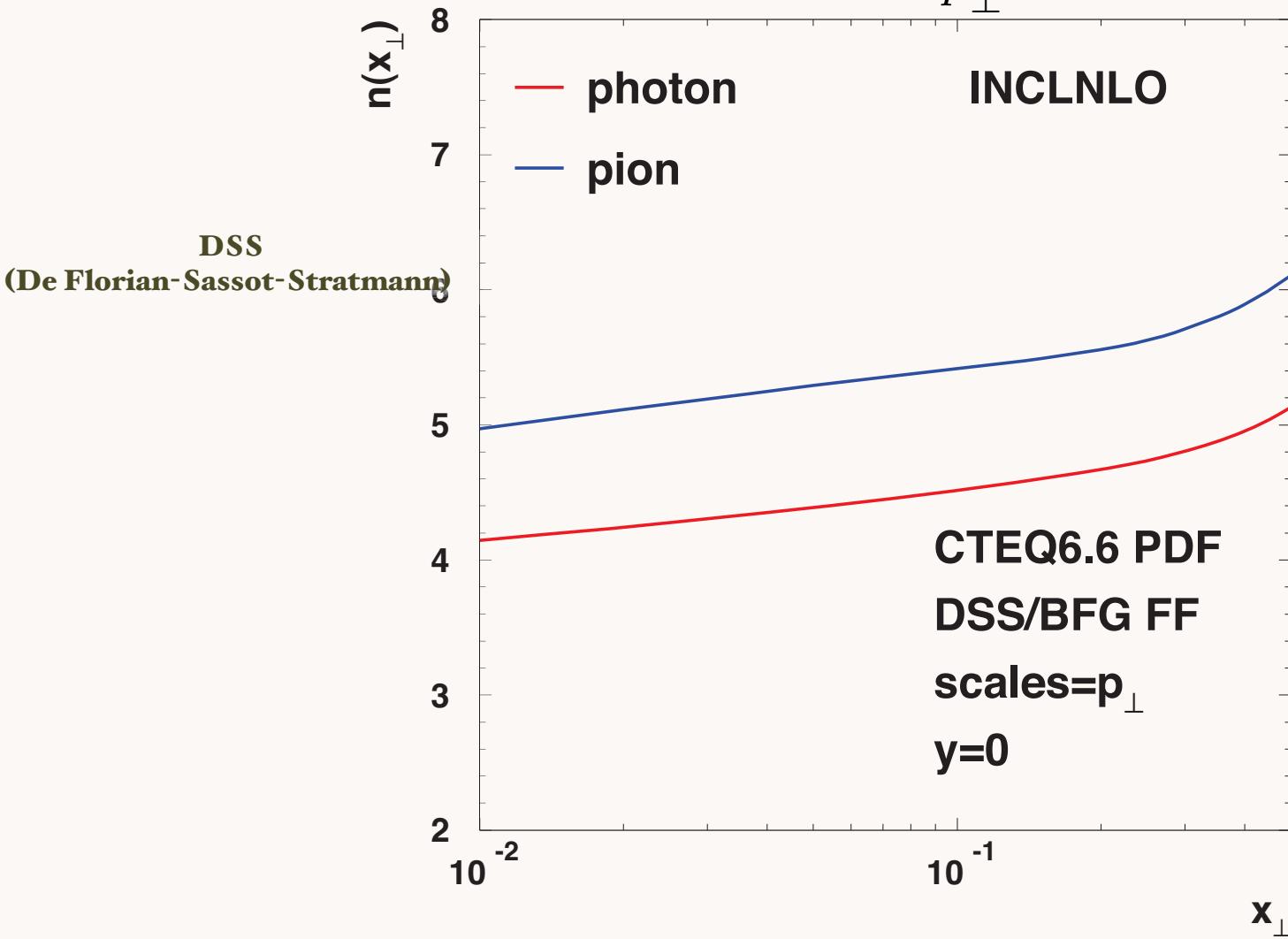
Fermilab, ISR data



Continuous
Rise of n_{eff}

QCD prediction: Modification of power fall-off due to DGLAP evolution and the Running Coupling

$$\frac{d\sigma}{d^3 p / E} = \frac{F(x_\perp, y)}{p_\perp^{n(x_\perp)}}$$



$pp \rightarrow \pi X$

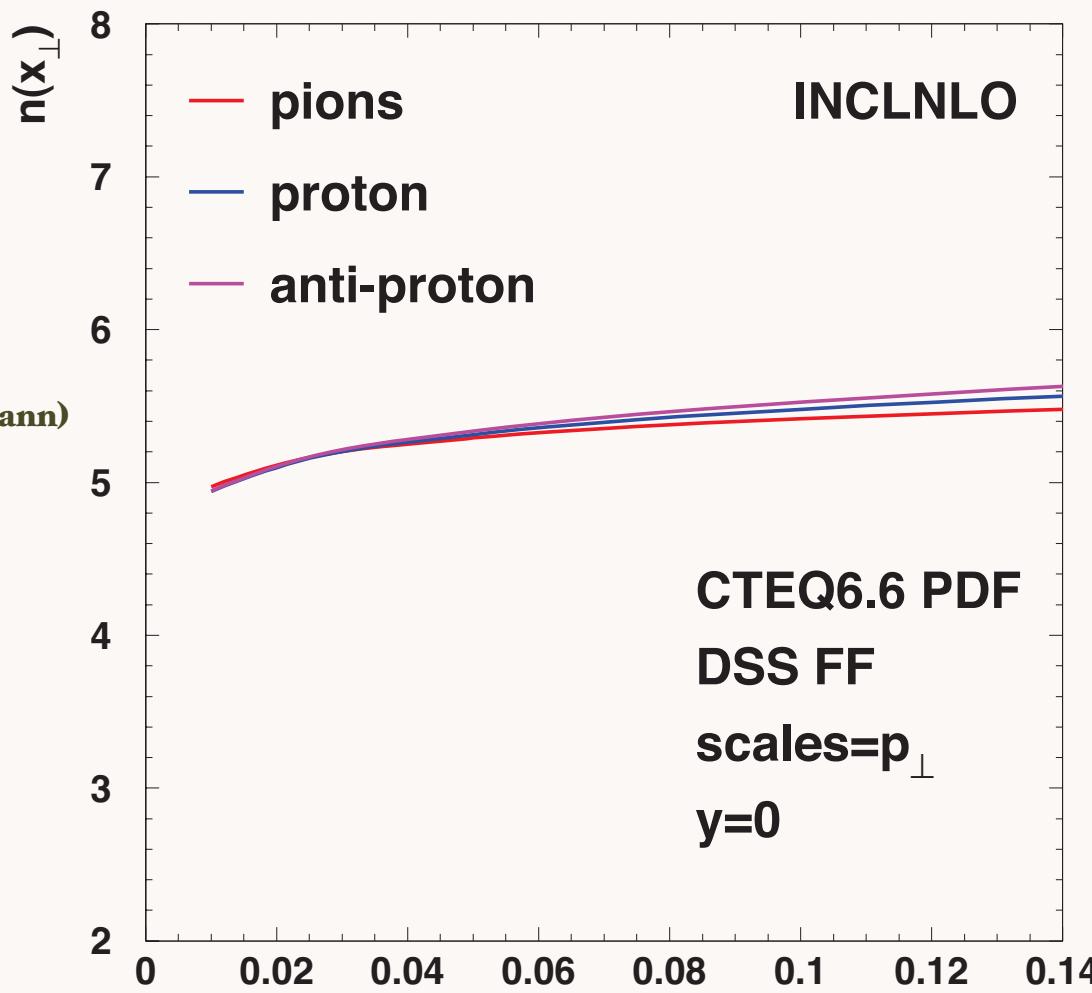
$pp \rightarrow \gamma X$

$5 < p_\perp < 20 \text{ GeV}$

$70 \text{ GeV} < \sqrt{s} < 4 \text{ TeV}$

$5 < p_\perp < 20 \text{ GeV}$ $70 \text{ GeV} < \sqrt{s} < 4 \text{ TeV}$

$pp \rightarrow \pi X$



$pp \rightarrow pX$

$pp \rightarrow \bar{p}X$

$$\frac{d\sigma}{d^3p/E} = \frac{F(x_\perp, y)}{p_\perp^{n(x_\perp)}}$$

Arleo, Aurenche

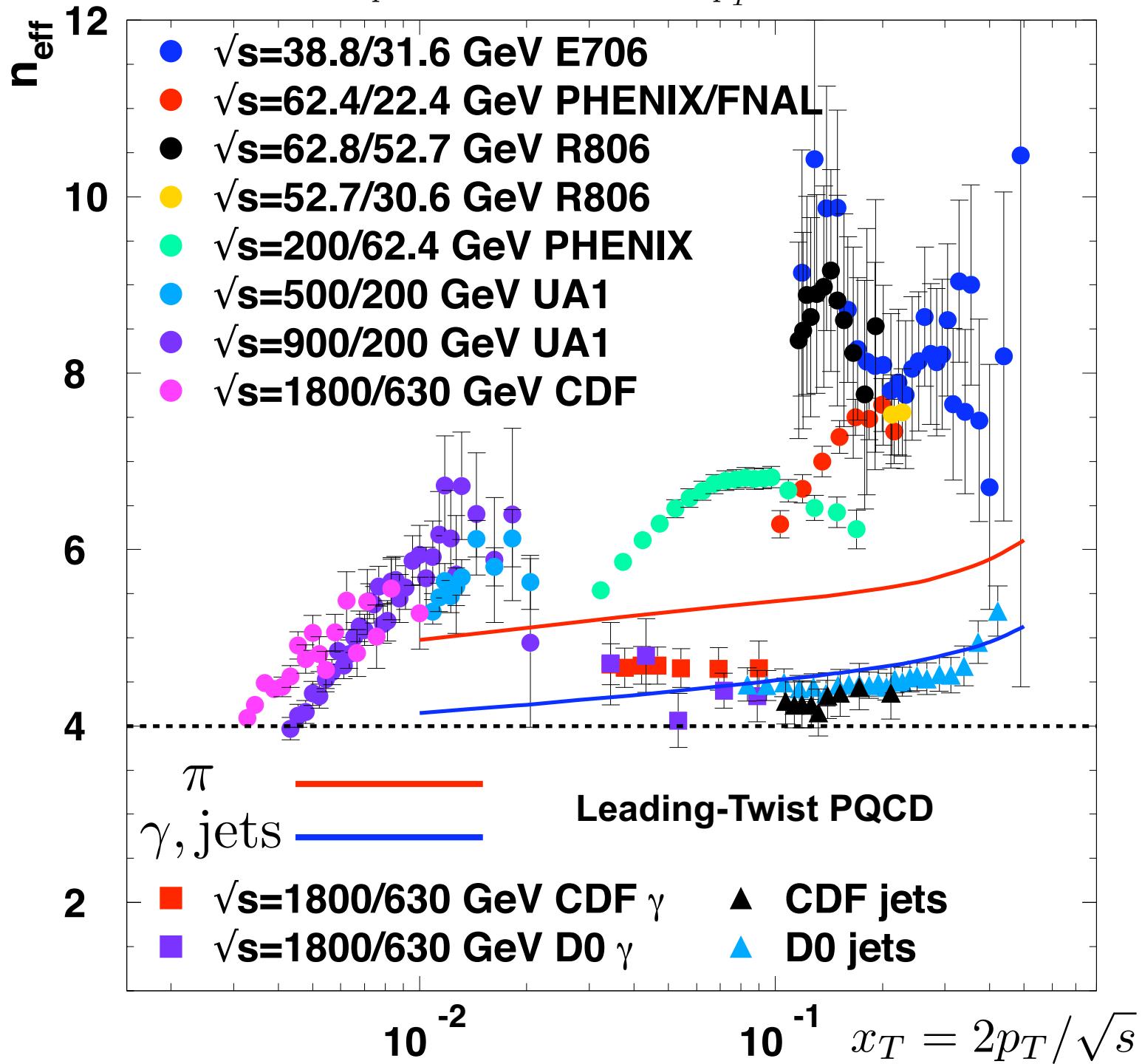
$$x_\perp = \frac{2p_\perp}{\sqrt{s}}$$

JTI Workshop ANL
April 16, 2009

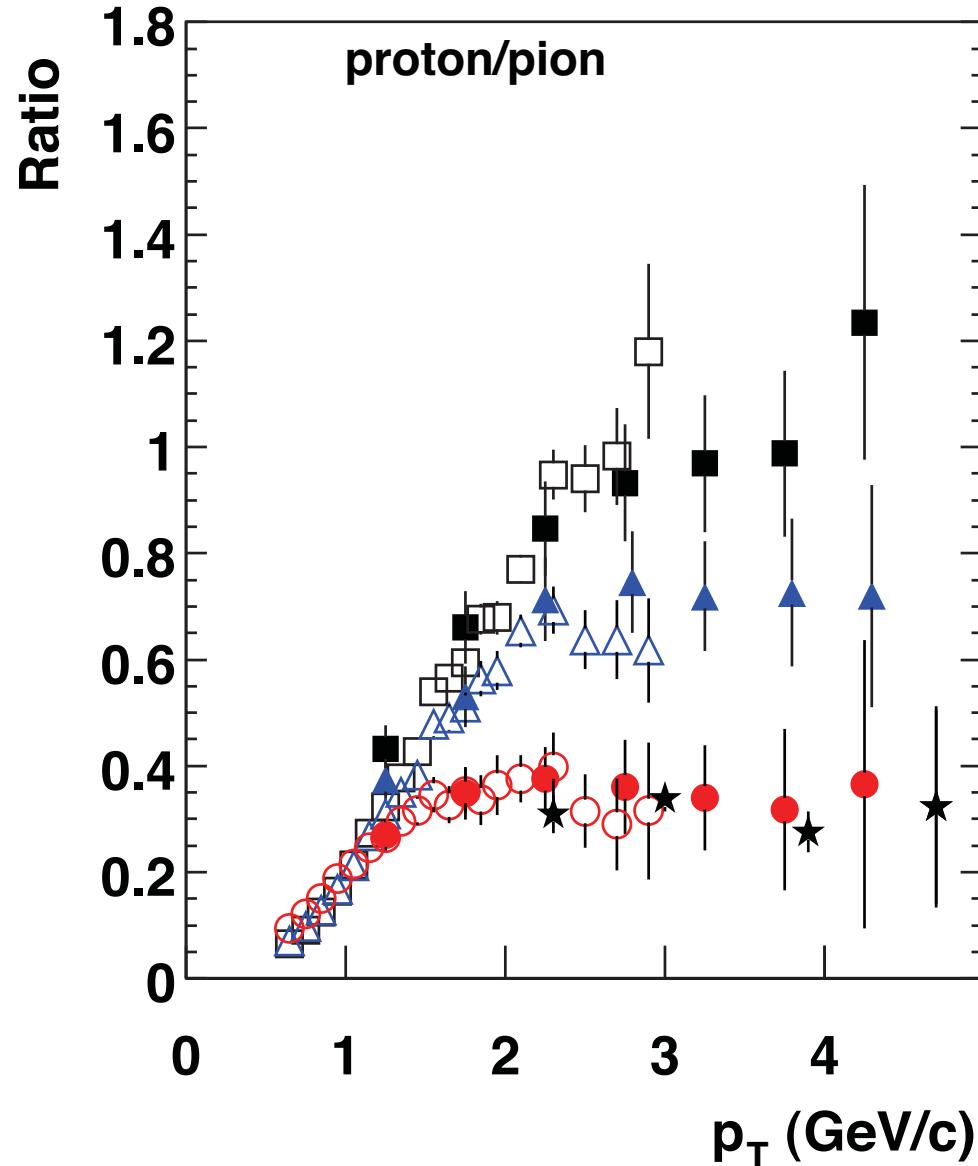
AdS/QCD and LF Holography

Stan Brodsky
SLAC

$$E \frac{d\sigma}{d^3 p}(pp \rightarrow HX) = \frac{F(x_T, \theta_{CM} = \pi/2)}{p_T^{n_{\text{eff}}}}$$



Baryon Anomaly: Particle ratio changes with centrality!



*Protons less absorbed
in nuclear collisions than pions*

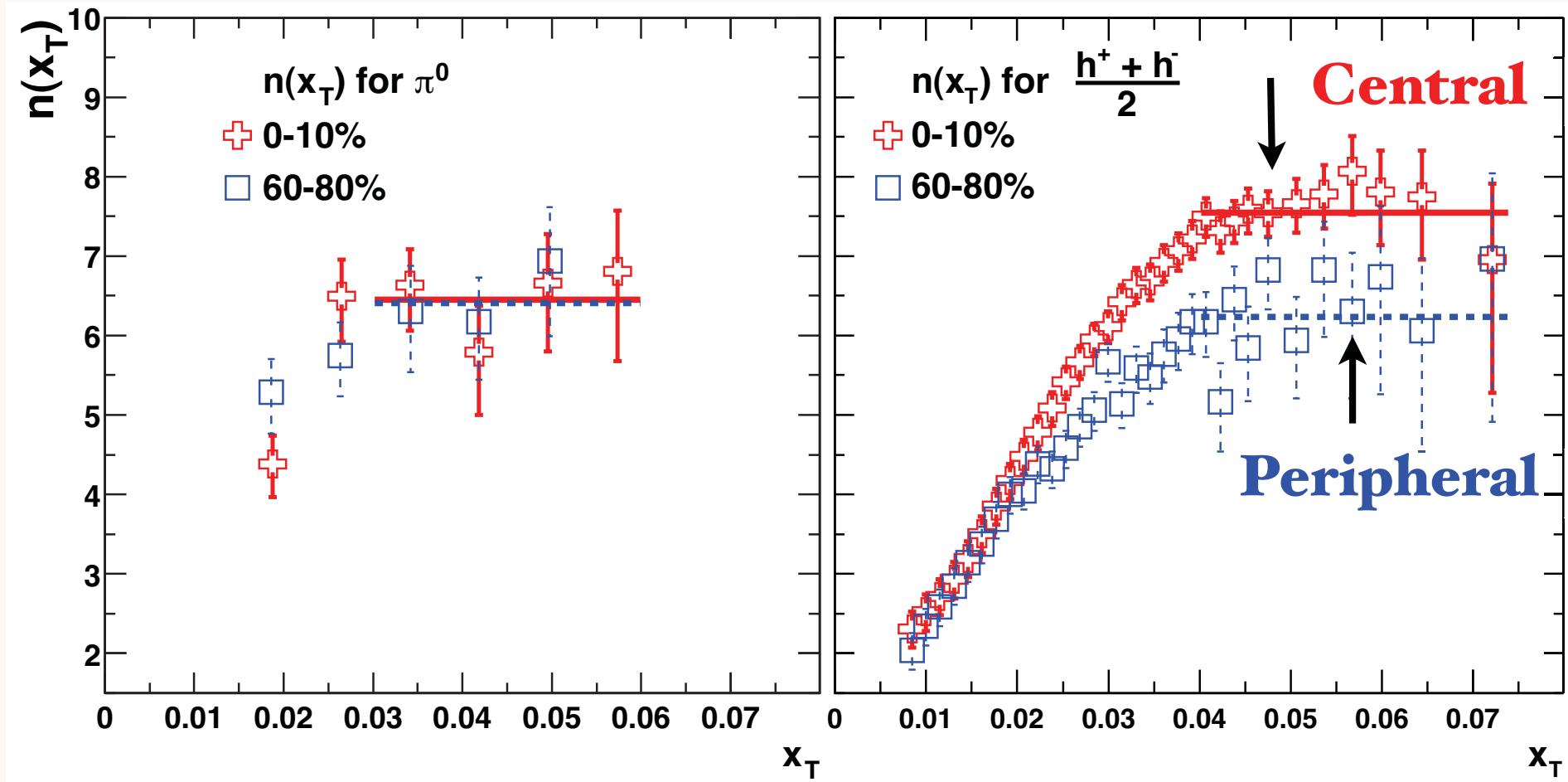
← Central

- ■ Au+Au 0-10%
- △ ▲ Au+Au 20-30%
- ● Au+Au 60-92%
- ★ p+p, $\sqrt{s} = 53$ GeV, ISR
- - - e⁺e⁻, gluon jets, DELPHI
- e⁺e⁻, quark jets, DELPHI

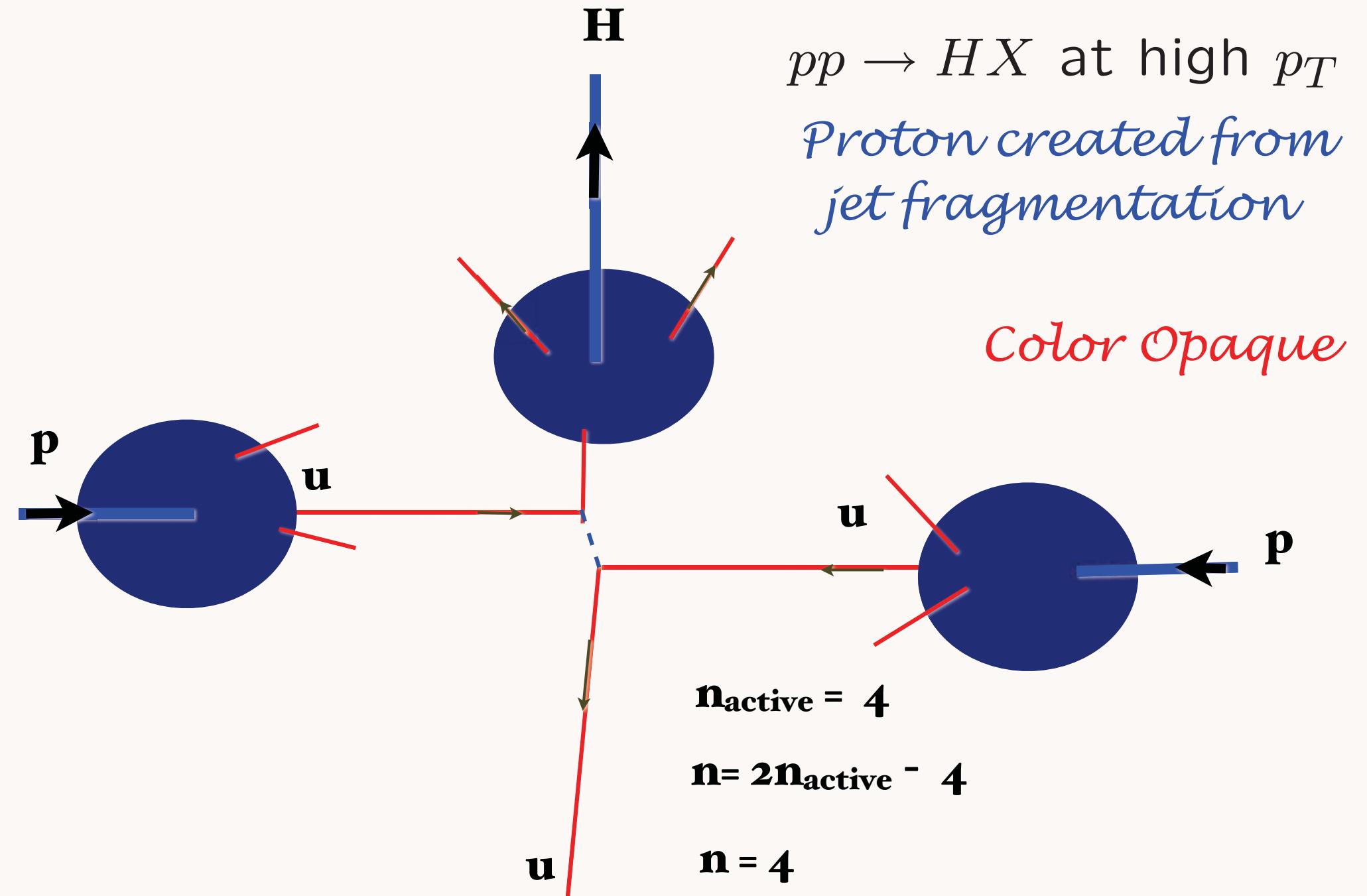
← Peripheral

Sickles, sjb

$$\sqrt{s_{NN}} = 130 \text{ and } 200 \text{ GeV}$$



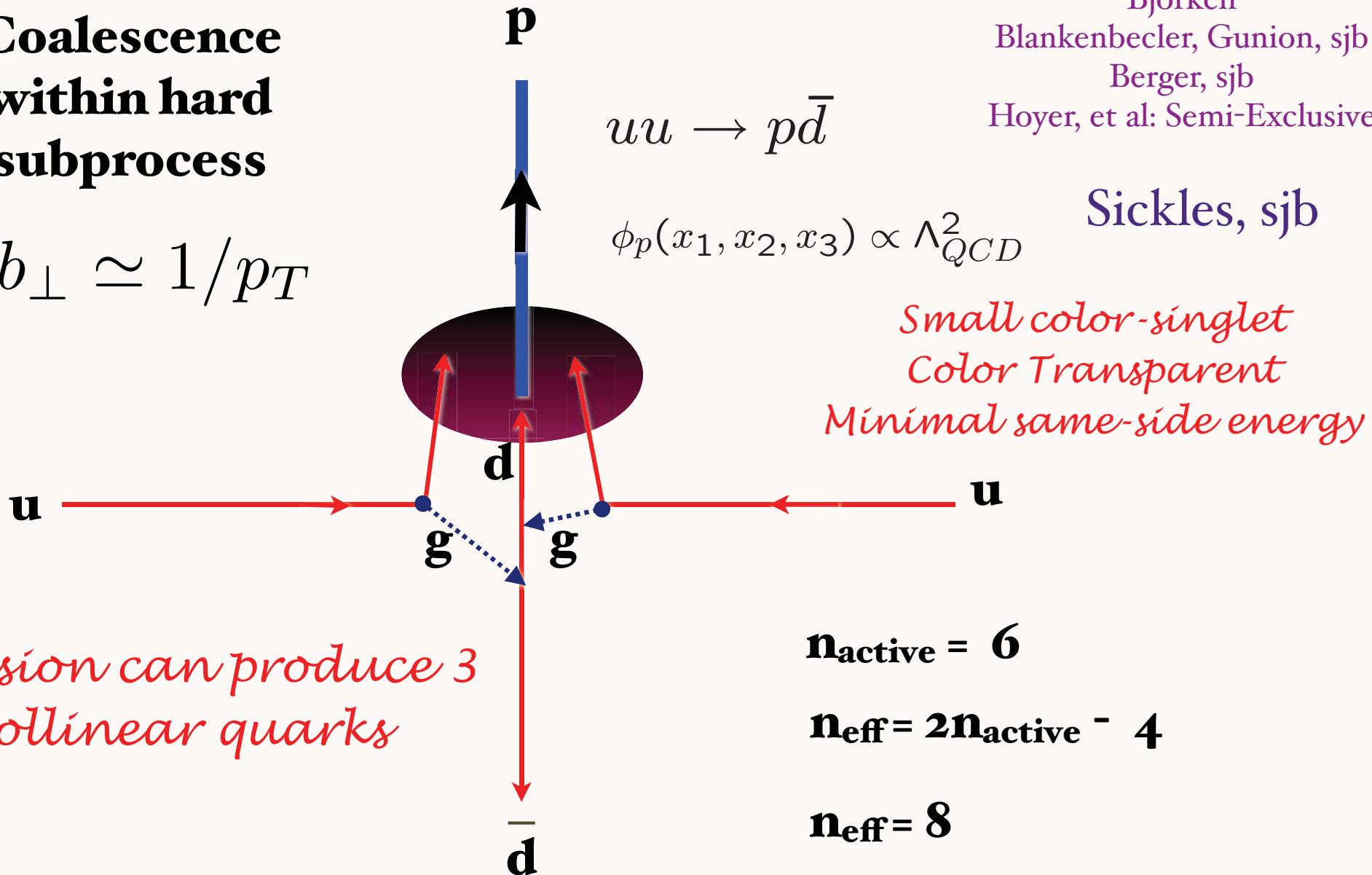
Proton power changes with centrality !



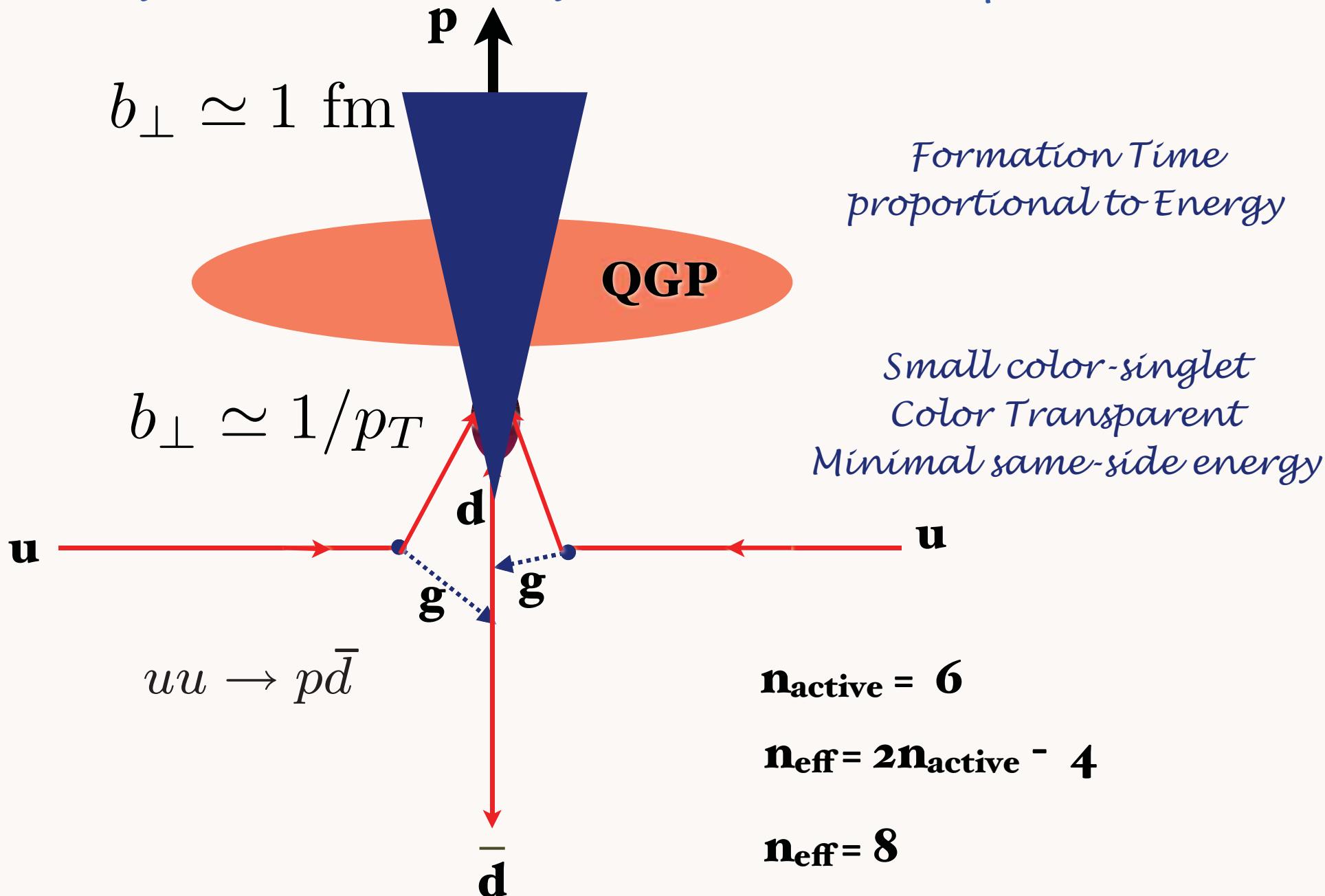
Baryon can be made directly within hard subprocess

Coalescence within hard subprocess

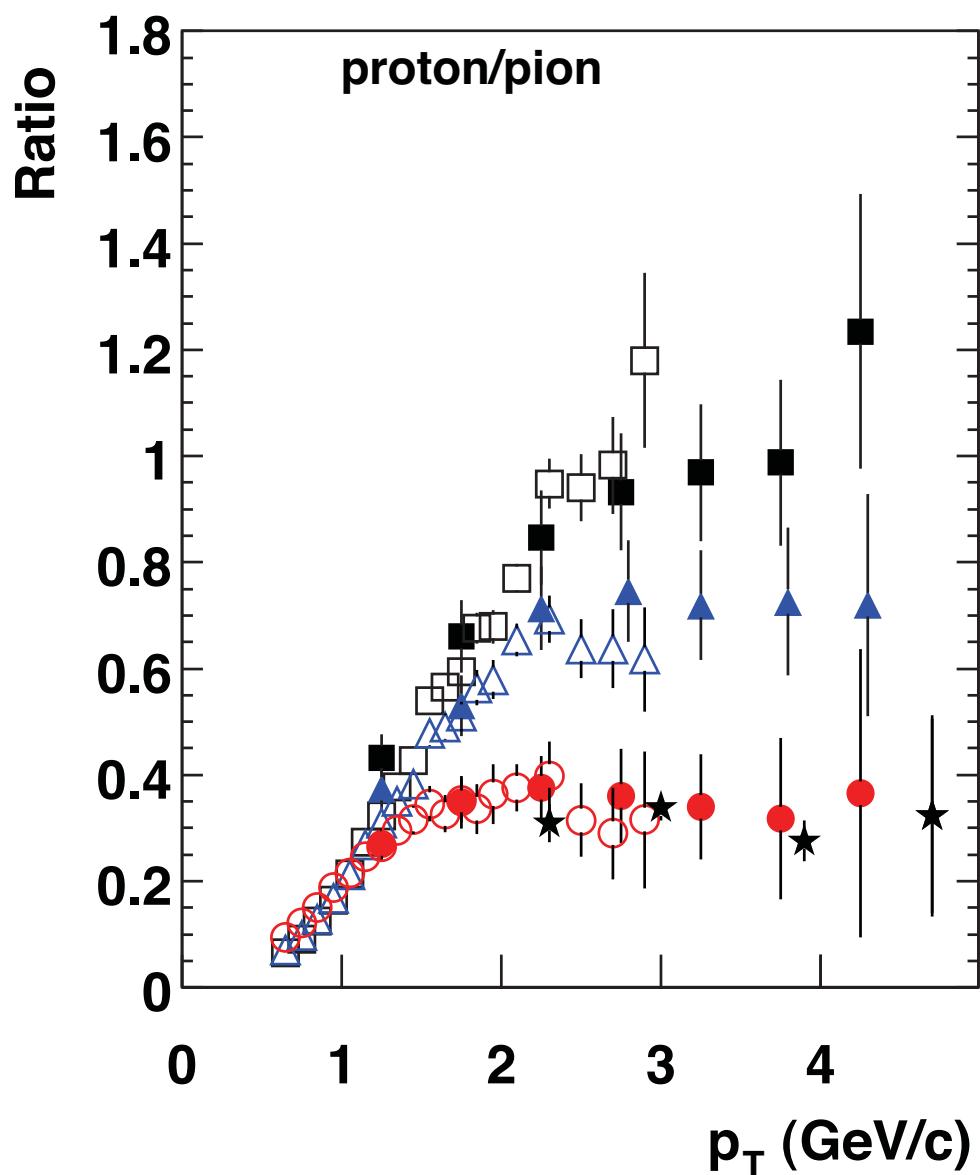
$$b_{\perp} \simeq 1/p_T$$



Baryon made directly within hard subprocess



Particle ratio changes with centrality!

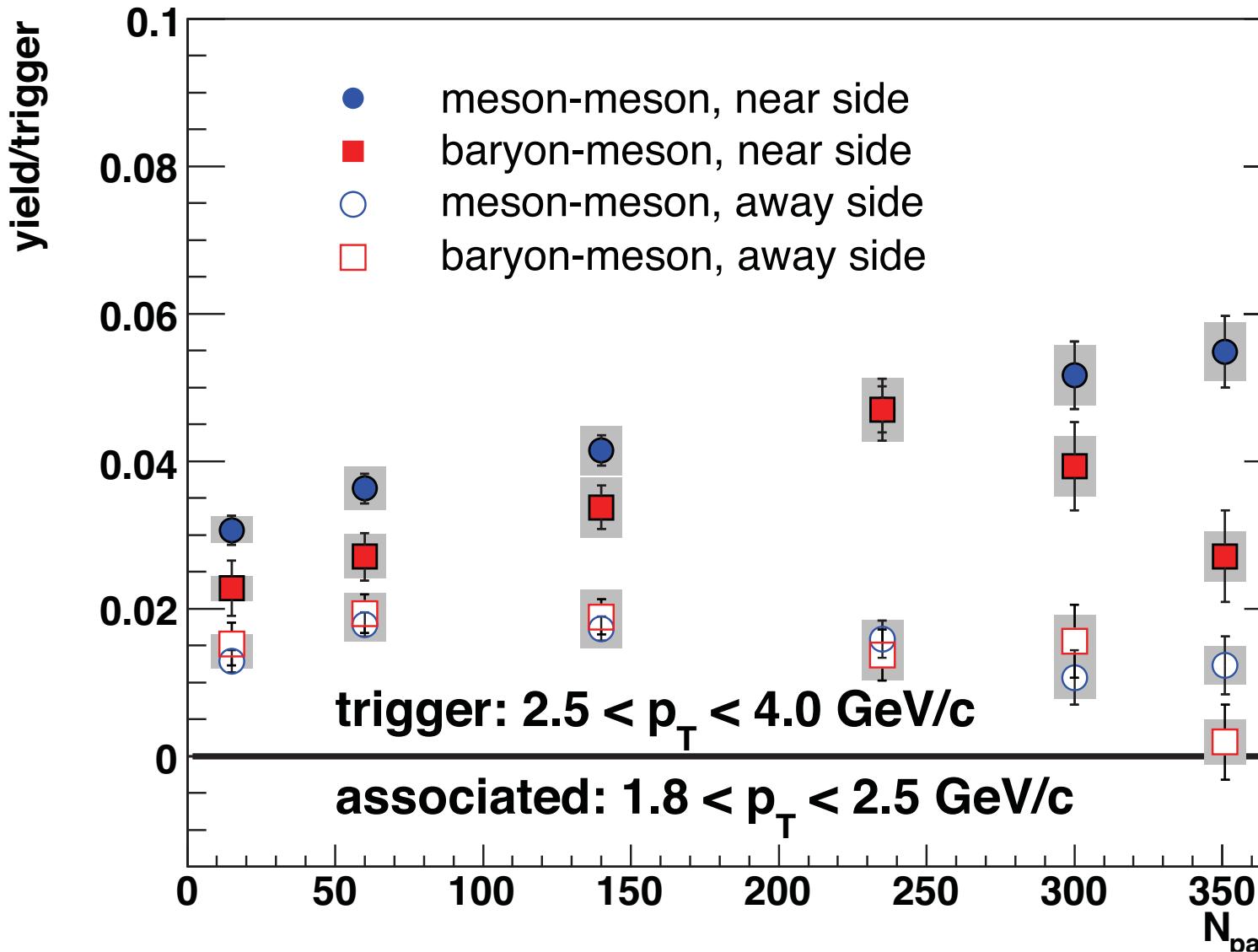


*Protons less absorbed
in nuclear collisions than pions
because of dominant,
color transparent higher twist process*

← Central

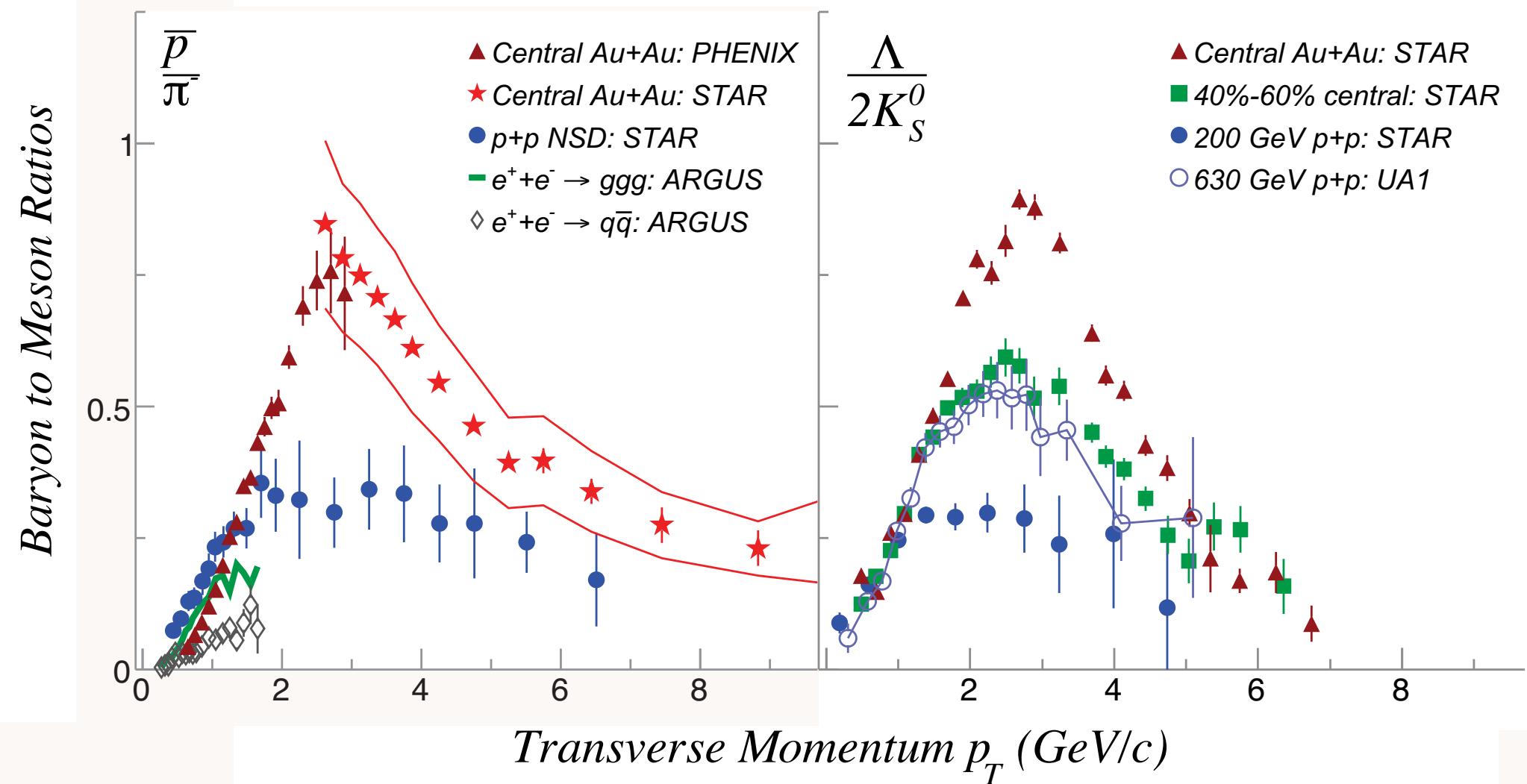
- ■ Au+Au 0-10%
- △ ▲ Au+Au 20-30%
- ● Au+Au 60-92%
- ★ p+p, $\sqrt{s} = 53$ GeV, ISR
- - - e⁺e⁻, gluon jets, DELPHI
- e⁺e⁻, quark jets, DELPHI

← Peripheral



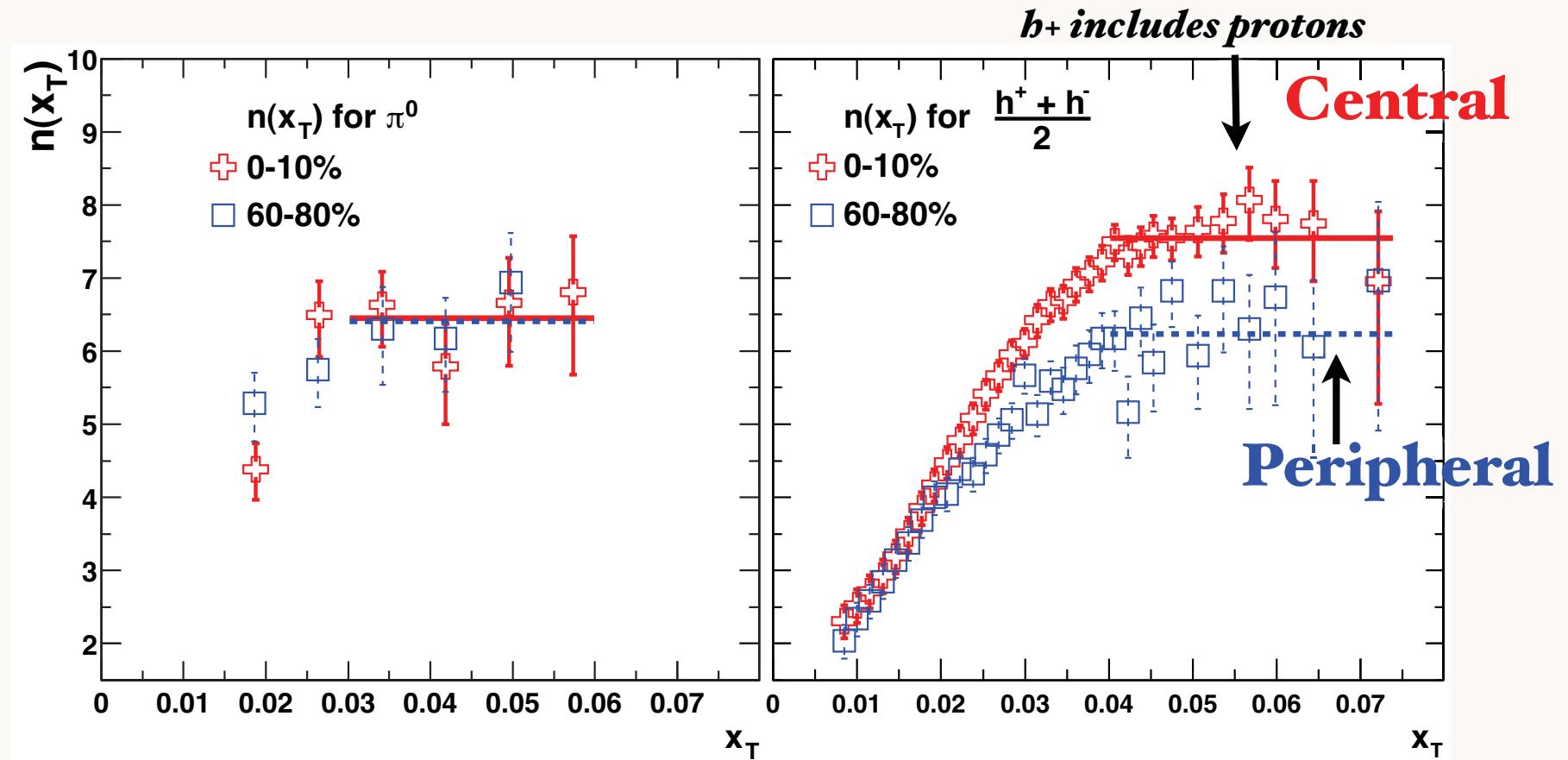
proton trigger:
same-side
particles
decreases with
centrality

Proton production more dominated by
color-transparent direct high- n_{eff} subprocesses



Power-law exponent $n(x_T)$ for π^0 and h spectra in central and peripheral Au+Au collisions at $\sqrt{s_{NN}} = 130$ and 200 GeV

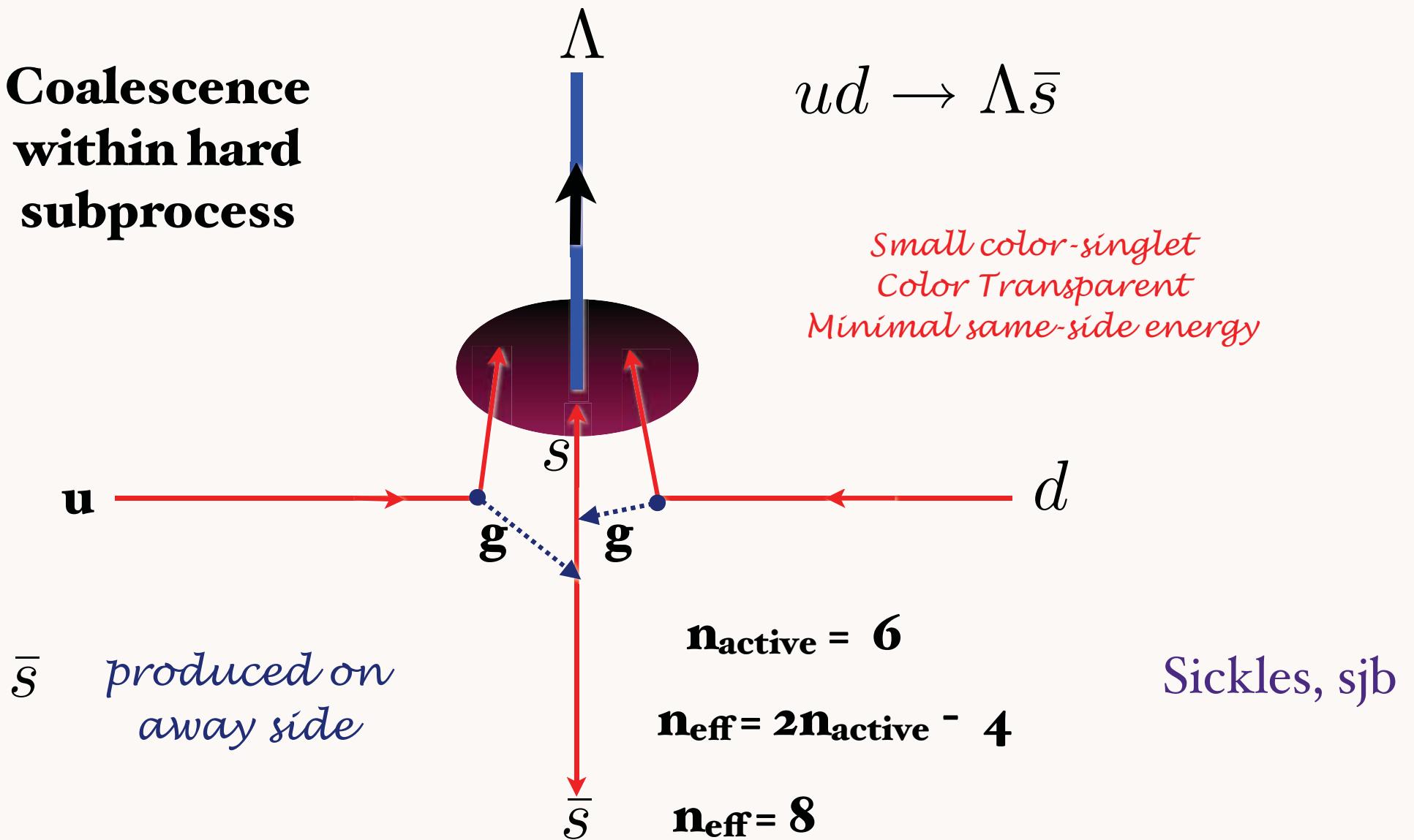
S. S. Adler, *et al.*, PHENIX Collaboration, *Phys. Rev. C* **69**, 034910 (2004) [nucl-ex/0308006].



Proton production dominated by
color-transparent direct high n_{eff} subprocesses

Lambda can be made directly within hard subprocess

Coalescence within hard subprocess



Baryon Anomaly: Evidence for Direct, Higher-Twist Subprocesses

- Explains anomalous power behavior at fixed x_T
- Protons more likely to come from direct higher-twist subprocess than pions
- Protons less absorbed than pions in central nuclear collisions because of color transparency
- Predicts increasing proton to pion ratio in central collisions
- Proton power n_{eff} increases with centrality since leading twist contribution absorbed
- Fewer same-side hadrons for proton trigger at high centrality
- Exclusive-inclusive connection at $x_T = 1$

Chiral Symmetry Breaking in AdS/QCD

We consider the action of the X field which encodes the effects of CSB in AdS/QCD:

$$S_X = \int d^4x dz \sqrt{g} \left(g^{\ell m} \partial_\ell X \partial_m X - \mu_X^2 X^2 \right), \quad (1)$$

with equations of motion

$$z^3 \partial_z \left(\frac{1}{z^3} \partial_z X \right) - \partial_\rho \partial^\rho X - \left(\frac{\mu_X R}{z} \right)^2 X = 0. \quad (2)$$

The zero mode has no variation along Minkowski coordinates

$$\partial_\mu X(x, z) = 0,$$

Ehrlich, Katz, Son, Stephanov

Babington, Erdmenger, Evans, Kirsch,
Guralnik, Thelfall

thus the equation of motion reduces to

$$[z^2 \partial_z^2 - 3z \partial_z + 3] X(z) = 0. \quad (3)$$

for $(\mu_X R)^2 = -3$, which corresponds to scaling dimension $\Delta_X = 3$. The solution is

$$X(z) = \langle X \rangle = Az + Bz^3, \quad (4)$$

where A and B are determined by the boundary conditions.

de Teramond, Shrock, sjb
(preliminary)

$A \propto m_q$ $B \propto \langle \bar{\psi} \psi \rangle$
Expectation value taken inside hadron

In presence of quark masses the Holographic LF wave equation is ($\zeta = z$)

$$\left[-\frac{d^2}{d\zeta^2} + V(\zeta) + \frac{X^2(\zeta)}{\zeta^2} \right] \phi(\zeta) = \mathcal{M}^2 \phi(\zeta), \quad (1)$$

and thus

$$\boxed{\delta M^2 = \left\langle \frac{X^2}{\zeta^2} \right\rangle.} \quad (2)$$

The parameter a is determined by the Weisberger term

$$a = \frac{2}{\sqrt{x}}.$$

Thus

$$\boxed{X(z) = \frac{m}{\sqrt{x}} z - \sqrt{x} \langle \bar{\psi} \psi \rangle z^3,} \quad (3)$$

and

$$\delta M^2 = \sum_i \left\langle \frac{m_i^2}{x_i} \right\rangle - 2 \sum_i m_i \langle \bar{\psi} \psi \rangle \langle z^2 \rangle + \langle \bar{\psi} \psi \rangle^2 \langle z^4 \rangle, \quad (4)$$

where we have used the sum over fractional longitudinal momentum $\sum_i x_i = 1$.

Mass shift from dynamics inside hadronic boundary

Chiral Symmetry Breaking in AdS/QCD

- Chiral symmetry breaking effect in AdS/QCD depends on weighted z^2 distribution, not constant condensate

$$\delta M^2 = -2m_q \langle \bar{\psi}\psi \rangle \times \int dz \phi^2(z) z^2$$

- z^2 weighting consistent with higher Fock states at periphery of hadron wavefunction
- AdS/QCD supports confined condensate picture

de Teramond, Shrock, sjb

Quark and Gluon condensates reside within hadrons, not vacuum

Casher and Susskind

Roberts et al.

Shrock and sjb

- **Bound-State Dyson-Schwinger Equations**
Roberts et al.
- **LF vacuum trivial up to $k^+ = 0$ zero modes**
- **Analogous to finite size superconductor**
- **Implications for cosmological constant --
Eliminates 45 orders of magnitude conflict**
Shrock and sjb

Pion mass and decay constant.

[Pieter Maris](#), [Craig D. Roberts \(Argonne, PHY\)](#) , [Peter C. Tandy \(Kent State U.\)](#) . ANL-PHY-8753-TH-97, KSUCNR-103-97, Jul 1997. 12pp.

Published in **Phys.Lett.B420:267-273,1998**.

e-Print: [nucl-th/9707003](#)

Pi- and K meson Bethe-Salpeter amplitudes.

[Pieter Maris](#), [Craig D. Roberts \(Argonne, PHY\)](#) . ANL-PHY-8788-TH-97, Aug 1997. 34pp.

Published in **Phys.Rev.C56:3369-3383,1997**.

e-Print: [nucl-th/9708029](#)

Concerning the quark condensate.

[K. Langfeld \(Tubingen U.\)](#) , [H. Markum \(Vienna, Tech. U.\)](#) , [R. Pullirsch \(Regensburg U.\)](#) , [C.D. Roberts \(Argonne, PHY & Rostock U.\)](#) , [S.M. Schmidt \(Tubingen U. & HGF, Bonn\)](#) . ANL-PHY-10460-TH-2002, MPG-VT-UR-239-02, Jan 2003. 7pp.

Published in **Phys.Rev.C67:065206,2003**.

e-Print: [nucl-th/0301024](#)

“In-Meson Condensate”

$$-\langle \bar{q}q \rangle_{\zeta}^{\pi} = f_{\pi} \langle 0 | \bar{q} \gamma_5 q | \pi \rangle .$$

Valid even for $m_q \rightarrow 0$

f_{π} nonzero

QCD Symmetries

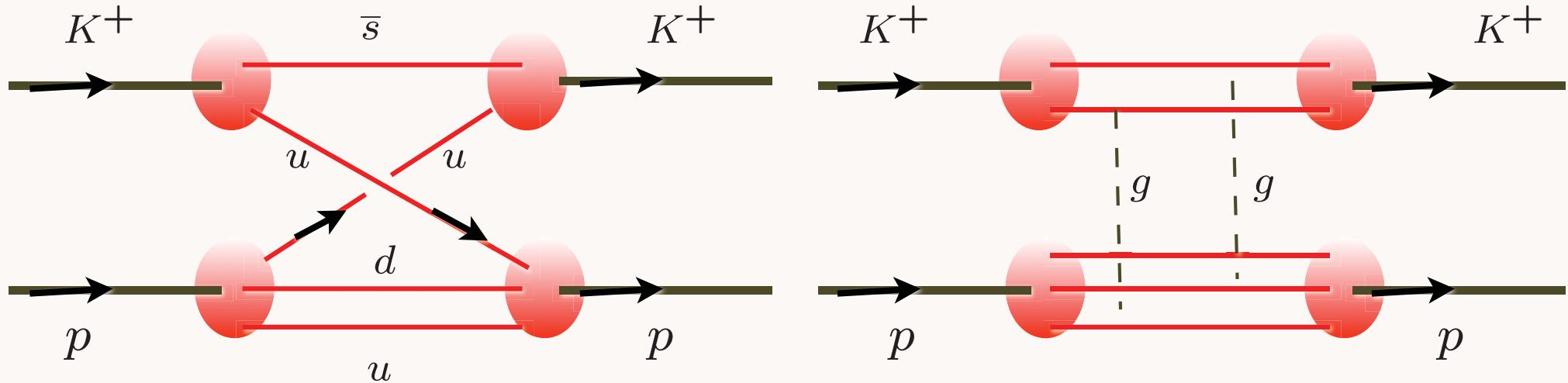
- Color Confinement: Maximum Wavelength of Quark and Gluons
- Conformal symmetry of QCD coupling in IR
- Provides Conformal Template
- Motivation for AdS/QCD
- QCD Condensates inside of hadronic LFWFs
- Technicolor: confined condensates inside of technihadrons -- alternative to Higgs
- Simple physical solution to cosmological constant conflict

Hadron Dynamics at the Amplitude Level

- LFWFs are the universal hadronic amplitudes which underlie structure functions, GPDs, exclusive processes, distribution amplitudes, direct subprocesses, hadronization.
- Relation of spin, momentum, and other distributions to physics of the hadron itself.
- Connections between observables, orbital angular momentum
- Role of FSI and ISIs--Sivers effect
- Higher Fock States give GMOR Relations, Chiral Symmetry Breaking

New Perspectives for QCD from AdS/CFT

- LFWFs: Fundamental frame-independent description of hadrons at amplitude level
- Holographic Model from AdS/CFT : Confinement at large distances and conformal behavior at short distances
- Model for LFWFs, meson and baryon spectra: many applications!
- New basis for diagonalizing Light-Front Hamiltonian
- Physics similar to MIT bag model, but covariant. No problem with support $0 < x < 1$.
- Quark Interchange dominant force at short distances



*Quark Interchange
(Spin exchange in atom-atom scattering)*

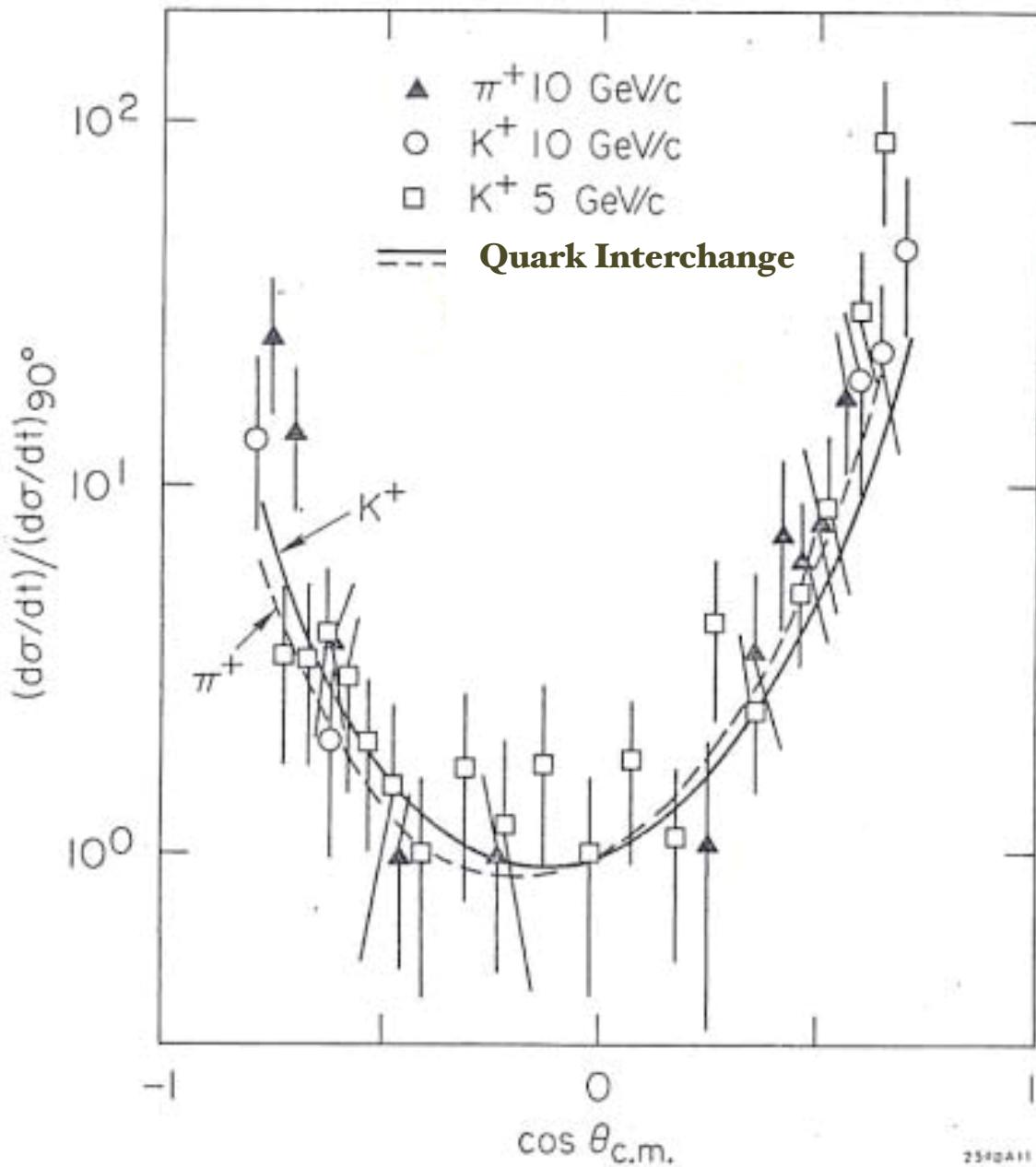
*Gluon Exchange
(Van der Waal -- Landshoff)*

$$\frac{d\sigma}{dt} = \frac{|M(s,t)|^2}{s^2}$$

$$M(t, u)_{\text{interchange}} \propto \frac{1}{ut^2}$$

$$M(s, t)_{\text{gluonexchange}} \propto s F(t)$$

MIT Bag Model (de Tar), large N_c , ('t Hooft), AdS/CFT all predict dominance of quark interchange:



*AdS/CFT explains why
quark interchange is
dominant interaction
at high momentum
transfer in exclusive
reactions*

$$M(t, u)_{\text{interchange}} \propto \frac{1}{ut^2}$$

Non-linear Regge behavior:

$$\alpha_R(t) \rightarrow -1$$

Why is quark-interchange dominant over gluon exchange?

Example: $M(K^+ p \rightarrow K^+ p) \propto \frac{1}{ut^2}$

Exchange of common u quark

$$M_{QIM} = \int d^2 k_\perp dx \psi_C^\dagger \psi_D^\dagger \Delta \psi_A \psi_B$$

Holographic model (Classical level):

Hadrons enter 5th dimension of AdS_5

Quarks travel freely within cavity as long as separation $z < z_0 = \frac{1}{\Lambda_{QCD}}$

LFWFs obey conformal symmetry producing quark counting rules.

Comparison of Exclusive Reactions at Large t

B. R. Baller,^(a) G. C. Blazey,^(b) H. Courant, K. J. Heller, S. Heppelmann,^(c) M. L. Marshak,
E. A. Peterson, M. A. Shupe, and D. S. Wahl^(d)

University of Minnesota, Minneapolis, Minnesota 55455

D. S. Barton, G. Bunce, A. S. Carroll, and Y. I. Makdisi

Brookhaven National Laboratory, Upton, New York 11973

and

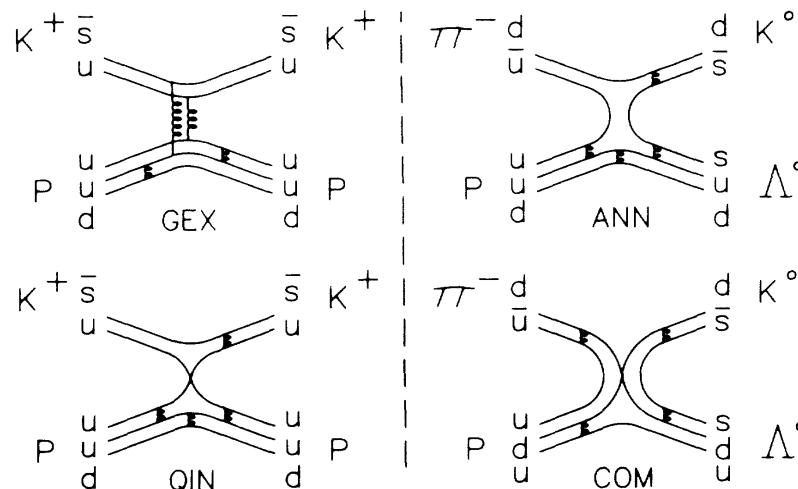
S. Gushue^(e) and J. J. Russell

Southeastern Massachusetts University, North Dartmouth, Massachusetts 02747

(Received 28 October 1987; revised manuscript received 3 February 1988)

Cross sections or upper limits are reported for twelve meson-baryon and two baryon-baryon reactions for an incident momentum of 9.9 GeV/c, near 90° c.m.: $\pi^\pm p \rightarrow p\pi^\pm, p\rho^\pm, \pi^+\Delta^\pm, K^+\Sigma^\pm, (\Lambda^0/\Sigma^0)K^0; K^\pm p \rightarrow pK^\pm; p^\pm p \rightarrow pp^\pm$. By studying the flavor dependence of the different reactions, we have been able to isolate the quark-interchange mechanism as dominant over gluon exchange and quark-antiquark annihilation.

- $\pi^\pm p \rightarrow p\pi^\pm,$
- $K^\pm p \rightarrow pK^\pm,$
- $\pi^\pm p \rightarrow p\rho^\pm,$
- $\pi^\pm p \rightarrow \pi^+\Delta^\pm,$
- $\pi^\pm p \rightarrow K^+\Sigma^\pm,$
- $\pi^- p \rightarrow \Lambda^0 K^0, \Sigma^0 K^0,$
- $p^\pm p \rightarrow pp^\pm.$



Features of Soft-Wall AdS/QCD

- Single-variable frame-independent radial Schrodinger equation
- Massless pion ($m_q = 0$)
- Regge Trajectories: universal slope in n and L
- Valid for all integer $J & S$. Spectrum is independent of S
- Dimensional Counting Rules for Hard Exclusive Processes
- Phenomenology: Space-like and Time-like Form Factors
- LF Holography: LFWFs; broad distribution amplitude
- No large N_c limit
- Add quark masses to LF kinetic energy
- Systematically improvable -- diagonalize H_{LF} on AdS basis

