

Result: Soft-Wall LFWF for massive constituents

$$\psi(x, \mathbf{k}_\perp) = \frac{4\pi c}{\kappa \sqrt{x(1-x)}} e^{-\frac{1}{2\kappa^2} \left(\frac{\mathbf{k}_\perp^2}{x(1-x)} + \frac{m_1^2}{x} + \frac{m_2^2}{1-x} \right)}$$

LF WF in impact space: soft-wall model with massive quarks

$$\psi(x, \mathbf{b}_\perp) = \frac{c\kappa}{\sqrt{\pi}} \sqrt{x(1-x)} e^{-\frac{1}{2}\kappa^2 x(1-x) \mathbf{b}_\perp^2 - \frac{1}{2\kappa^2} \left[\frac{m_1^2}{x} + \frac{m_2^2}{1-x} \right]}$$

$$z \rightarrow \zeta \rightarrow \chi$$

$$\chi^2 = b^2 x(1-x) + \frac{1}{\kappa^4} \left[\frac{m_1^2}{x} + \frac{m_2^2}{1-x} \right]$$

J/ψ

$\psi_{J/\psi}(x, b)$

LFWF peaks at

$$x_i = \frac{m_{\perp i}}{\sum_j^n m_{\perp j}}$$

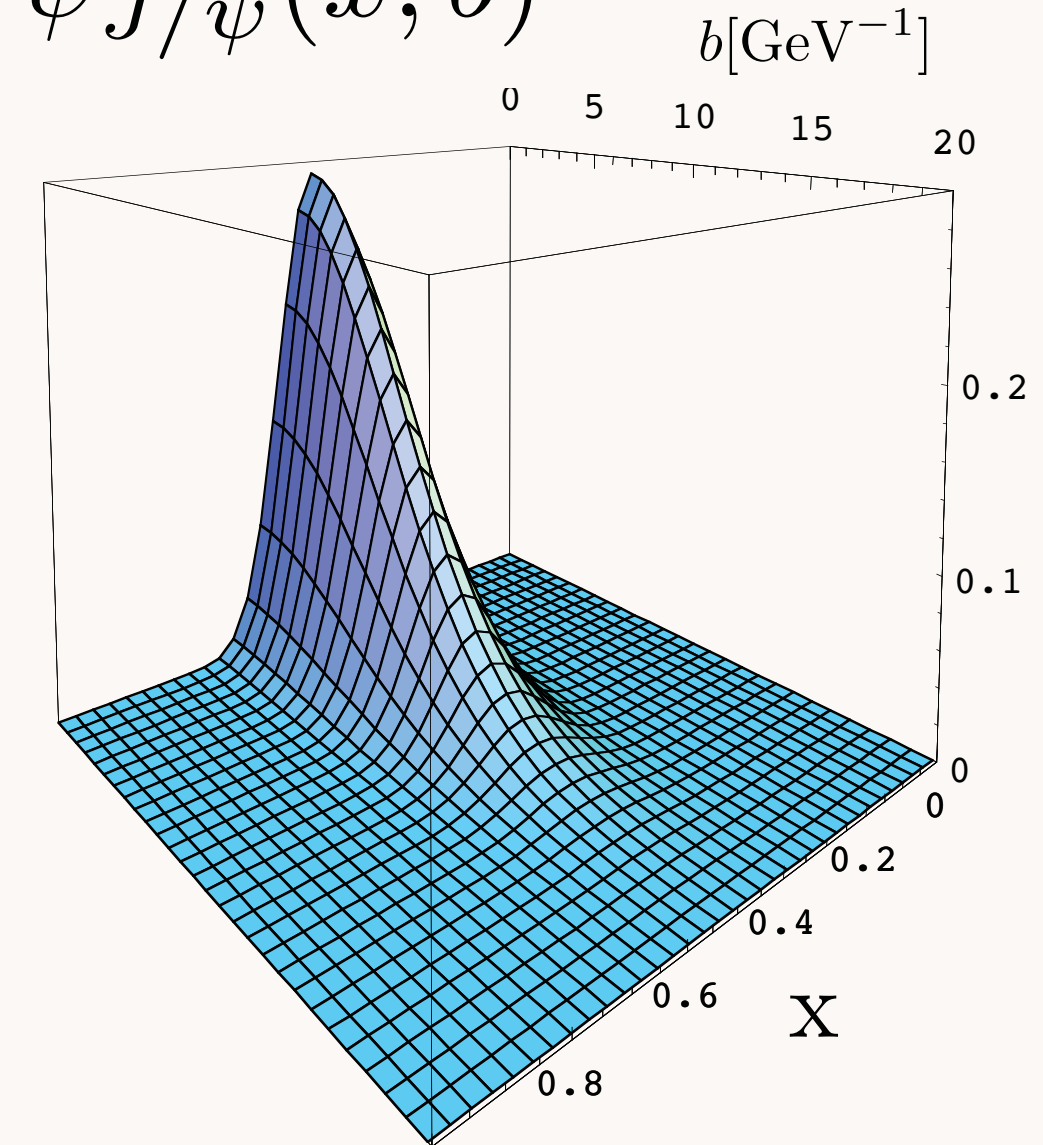
where

$$m_{\perp i} = \sqrt{m^2 + k_{\perp}^2}$$

*minimum of LF
energy
denominator*

$$\kappa = 0.375 \text{ GeV}$$

Landau Congress
Moscow, June 20, 2008



$$m_a = m_b = 1.25 \text{ GeV}$$

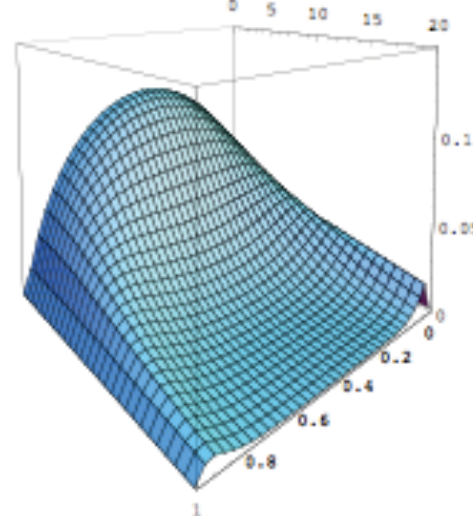
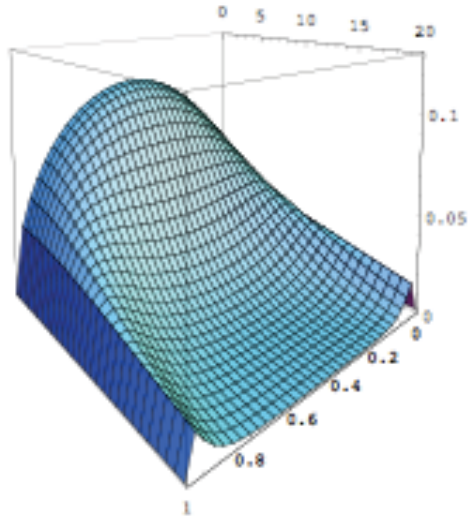
AdS/QCD
96

Stan Brodsky
SLAC & IPPP

$$|\pi^+\rangle = |u\bar{d}\rangle$$

$$m_u = 2 \text{ MeV}$$

$$m_d = 5 \text{ MeV}$$

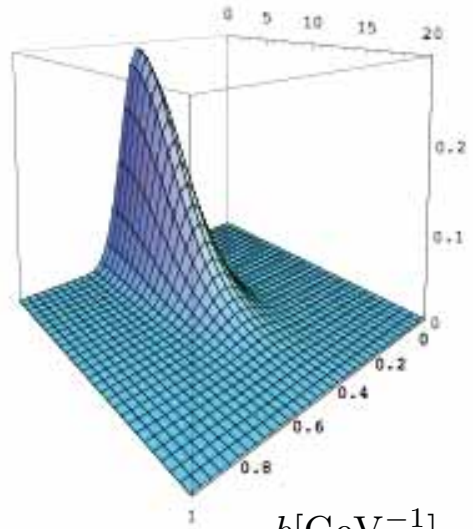
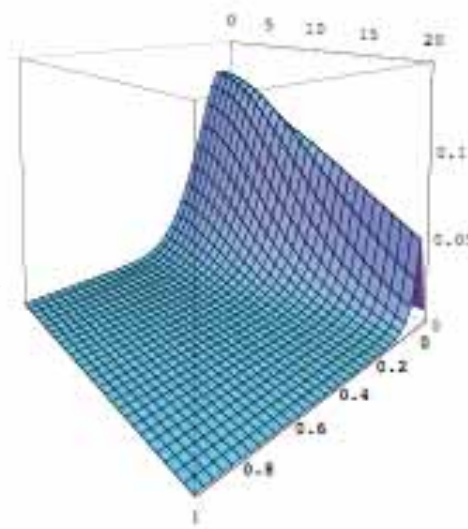


$$|K^+\rangle = |u\bar{s}\rangle$$

$$m_s = 95 \text{ MeV}$$

$$|D^+\rangle = |c\bar{d}\rangle$$

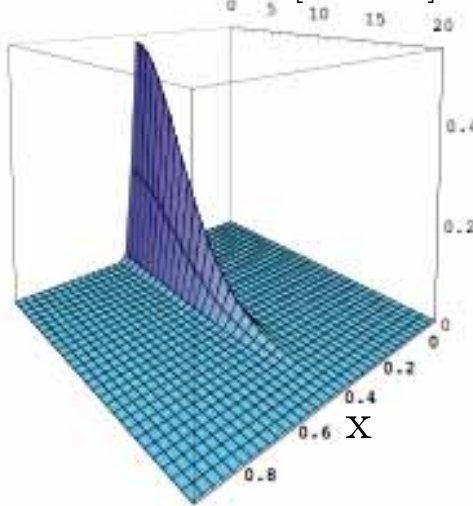
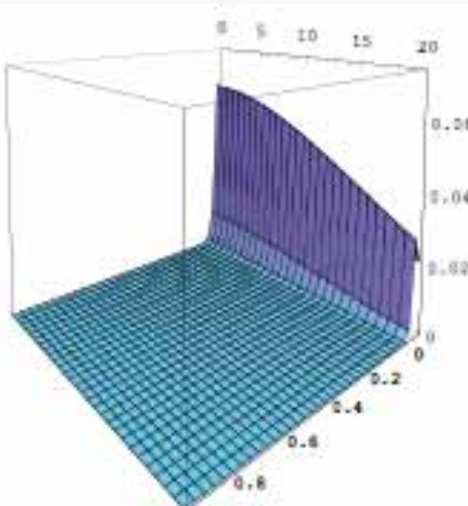
$$m_c = 1.25 \text{ GeV}$$



$$|\eta_c\rangle = |c\bar{c}\rangle$$

$$|B^+\rangle = |u\bar{b}\rangle$$

$$m_b = 4.2 \text{ GeV}$$

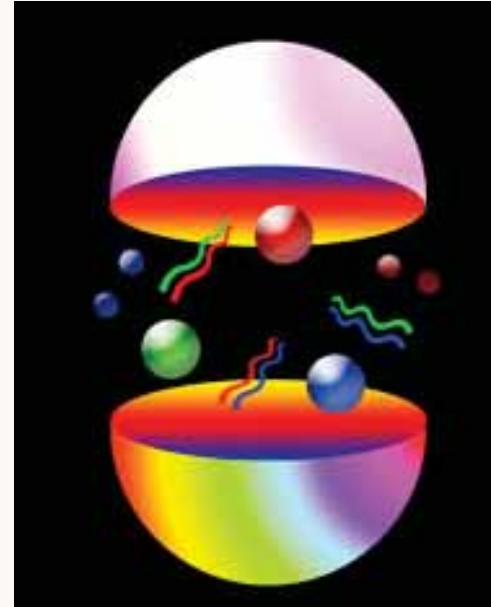


$$|\eta_b\rangle = |b\bar{b}\rangle$$

$$\kappa = 375 \text{ MeV}$$

- Baryons Spectrum in "bottom-up" holographic QCD
GdT and Brodsky: hep-th/0409074, hep-th/0501022.

Baryons in AdS/CFT



- Action for massive fermionic modes on AdS_{d+1} :

$$S[\bar{\Psi}, \Psi] = \int d^{d+1}x \sqrt{g} \bar{\Psi}(x, z) \left(i\Gamma^\ell D_\ell - \mu \right) \Psi(x, z).$$

- Equation of motion: $(i\Gamma^\ell D_\ell - \mu) \Psi(x, z) = 0$

$$\left[i \left(z\eta^{\ell m} \Gamma_\ell \partial_m + \frac{d}{2} \Gamma_z \right) + \mu R \right] \Psi(x^\ell) = 0.$$

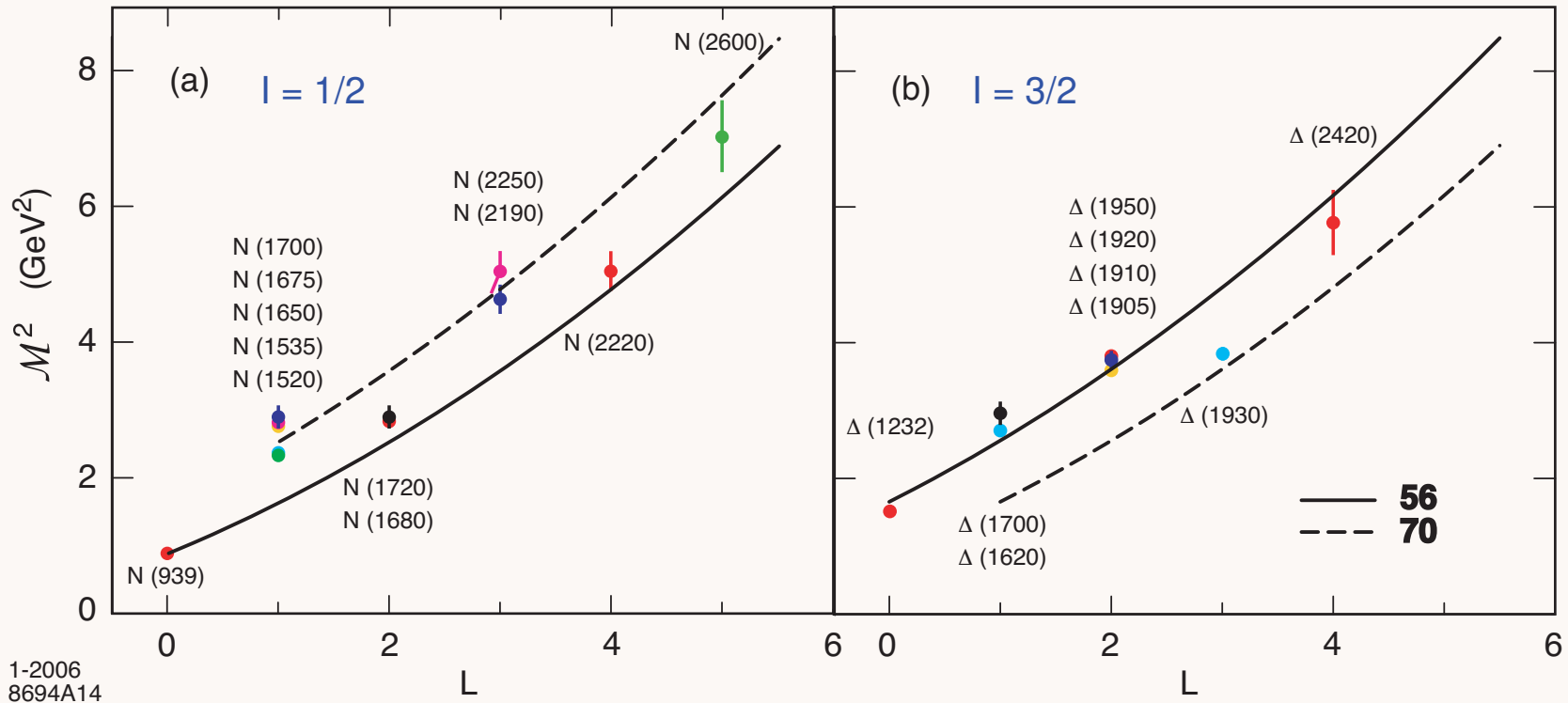


Fig: Light baryon orbital spectrum for $\Lambda_{QCD} = 0.25$ GeV in the HW model. The 56 trajectory corresponds to L even $P = +$ states, and the 70 to L odd $P = -$ states.

$SU(6)$	S	L	Baryon State
56	$\frac{1}{2}$	0	$N_{\frac{1}{2}}^{1+}$ (939)
	$\frac{3}{2}$	0	$\Delta_{\frac{3}{2}}^{3+}$ (1232)
70	$\frac{1}{2}$	1	$N_{\frac{1}{2}}^{1-}$ (1535) $N_{\frac{3}{2}}^{3-}$ (1520)
	$\frac{3}{2}$	1	$N_{\frac{1}{2}}^{1-}$ (1650) $N_{\frac{3}{2}}^{3-}$ (1700) $N_{\frac{5}{2}}^{5-}$ (1675)
	$\frac{1}{2}$	1	$\Delta_{\frac{1}{2}}^{1-}$ (1620) $\Delta_{\frac{3}{2}}^{3-}$ (1700)
56	$\frac{1}{2}$	2	$N_{\frac{3}{2}}^{3+}$ (1720) $N_{\frac{5}{2}}^{5+}$ (1680)
	$\frac{3}{2}$	2	$\Delta_{\frac{1}{2}}^{1+}$ (1910) $\Delta_{\frac{3}{2}}^{3+}$ (1920) $\Delta_{\frac{5}{2}}^{5+}$ (1905) $\Delta_{\frac{7}{2}}^{7+}$ (1950)
70	$\frac{1}{2}$	3	$N_{\frac{5}{2}}^{5-}$ $N_{\frac{7}{2}}^{7-}$
	$\frac{3}{2}$	3	$N_{\frac{3}{2}}^{3-}$ $N_{\frac{5}{2}}^{5-}$ $N_{\frac{7}{2}}^{7-}$ (2190) $N_{\frac{9}{2}}^{9-}$ (2250)
	$\frac{1}{2}$	3	$\Delta_{\frac{5}{2}}^{5-}$ (1930) $\Delta_{\frac{7}{2}}^{7-}$
56	$\frac{1}{2}$	4	$N_{\frac{7}{2}}^{7+}$ $N_{\frac{9}{2}}^{9+}$ (2220)
	$\frac{3}{2}$	4	$\Delta_{\frac{5}{2}}^{5+}$ $\Delta_{\frac{7}{2}}^{7+}$ $\Delta_{\frac{9}{2}}^{9+}$ $\Delta_{\frac{11}{2}}^{11+}$ (2420)
70	$\frac{1}{2}$	5	$N_{\frac{9}{2}}^{9-}$ $N_{\frac{11}{2}}^{11-}$ (2600)
	$\frac{3}{2}$	5	$N_{\frac{7}{2}}^{7-}$ $N_{\frac{9}{2}}^{9-}$ $N_{\frac{11}{2}}^{11-}$ $N_{\frac{13}{2}}^{13-}$

Space-Like Dirac Proton Form Factor

- Consider the spin non-flip form factors

$$F_+(Q^2) = g_+ \int d\zeta J(Q, \zeta) |\psi_+(\zeta)|^2,$$

$$F_-(Q^2) = g_- \int d\zeta J(Q, \zeta) |\psi_-(\zeta)|^2,$$

where the effective charges g_+ and g_- are determined from the spin-flavor structure of the theory.

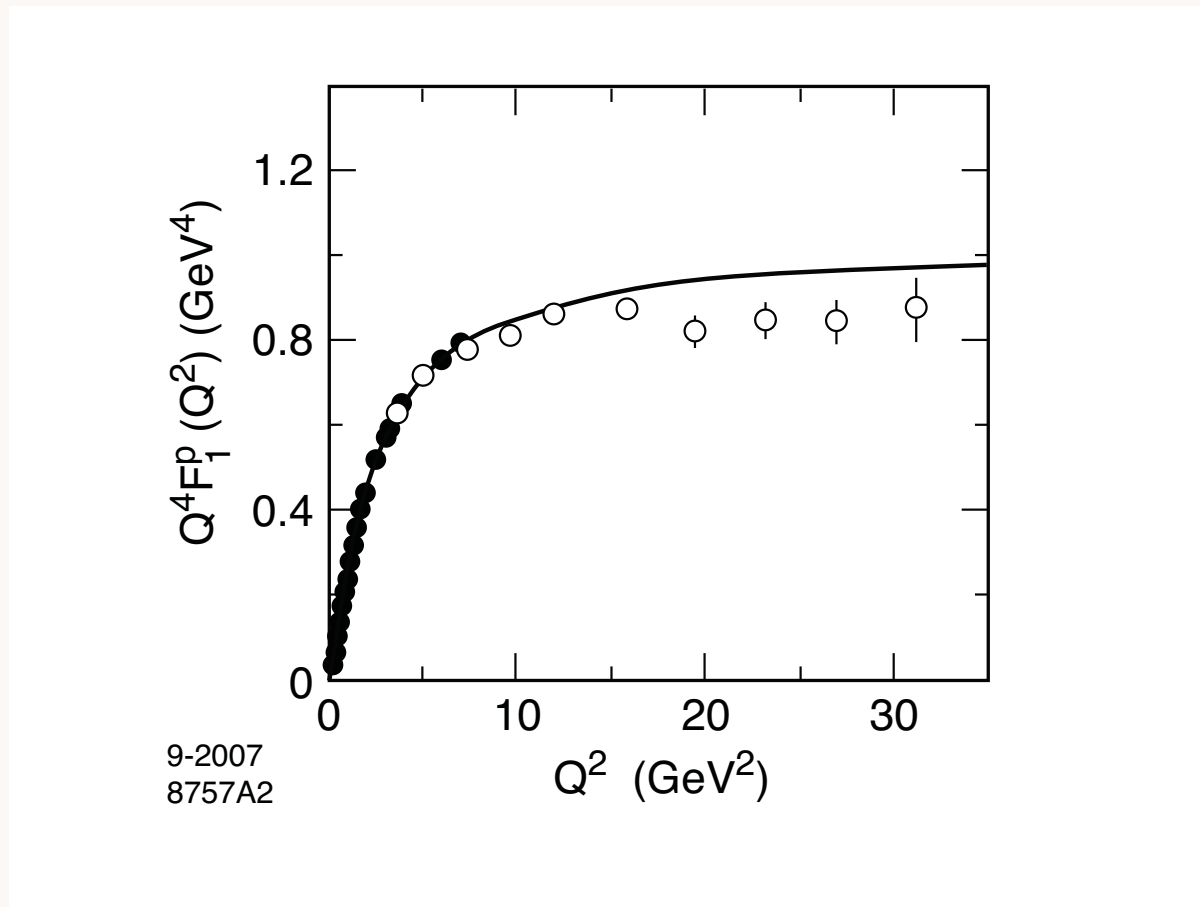
- Choose the struck quark to have $S^z = +1/2$. The two AdS solutions $\psi_+(\zeta)$ and $\psi_-(\zeta)$ correspond to nucleons with $J^z = +1/2$ and $-1/2$.
- For $SU(6)$ spin-flavor symmetry

$$F_1^p(Q^2) = \int d\zeta J(Q, \zeta) |\psi_+(\zeta)|^2,$$

$$F_1^n(Q^2) = -\frac{1}{3} \int d\zeta J(Q, \zeta) [|\psi_+(\zeta)|^2 - |\psi_-(\zeta)|^2],$$

where $F_1^p(0) = 1$, $F_1^n(0) = 0$.

- Scaling behavior for large Q^2 : $Q^4 F_1^p(Q^2) \rightarrow \text{constant}$ Proton $\tau = 3$

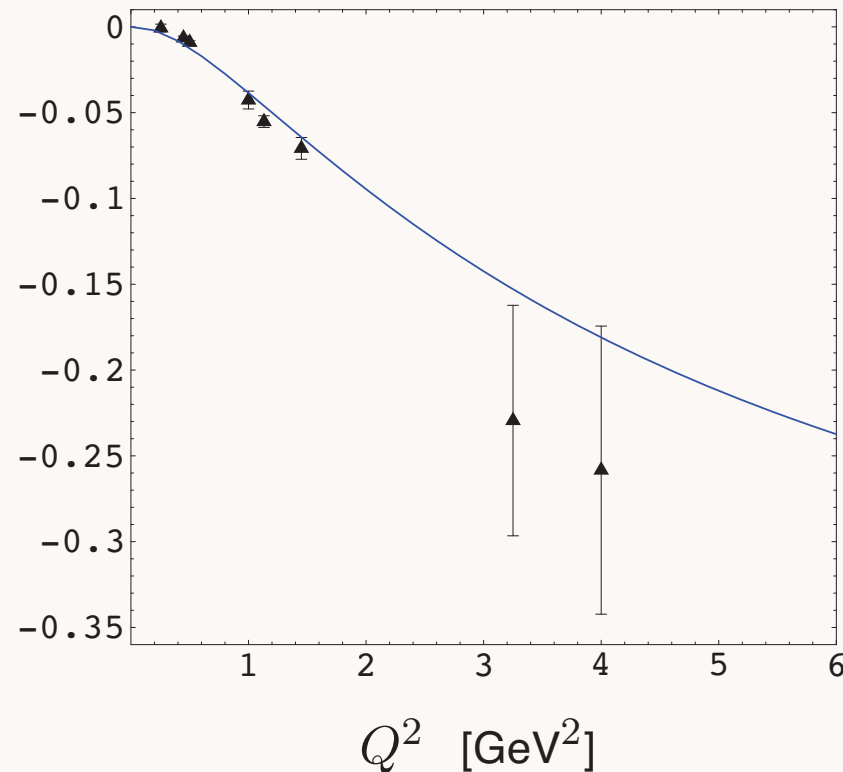


SW model predictions for $\kappa = 0.424$ GeV. Data analysis from: M. Diehl *et al.* Eur. Phys. J. C **39**, 1 (2005).

Dirac Neutron Form Factor (Valence Approximation)

Truncated Space Confinement

$$Q^4 F_1^n(Q^2) \text{ [GeV}^4\text{]}$$



Prediction for $Q^4 F_1^n(Q^2)$ for $\Lambda_{\text{QCD}} = 0.21$ GeV in the hard wall approximation. Data analysis from Diehl (2005).

Landau Congress
Moscow, June 20, 2008

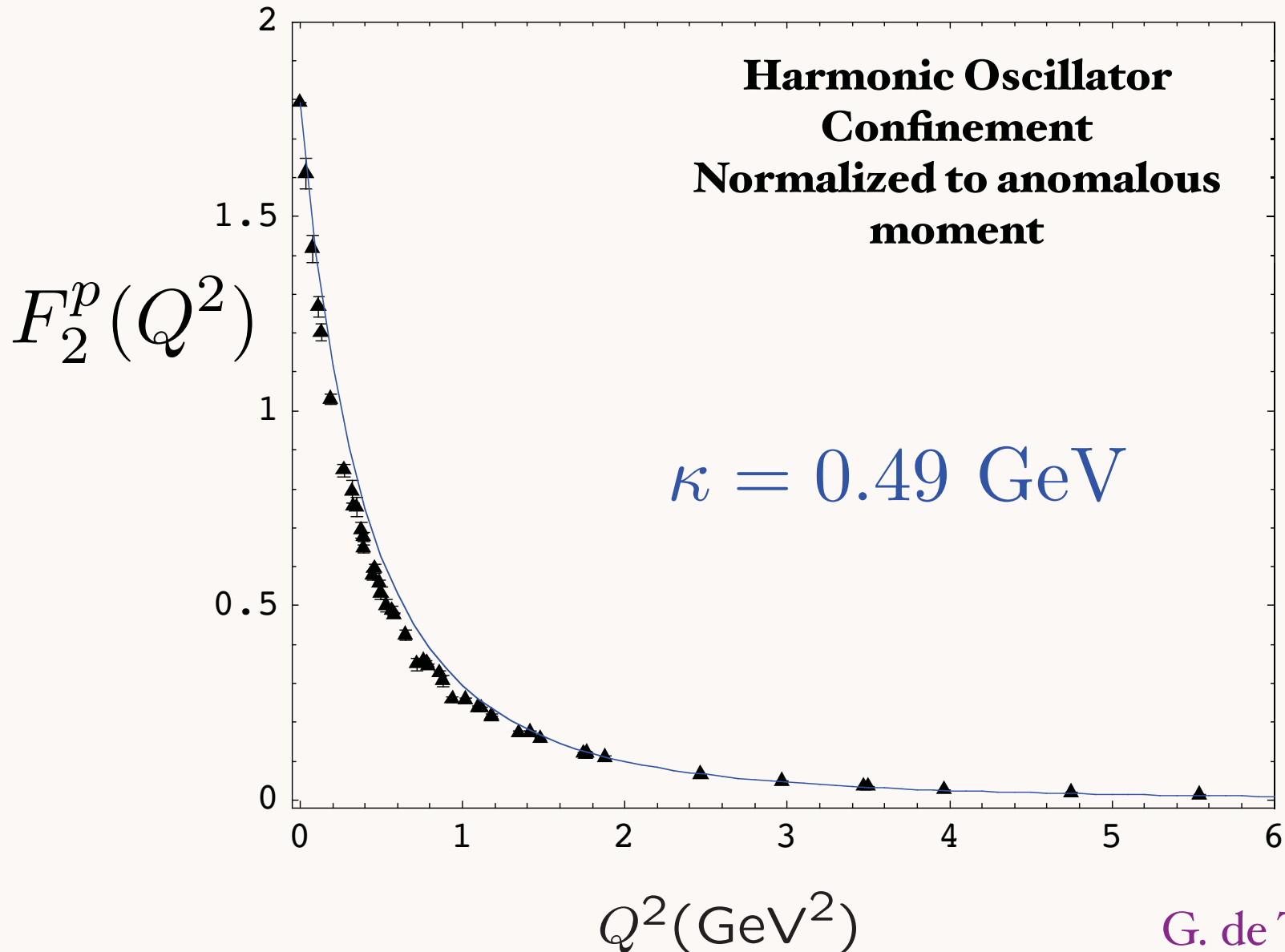
AdS/QCD
103

Stan Brodsky
SLAC & IPPP

Spacelike Pauli Form Factor

Preliminary

From overlap of $L = 1$ and $L = 0$ LFWFs



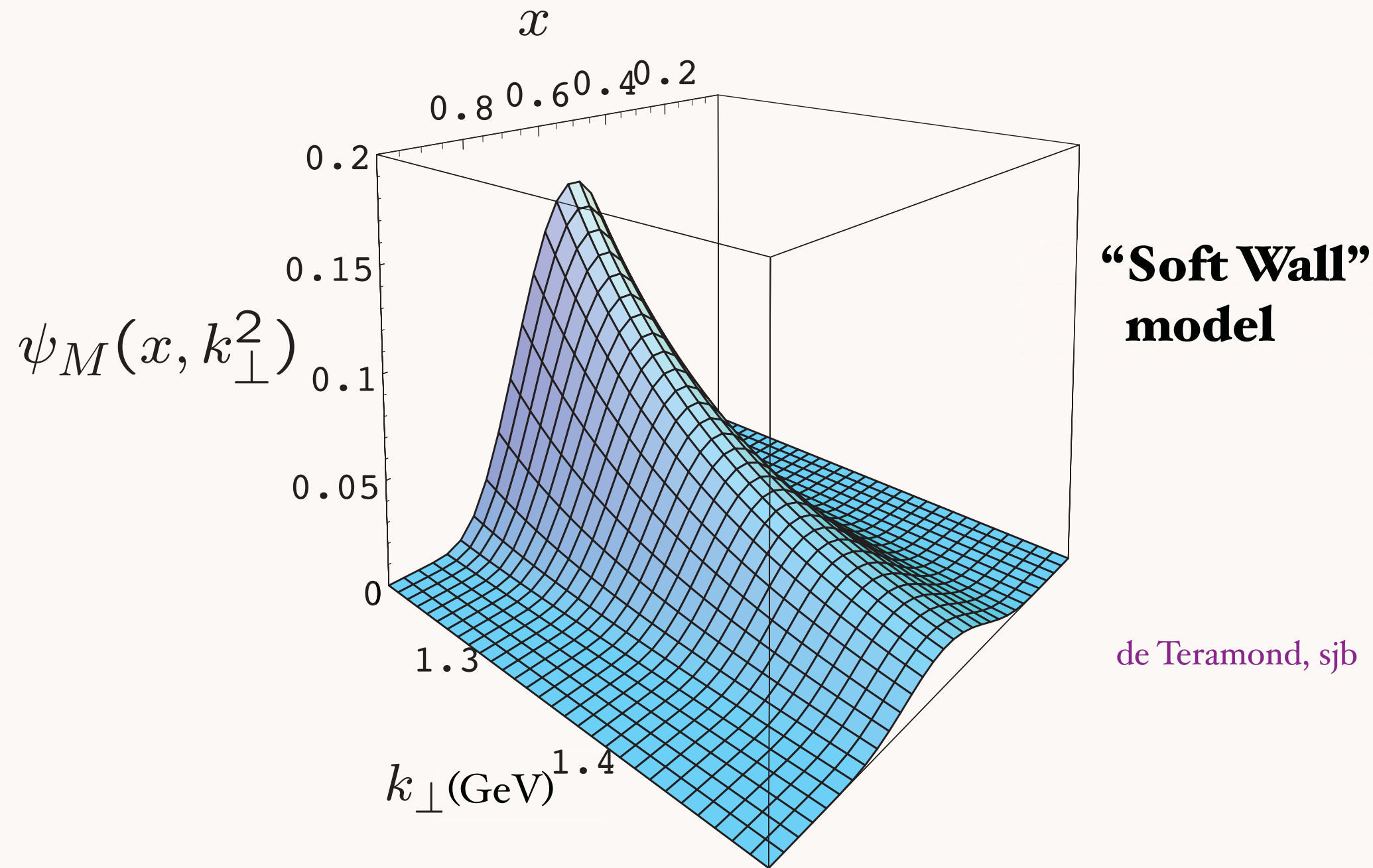
G. de Teramond, sjb

Landau Congress
Moscow, June 20, 2008

AdS/QCD
104

Stan Brodsky
SLAC & IPPP

Prediction from AdS/CFT: Meson LFWF



**“Soft Wall”
model**

de Teramond, sjb

Landau Congress
Moscow, June 20, 2008

AdS/QCD
105

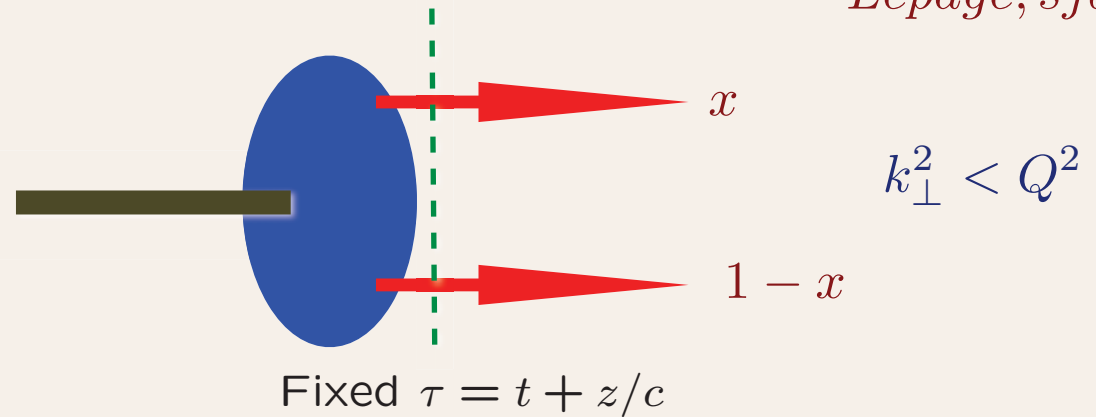
Stan Brodsky
SLAC & IPPP

Hadron Distribution Amplitudes

Lepage, sjb

$$\phi_H(x_i, Q)$$

$$\sum_i x_i = 1$$



- Fundamental gauge invariant non-perturbative input to hard exclusive processes, heavy hadron decays. Defined for Mesons, Baryons

- Evolution Equations from PQCD, OPE, Conformal Invariance

Lepage, sjb

Frishman, Lepage, Sachrajda, sjb

Peskin Braun

Efremov, Radyushkin Chernyak etal

- Compute from valence light-front wavefunction in light-cone gauge

$$\phi_M(x, Q) = \int^Q d^2 \vec{k} \psi_{q\bar{q}}(x, \vec{k}_{\perp})$$

Prediction from AdS/CFT: Meson LFWF

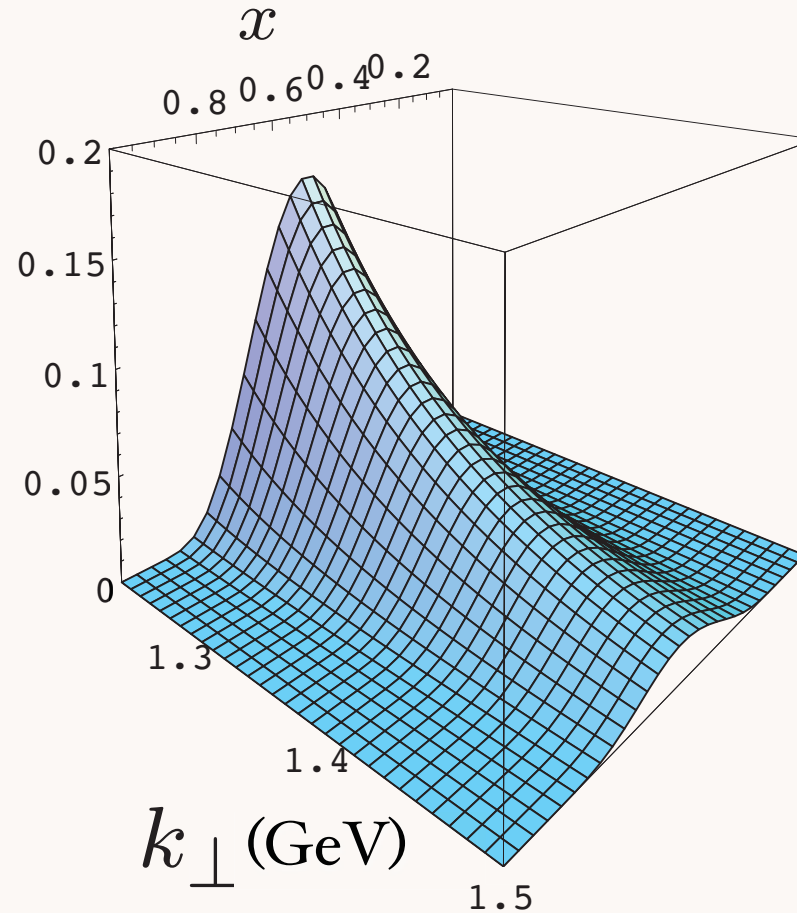
de Teramond, sjb

**“Soft Wall”
model**

$$\kappa = 0.375 \text{ GeV}$$

massless quarks

$$\psi_M(x, k_{\perp}^2)$$



$$\psi_M(x, k_{\perp}) = \frac{4\pi}{\kappa \sqrt{x(1-x)}} e^{-\frac{k_{\perp}^2}{2\kappa^2 x(1-x)}}$$

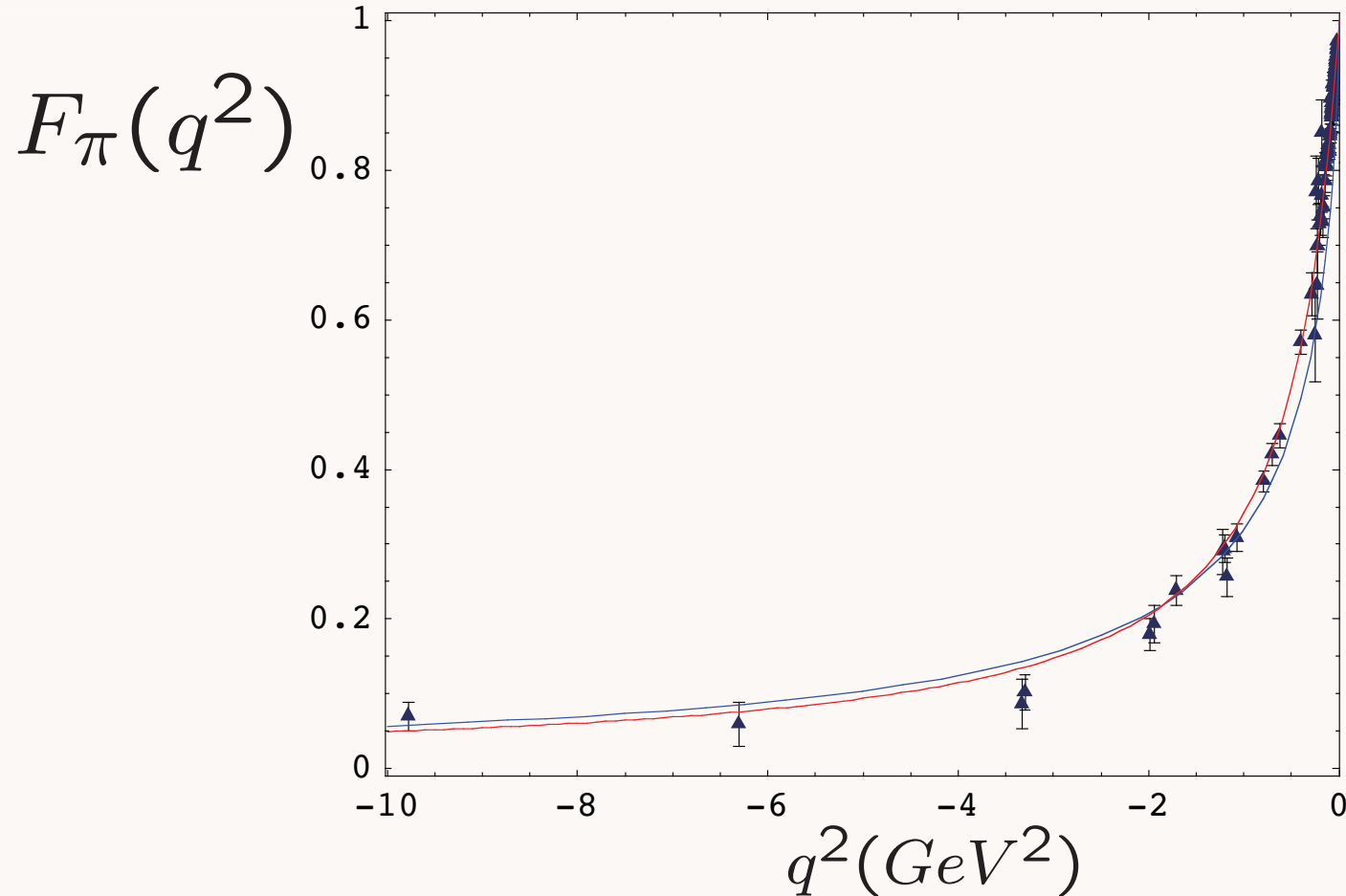
$$\phi_M(x, Q_0) \propto \sqrt{x(1-x)}$$

Landau Congress
Moscow, June 20, 2008

AdS/QCD
107

Stan Brodsky
SLAC & IPPP

Spacelike pion form factor from AdS/CFT



Data Compilation
Baldini, Kloe and Volmer

— Soft Wall: Harmonic Oscillator Confinement

— Hard Wall: Truncated Space Confinement

One parameter - set by pion decay constant

Landau Congress
Moscow, June 20, 2008

AdS/QCD
108

de Teramond, sjb
See also: Radyushkin
Stan Brodsky
SLAC & IPPP

Light-Front Wavefunctions

Dirac's Front Form: Fixed $\tau = t + z/c$

$$\Psi(x, k_{\perp}) \quad x_i = \frac{k_i^+}{P^+}$$

Invariant under boosts. Independent of P^{μ}

$$H_{LF}^{QCD} |\psi\rangle = M^2 |\psi\rangle$$

*Remarkable new insights from AdS/CFT,
the duality between conformal field theory
and Anti-de Sitter Space*

How can we systematically improve AdS/QCD?

AdS/QCD: Semiclassical model

No Particle Creation

Valence Fock State only

$$|p, S_z\rangle = \sum_{n=3} \Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i) |n; \vec{k}_{\perp i}, \lambda_i\rangle$$

sum over states with $n=3, 4, \dots$ constituents

The Light Front Fock State Wavefunctions

$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

are boost invariant; they are independent of the hadron's energy and momentum P^μ .

The light-cone momentum fraction

$$x_i = \frac{k_i^+}{p^+} = \frac{k_i^0 + k_i^z}{P^0 + P^z}$$

are boost invariant.

$$\sum_i^n k_i^+ = P^+, \quad \sum_i^n x_i = 1, \quad \sum_i^n \vec{k}_i^\perp = \vec{0}^\perp.$$

Intrinsic heavy quarks

$$\bar{u}(x) \neq \bar{d}(x)$$

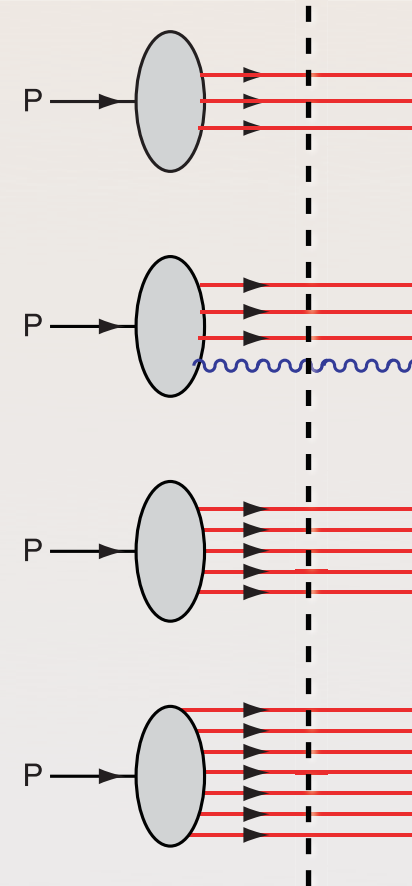
Mueller: BFKL DYNAMICS

$$\bar{s}(x) \neq s(x)$$

Landau Congress
Moscow, June 20, 2008

AdS/QCD
III

Stan Brodsky
SLAC & IPPP



Fixed LF time

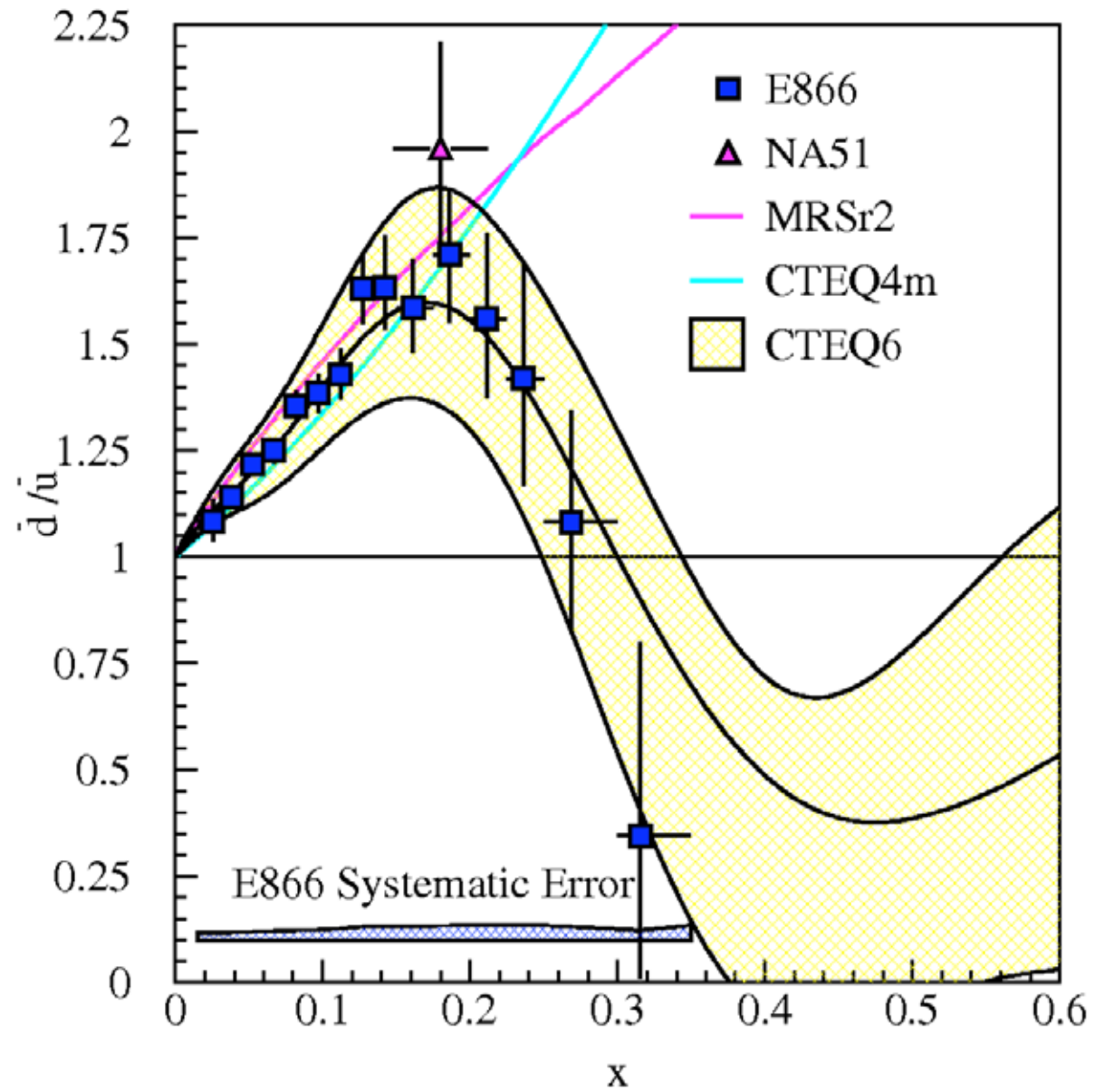
Light Antiquark Flavor Asymmetry

- Naïve Assumption from gluon splitting:

$$\bar{d}(x) = \bar{u}(x)$$

- E866/NuSea (Drell-Yan)

$\bar{d}(x)/\bar{u}(x)$ for $0.015 \leq x \leq 0.35$



Light-Front QCD

Heisenberg Matrix Formulation

Physical gauge: $A^+ = 0$

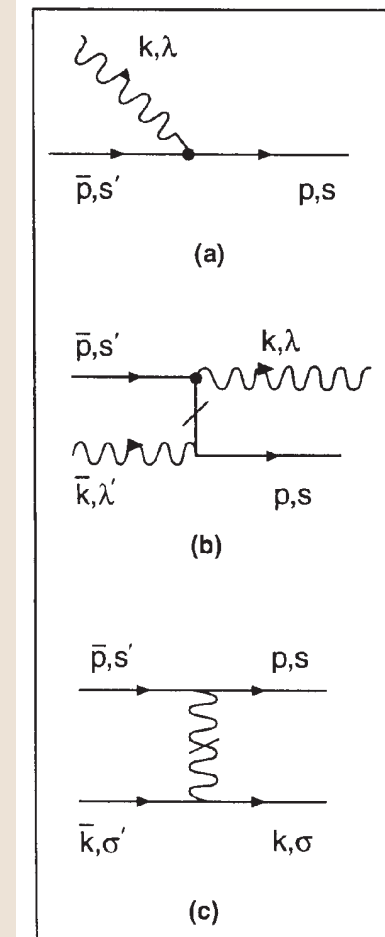
$$L^{QCD} \rightarrow H_{LF}^{QCD}$$

$$H_{LF}^{QCD} = \sum_i \left[\frac{m^2 + k_{\perp}^2}{x} \right]_i + H_{LF}^{int}$$

H_{LF}^{int} : Matrix in Fock Space

$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

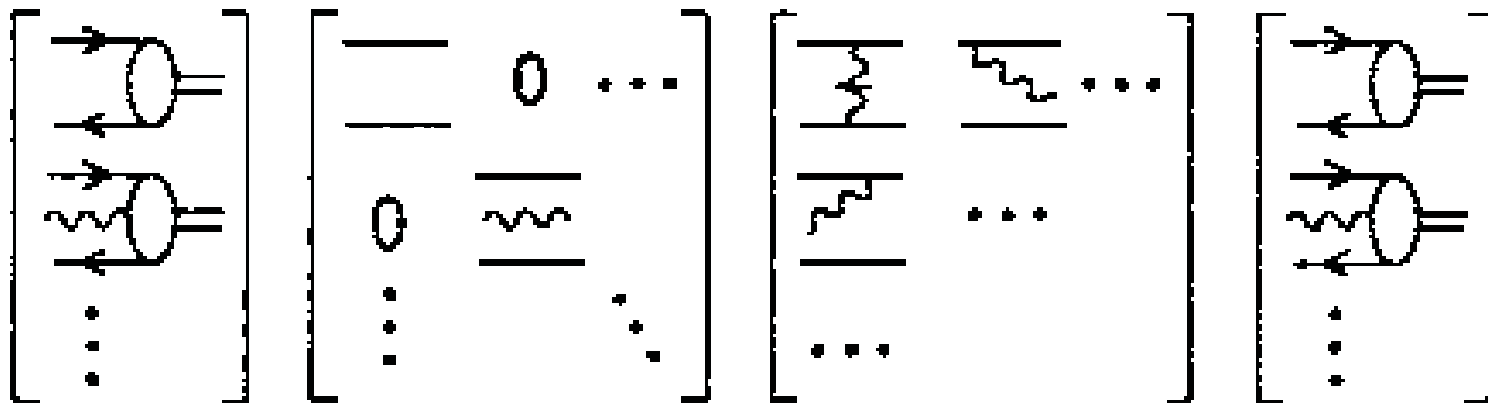
Eigenvalues and Eigensolutions give Hadron Spectrum and Light-Front wavefunctions



DLCQ: Periodic BC in x^- . Discrete k^+ ; frame-independent truncation

LIGHT-FRONT SCHRÖDINGER EQUATION

$$\left(M_\pi^2 - \sum_i \frac{\vec{k}_{\perp i}^2 + m_i^2}{x_i} \right) \begin{bmatrix} \psi_{q\bar{q}/\pi} \\ \psi_{q\bar{q}g/\pi} \\ \vdots \end{bmatrix} = \begin{bmatrix} \langle q\bar{q} | V | q\bar{q} \rangle & \langle q\bar{q} | V | q\bar{q}g \rangle & \cdots \\ \langle q\bar{q}g | V | q\bar{q} \rangle & \langle q\bar{q}g | V | q\bar{q}g \rangle & \cdots \\ \vdots & \vdots & \ddots \end{bmatrix} \begin{bmatrix} \psi_{q\bar{q}/\pi} \\ \psi_{q\bar{q}g/\pi} \\ \vdots \end{bmatrix}$$



$$A^+ = 0$$

G.P. Lepage, sjb

Landau Congress
Moscow, June 20, 2008

AdS/QCD
114

Stan Brodsky
SLAC & IPPP

Light-Front QCD

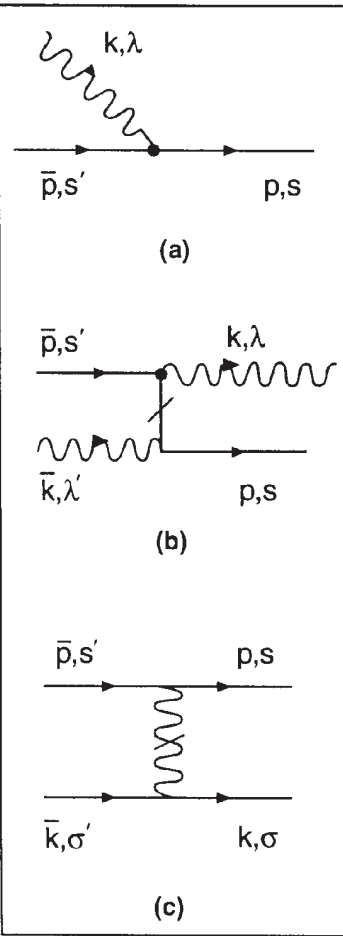
Heisenberg Matrix Formulation

$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

DLCQ

Discretized Light-Cone Quantization

n	Sector	1 q \bar{q}	2 gg	3 q \bar{q} g	4 q \bar{q} q \bar{q}	5 ggg	6 q \bar{q} gg	7 q \bar{q} q \bar{q} g	8 q \bar{q} q \bar{q} q \bar{q}	9 gggg	10 q \bar{q} ggg	11 q \bar{q} q \bar{q} gg	12 q \bar{q} q \bar{q} q \bar{q} g	13 q \bar{q} q \bar{q} q \bar{q} q \bar{q}
1	q \bar{q}				
2	gg			
3	q \bar{q} g							
4	q \bar{q} q \bar{q}	
5	ggg
6	q \bar{q} gg						
7	q \bar{q} q \bar{q} g
8	q \bar{q} q \bar{q} q \bar{q}			
9	gggg
10	q \bar{q} ggg
11	q \bar{q} q \bar{q} gg
12	q \bar{q} q \bar{q} q \bar{q} g			
13	q \bar{q} q \bar{q} q \bar{q} q \bar{q}		



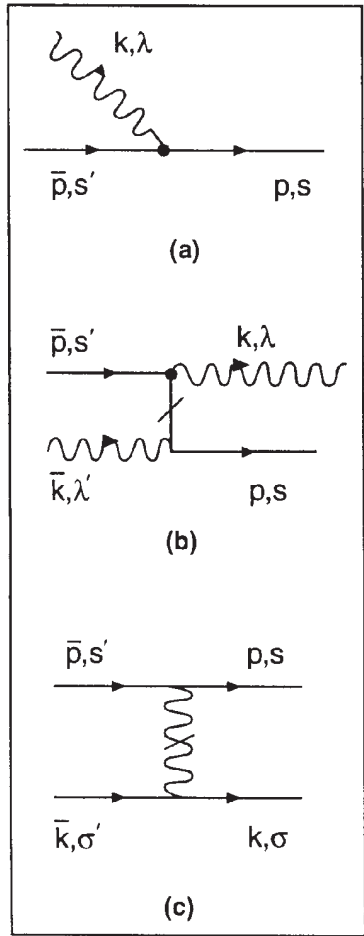
Eigenvalues and Eigensolutions give Hadron Spectrum and Light-Front wavefunctions

H.C. Pauli & sjb

DLCQ: Frame-independent, No fermion doubling; Minkowski Space

Light-Front QCD Heisenberg Equation

$$H_{LC}^{QCD} |\Psi_h\rangle = M_h^2 |\Psi_h\rangle$$



n	Sector	1 q \bar{q}	2 gg	3 q \bar{q} g	4 q \bar{q} q \bar{q}	5 gg g	6 q \bar{q} gg	7 q \bar{q} q \bar{q} g	8 q \bar{q} q \bar{q} q \bar{q}	9 gg gg	10 q \bar{q} gg g	11 q \bar{q} q \bar{q} gg	12 q \bar{q} q \bar{q} q \bar{q} g	13 q \bar{q} q \bar{q} q \bar{q} q \bar{q}
1	q \bar{q}				
2	gg			
3	q \bar{q} g							
4	q \bar{q} q \bar{q}	
5	gg g
6	q \bar{q} gg								.				.	.
7	q \bar{q} q \bar{q} g
8	q \bar{q} q \bar{q} q \bar{q}			
9	gg gg
10	q \bar{q} gg g
11	q \bar{q} q \bar{q} gg
12	q \bar{q} q \bar{q} q \bar{q} g				
13	q \bar{q} q \bar{q} q \bar{q} q \bar{q}			

Use AdS/QCD basis functions

Landau Congress
Moscow, June 20, 2008

AdS/QCD
116

Stan Brodsky
SLAC & IPPP

*Use AdS/CFT orthonormal LFWFs
as a basis for diagonalizing
the QCD LF Hamiltonian*

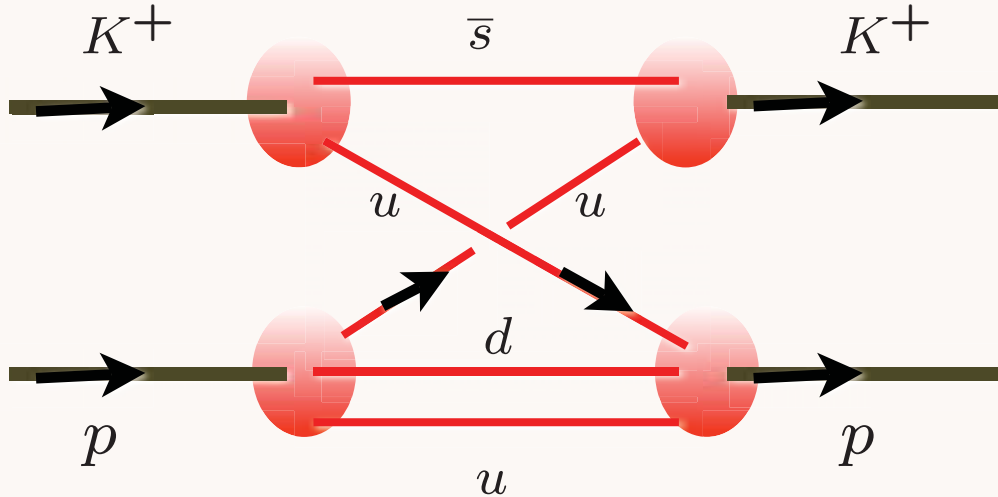
- Good initial approximant: generates all Fock states
- Better than plane wave basis Pauli, Hornbostel, Hiller,
McCartor, sjb
- DLCQ discretization -- highly successful I+I
- Use independent HO LFWFs, remove CM motion Vary, Harinandrath, Maris, sjb
- Similar to Shell Model calculations

Holographic Connection between LF and AdS/CFT

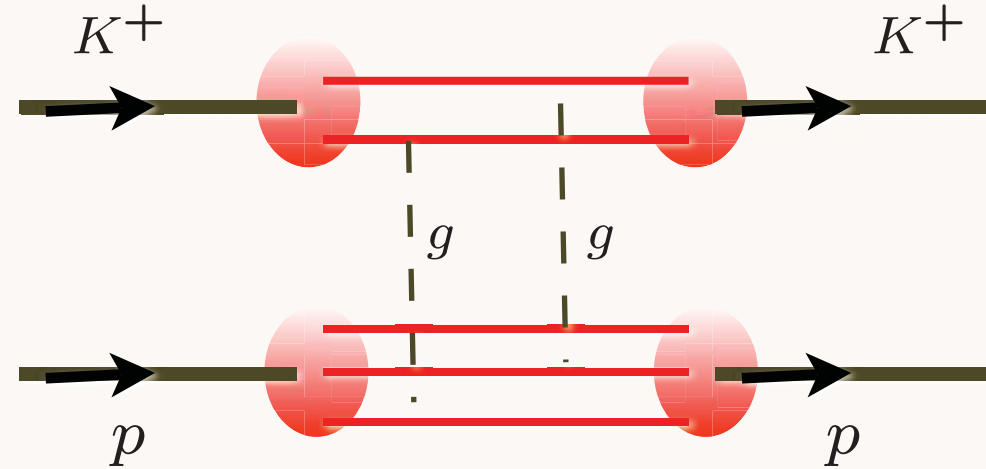
- Predictions for hadronic spectra, light-front wavefunctions, interactions
- Deduce meson and baryon wavefunctions, distribution amplitude, structure function from holographic constraint
- Identification of Orbital Angular Momentum Casimir for $SO(2)$: LF Rotations
- Extension to massive quarks

New Perspectives for QCD from AdS/CFT

- LFWFs: Fundamental frame-independent description of hadrons at amplitude level
- Holographic Model from AdS/CFT : Confinement at large distances and conformal behavior at short distances
- Model for LFWFs, meson and baryon spectra: many applications!
- New basis for diagonalizing Light-Front Hamiltonian
- Physics similar to MIT bag model, but covariant. No problem with support $0 < x < 1$.
- Quark Interchange dominant force at short distances



Quark Interchange
(Spin exchange in atom-atom scattering)



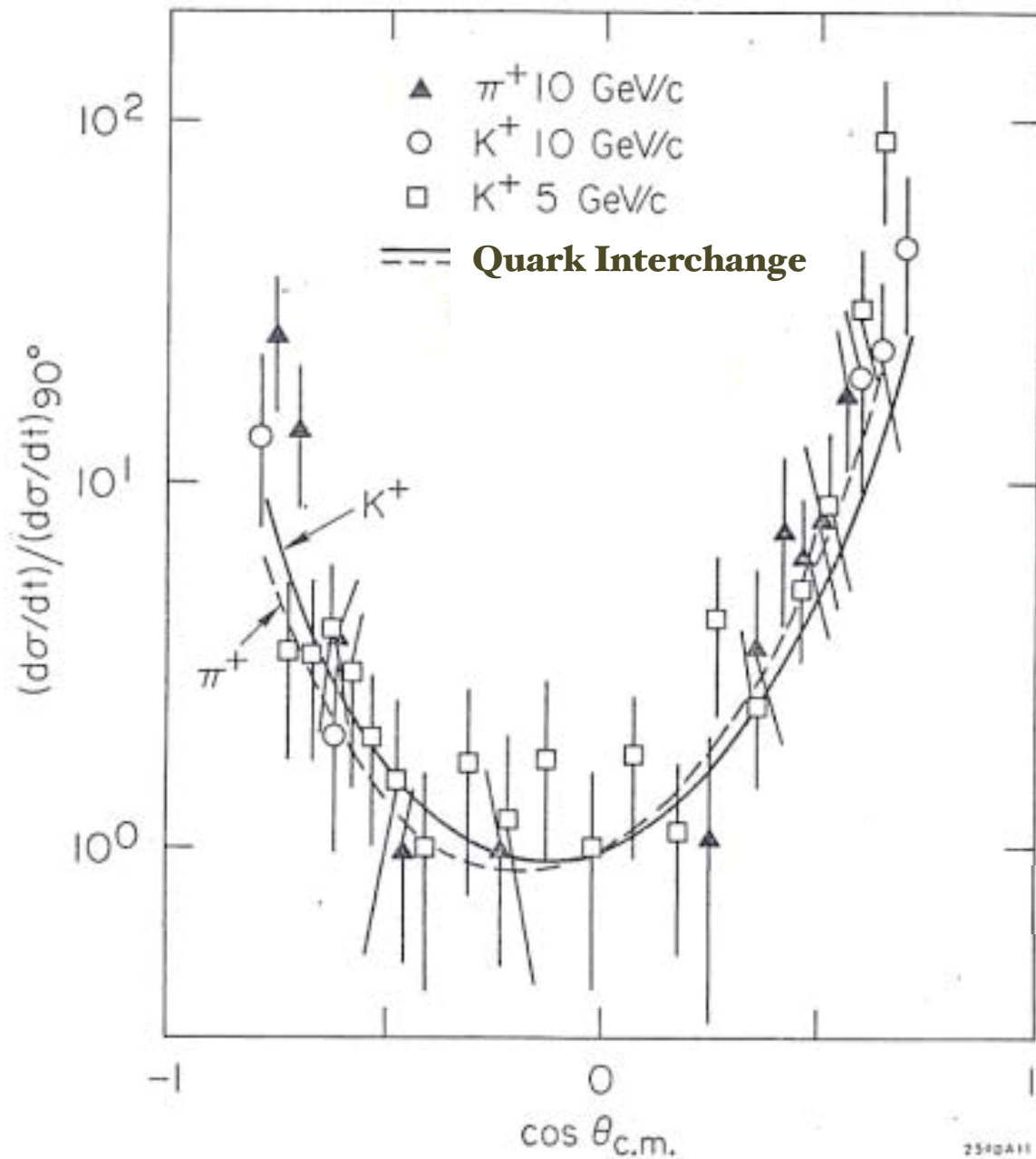
Gluon Exchange
(Van der Waal -- Landshoff)

$$\frac{d\sigma}{dt} = \frac{|M(s,t)|^2}{s^2}$$

$$M(t, u)_{\text{interchange}} \propto \frac{1}{ut^2}$$

$$M(s, t)_{\text{gluonexchange}} \propto sF(t)$$

MIT Bag Model (de Tar), large \$N_c\$, ('t Hooft), AdS/CFT
all predict dominance of quark interchange:



AdS/CFT explains why quark interchange is dominant interaction at high momentum transfer in exclusive reactions

$$M(t, u)_{\text{interchange}} \propto \frac{1}{ut^2}$$

Non-linear Regge behavior:

$$\alpha_R(t) \rightarrow -1$$

Comparison of Exclusive Reactions at Large t

B. R. Baller,^(a) G. C. Blazey,^(b) H. Courant, K. J. Heller, S. Heppelmann,^(c) M. L. Marshak,
E. A. Peterson, M. A. Shupe, and D. S. Wahl^(d)
University of Minnesota, Minneapolis, Minnesota 55455

D. S. Barton, G. Bunce, A. S. Carroll, and Y. I. Makdisi
Brookhaven National Laboratory, Upton, New York 11973

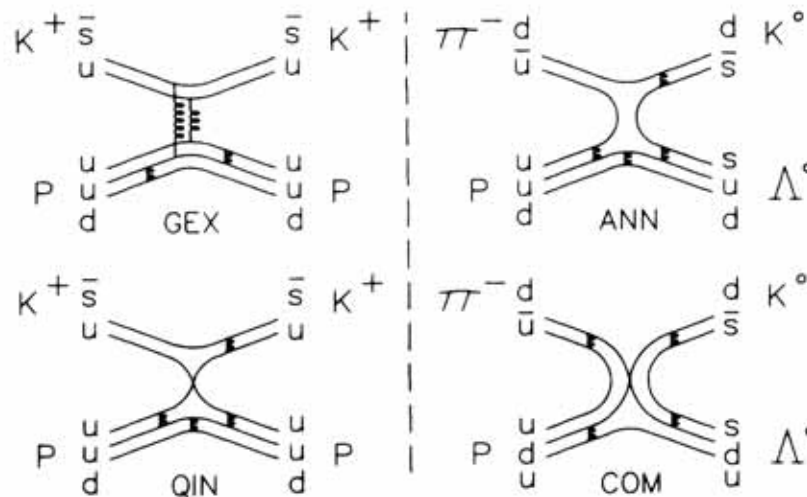
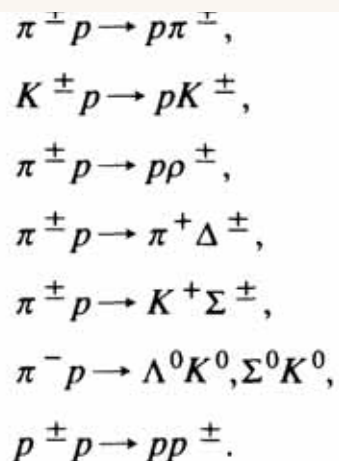
and

S. Gushue^(e) and J. J. Russell

Southeastern Massachusetts University, North Dartmouth, Massachusetts 02747

(Received 28 October 1987; revised manuscript received 3 February 1988)

Cross sections or upper limits are reported for twelve meson-baryon and two baryon-baryon reactions for an incident momentum of 9.9 GeV/c, near 90° c.m.: $\pi^\pm p \rightarrow p\pi^\pm, p\rho^\pm, \pi^+\Delta^\pm, K^+\Sigma^\pm, (\Lambda^0/\Sigma^0)K^0$; $K^\pm p \rightarrow pK^\pm$; $p^\pm p \rightarrow pp^\pm$. By studying the flavor dependence of the different reactions, we have been able to isolate the quark-interchange mechanism as dominant over gluon exchange and quark-antiquark annihilation.



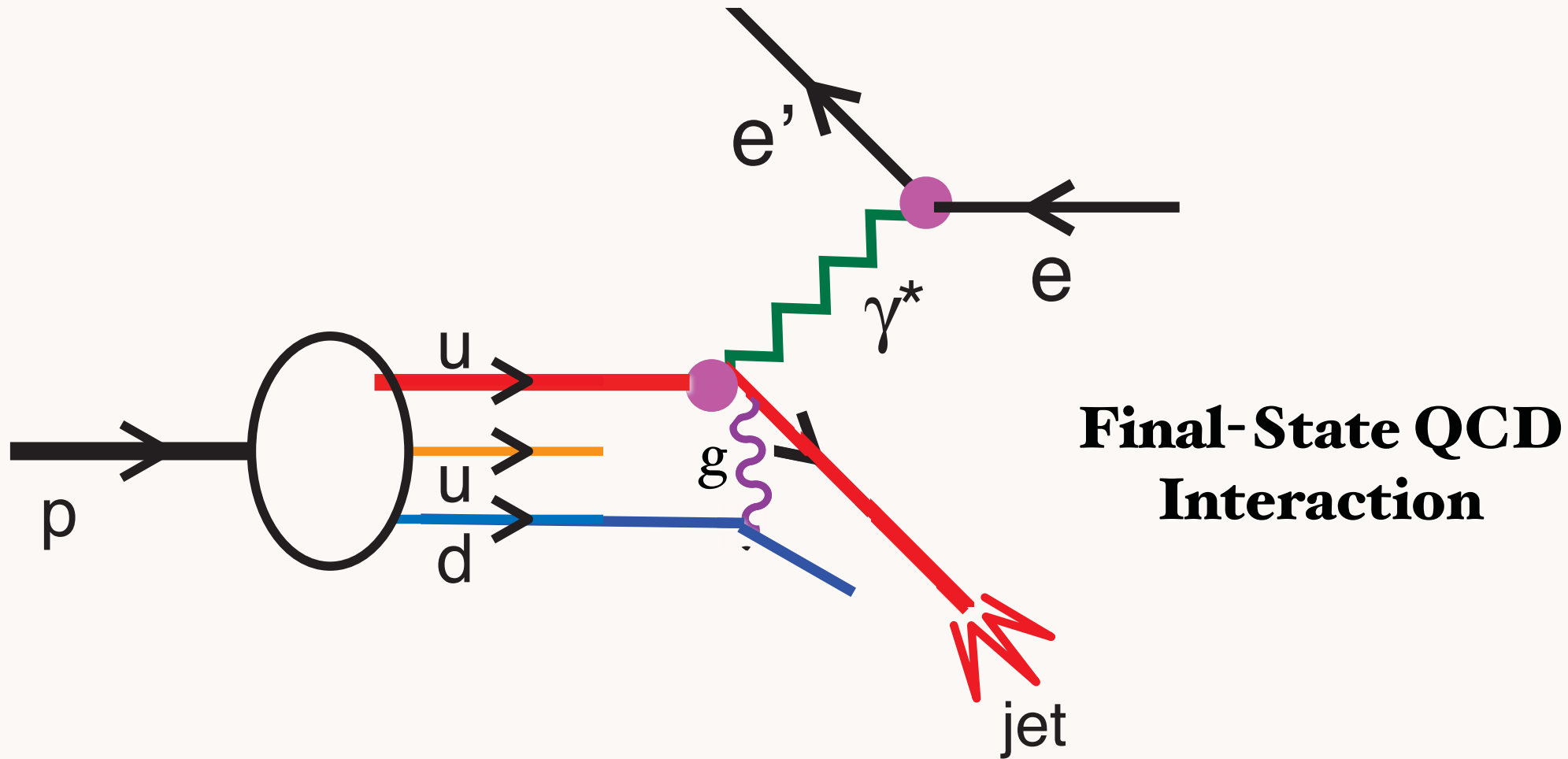
New Perspectives on QCD Phenomena from AdS/CFT

- **AdS/CFT:** Duality between string theory in Anti-de Sitter Space and Conformal Field Theory
- New Way to Implement Conformal Symmetry
- Holographic Model: Conformal Symmetry at Short Distances, Confinement at large distances
- Remarkable predictions for hadronic spectra, wavefunctions, interactions
- AdS/CFT provides novel insights into the quark structure of hadrons

Hadron Dynamics at the Amplitude Level

- LFWFS are the universal hadronic amplitudes which underlie structure functions, GPDs, exclusive processes, distribution amplitudes, direct subprocesses, hadronization.
- Relation of spin, momentum, and other distributions to physics of the hadron itself.
- Connections between observables, orbital angular momentum
- Role of FSI and ISIs: Diffractive DIS, Sivers effect

Deep Inelastic Electron-Proton Scattering



*Conventional wisdom:
Final-state interactions of struck quark can be neglected*

Single-spin asymmetries

Leading Twist Sivers Effect

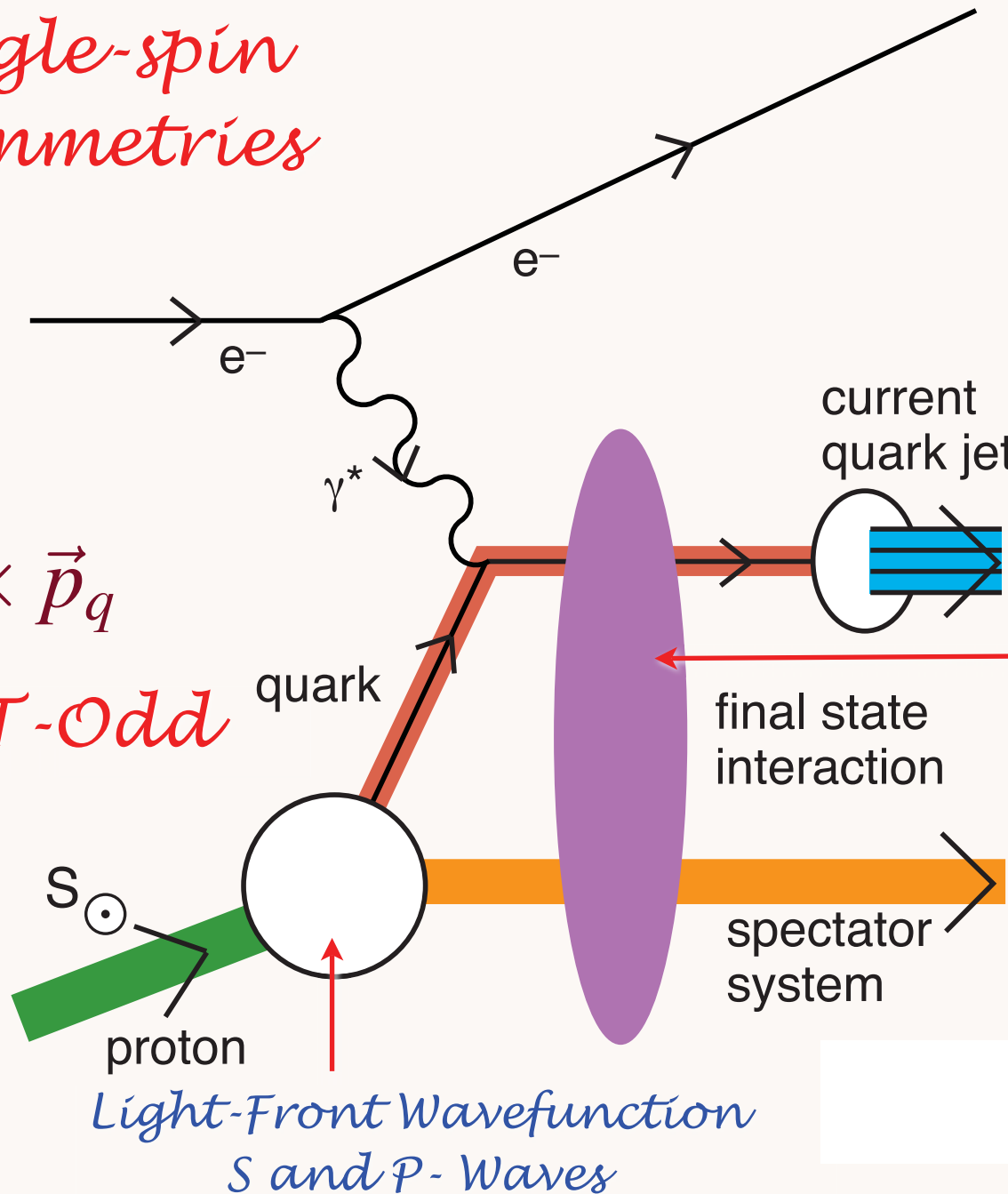
Hwang,
Schmidt, sjb

Collins, Burkardt
Ji, Yuan

*QCD S- and P-
Coulomb Phases
--Wilson Line*

$$i \vec{S}_p \cdot \vec{q} \times \vec{p}_q$$

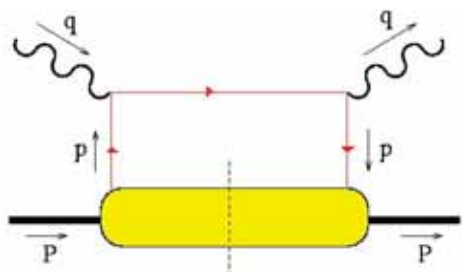
Pseudo-T-Odd



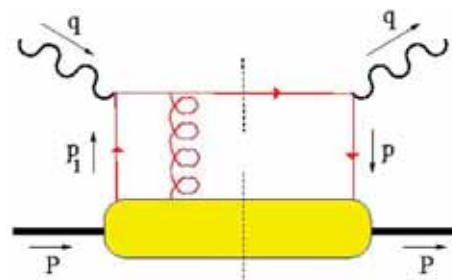
Landau Congress
Moscow, June 20, 2008

AdS/QCD
126

Stan Brodsky
SLAC & IPPP



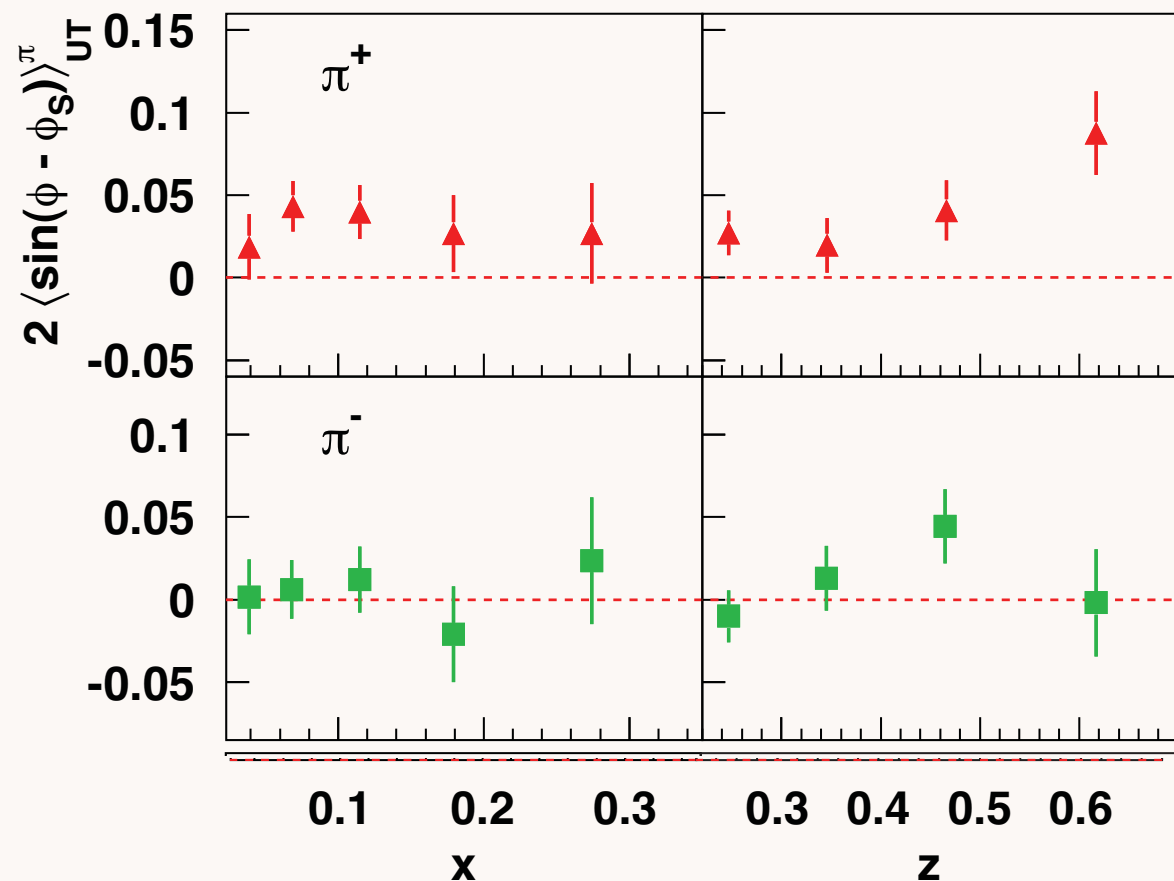
can interfere with



and produce a T-odd effect!
(also need $L_z \neq 0$)

HERMES coll., A. Airapetian et al., Phys. Rev. Lett. 94 (2005) 012002.

Sivers asymmetry from HERMES



- First evidence for non-zero Sivers function!
- \Rightarrow presence of non-zero **quark orbital angular momentum!**
- **Positive** for π^+ ...
Consistent with zero for π^- ...

Gamberg: Hermes data compatible with BHS model

Schmidt, Lu: Hermes charge pattern follow quark contributions to anomalous moment

**Stan Brodsky
SLAC & IPPP**

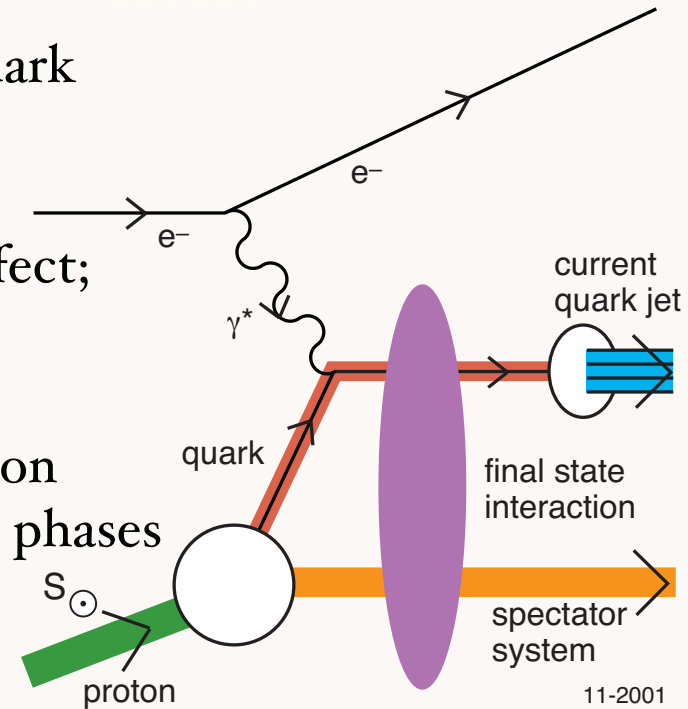
**Landau Congress
Moscow, June 20, 2008**

**AdS/QCD
127**

Final-State Interactions Produce Pseudo-T-Odd (Sivers Effect)

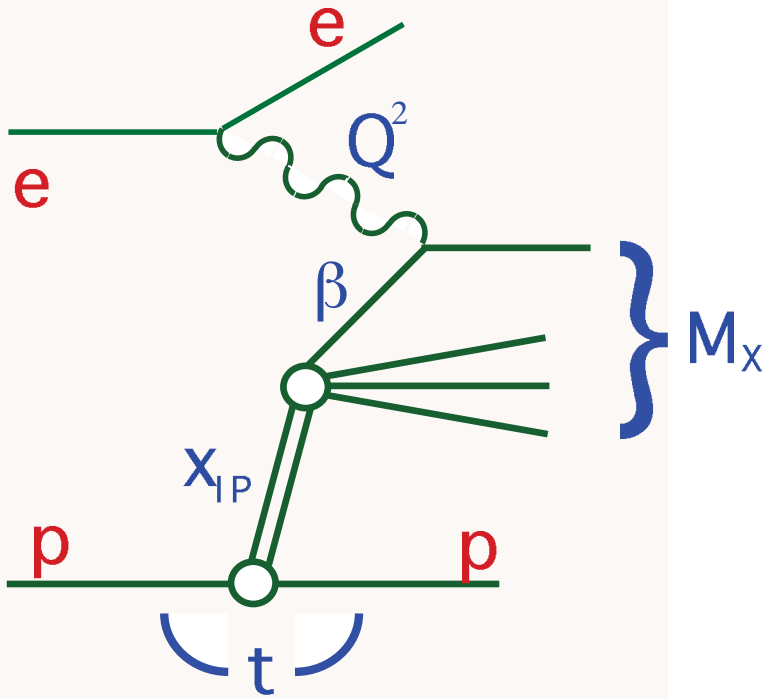
- Leading-Twist Bjorken Scaling!
- Requires nonzero orbital angular momentum of quark
- Arises from the interference of Final-State QCD Coulomb phases in S- and P- waves; Wilson line effect; gauge independent
- Relate to the quark contribution to the target proton anomalous magnetic moment and final-state QCD phases
- QCD phase at soft scale!
- New window to QCD coupling and running gluon mass in the IR
- QED S and P Coulomb phases infinite -- difference of phases finite!

$$i \vec{S} \cdot \vec{p}_{jet} \times \vec{q}$$

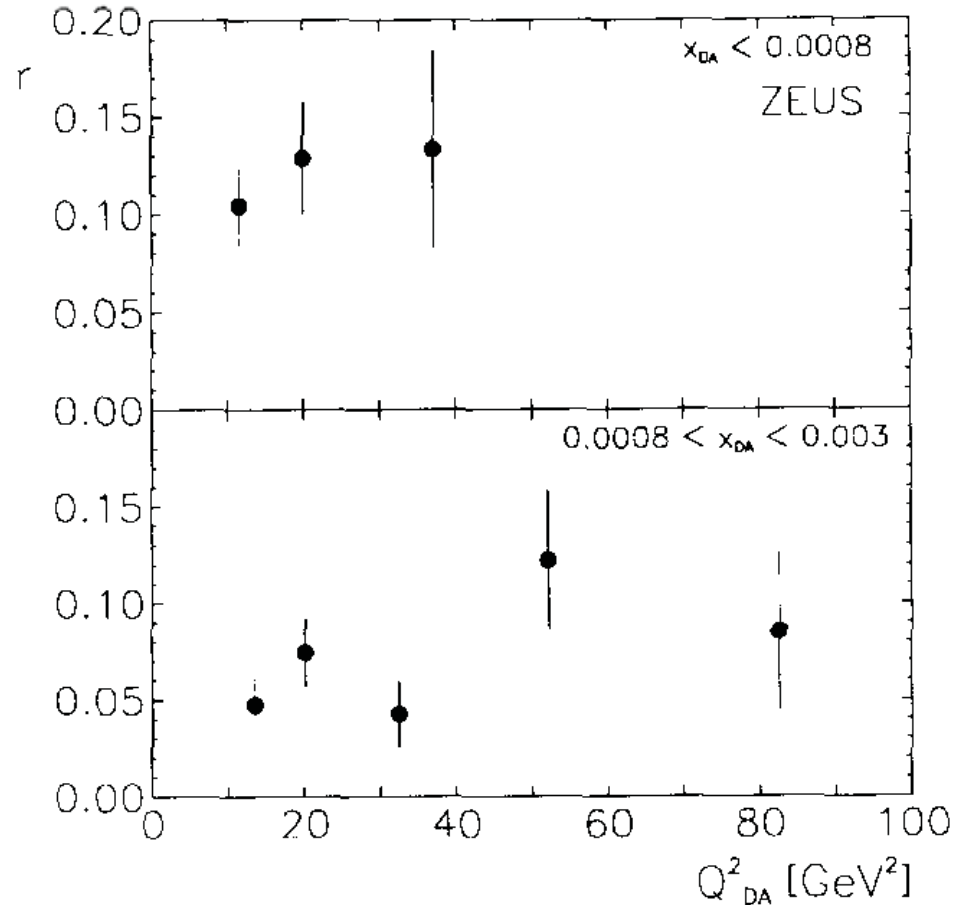


11-2001
8624A06

Remarkable observation at HERA



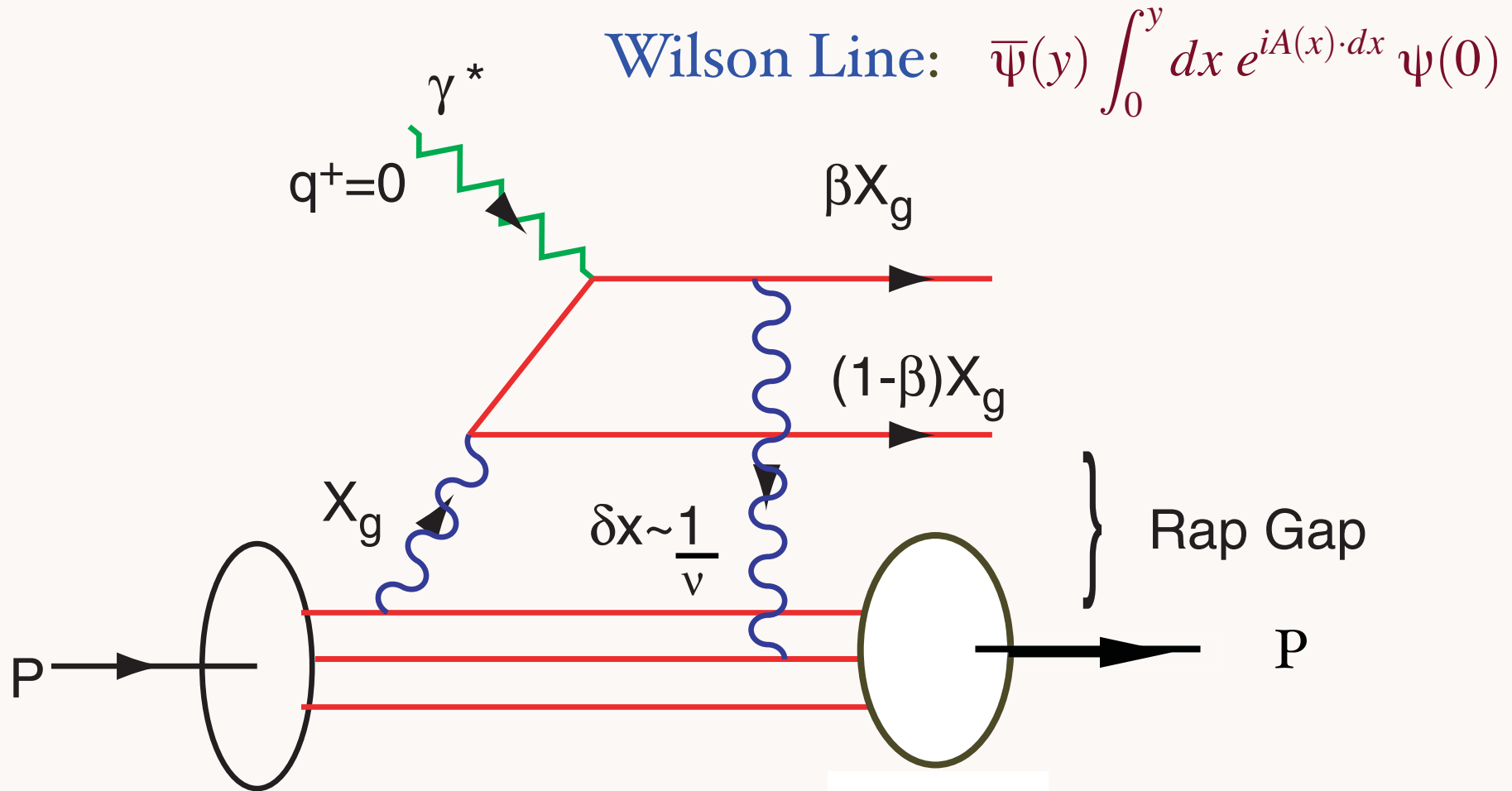
10% to 15%
of DIS events
are
diffractive!



Fraction r of events with a large rapidity gap, $\eta_{\max} < 1.5$, as a function of Q_{DA}^2 for two ranges of x_{DA} . No acceptance corrections have been applied.

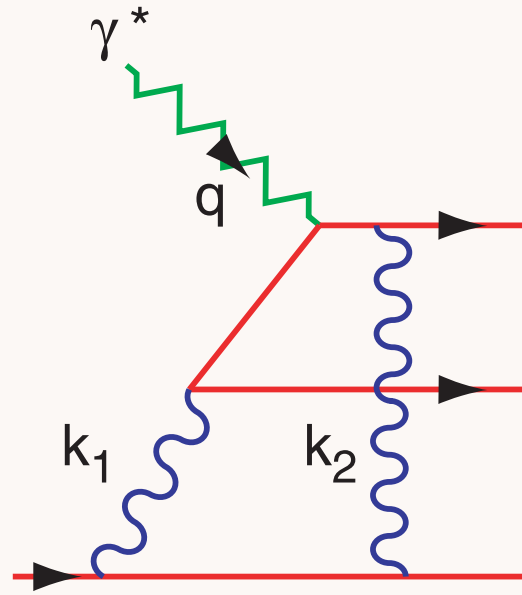
M. Derrick et al. [ZEUS Collaboration], Phys. Lett. B 315, 481 (1993).

QCD Mechanism for Rapidity Gaps

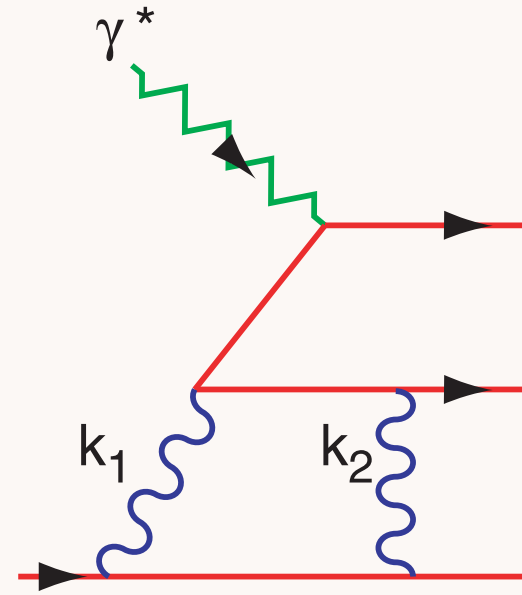


Reproduces lab-frame color dipole approach

Final State Interactions in QCD



Feynman Gauge



Light-Cone Gauge

Result is Gauge Independent

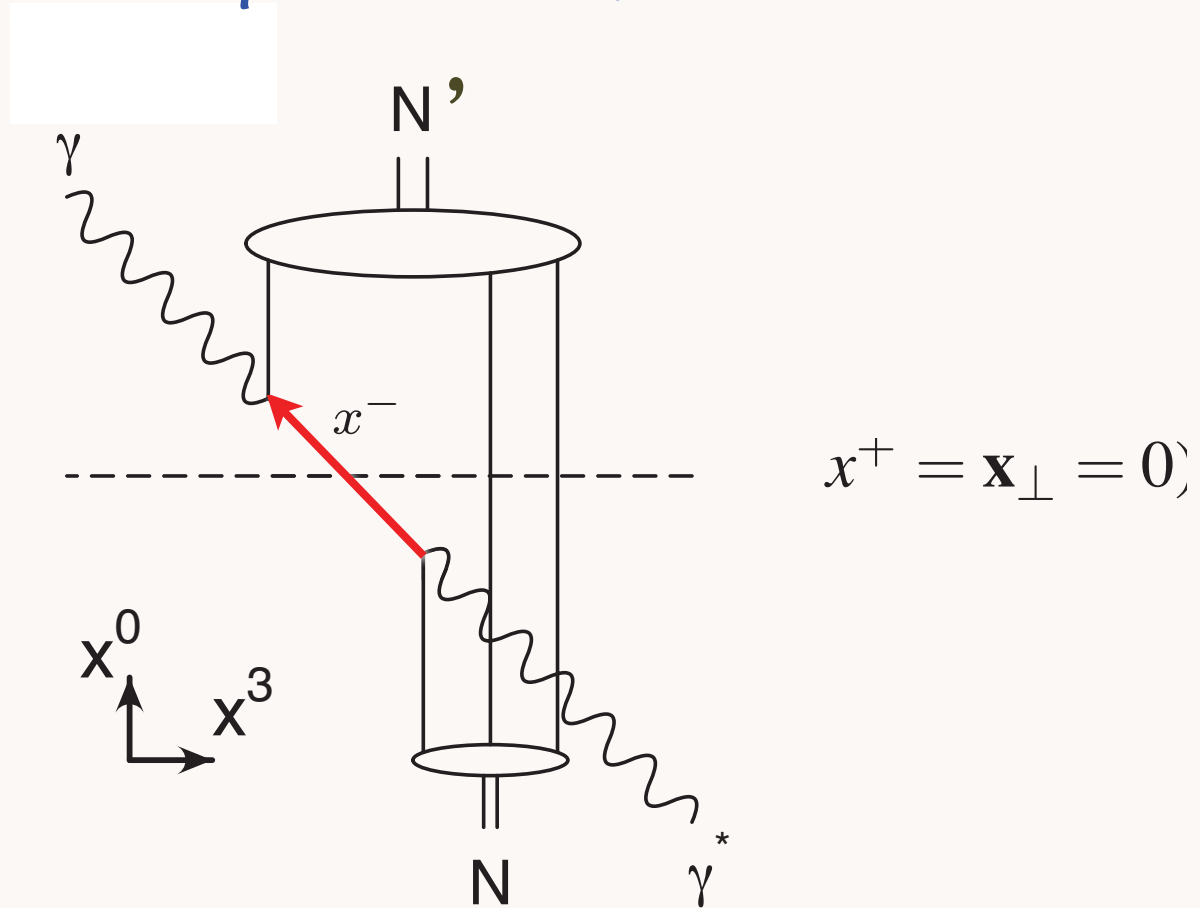
Some Applications of Light-Front Wavefunctions

- Exact formulae for form factors, quark and gluon distributions; vanishing anomalous gravitational moment; edm connection to anm
- Deeply Virtual Compton Scattering, generalized parton distributions, angular momentum sum rules
- Exclusive weak decay amplitudes
- Single spin asymmetries: Role of ISI and FSI
- Factorization theorems, DGLAP, BFKL, ERBL Evolution
- Quark interchange amplitude
- Relation of spin, momentum, and other distributions to physics of the hadron itself.

Space-time picture of DVCS

P. Hoyer

$$\sigma = \frac{1}{2}x^- P^+$$



The position of the struck quark differs by x^- in the two wave functions

**Measure x^- distribution from DVCS:
Take Fourier transform of skewness,
the longitudinal momentum transfer**

$$\xi = \frac{Q^2}{2p \cdot q}$$

S. J. Brodsky^a, D. Chakrabarti^b, A. Harindranath^c, A. Mukherjee^d, J. P. Vary^{e,a,f}

Landau Congress
Moscow, June 20, 2008

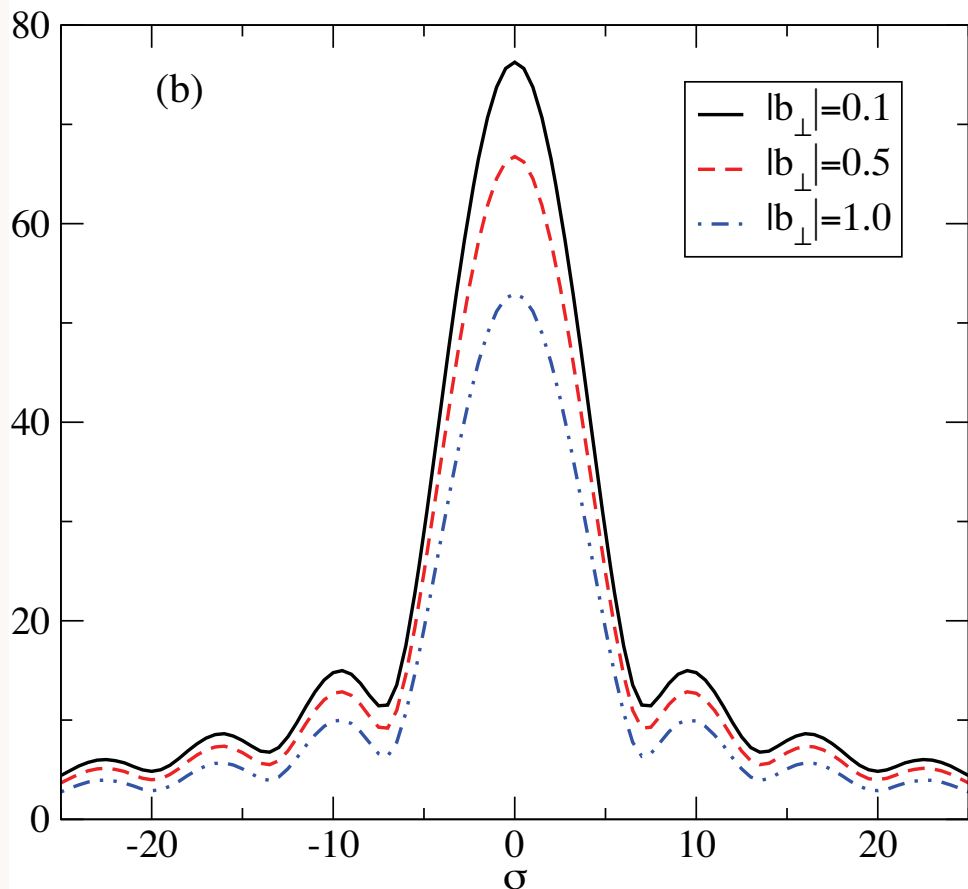
AdS/QCD
133

Stan Brodsky
SLAC & IPPP

Hadron Optics

$$A(\sigma, \vec{b}_\perp) = \frac{1}{2\pi} \int d\xi e^{i\frac{1}{2}\xi\sigma} \tilde{A}(\xi, \vec{b}_\perp)$$

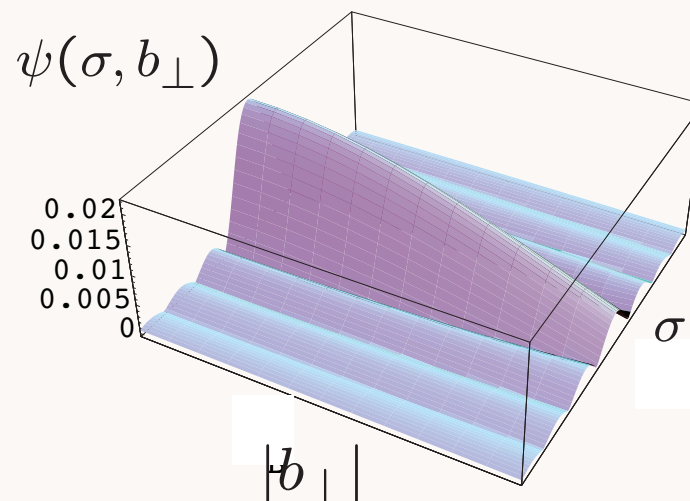
$$\sigma = \frac{1}{2}x^-P^+ \quad \xi = \frac{Q^2}{2p \cdot q}$$



The Fourier Spectrum of the DVCS amplitude in σ space for different fixed values of $|b_\perp|$.
GeV units

DVCS Amplitude using holographic QCD meson LFWF

$$\Lambda_{QCD} = 0.32$$



String Theory



AdS/CFT

Mapping of Poincare' and Conformal $SO(4,2)$ symmetries of 3+1 space to AdS5 space

Goal: First Approximant to QCD

Counting rules for Hard Exclusive Scattering
Regge Trajectories
QCD at the Amplitude Level



AdS/QCD

Conformal behavior at short distances + Confinement at large distance

Semi-Classical QCD / Wave Equations



Holography

Boost Invariant 3+1 Light-Front Wave Equations

$J=0, 1, 1/2, 3/2$ plus L

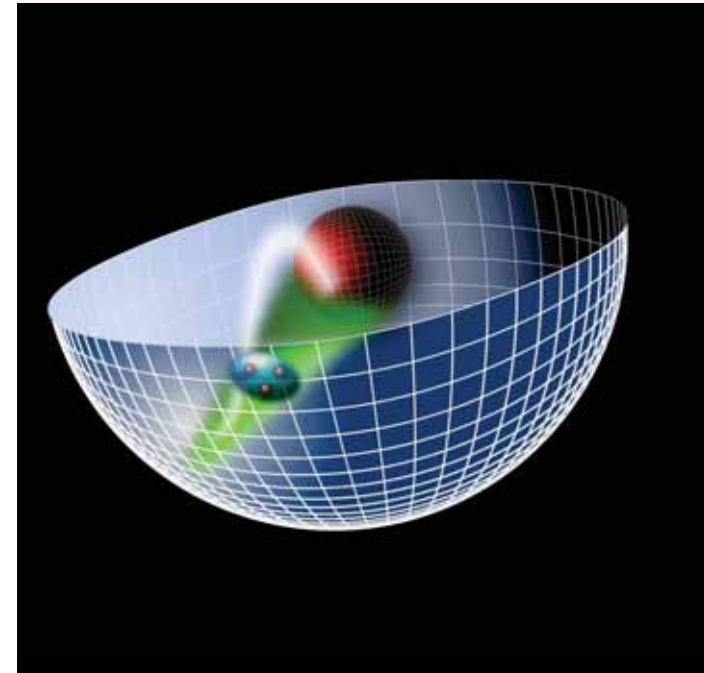
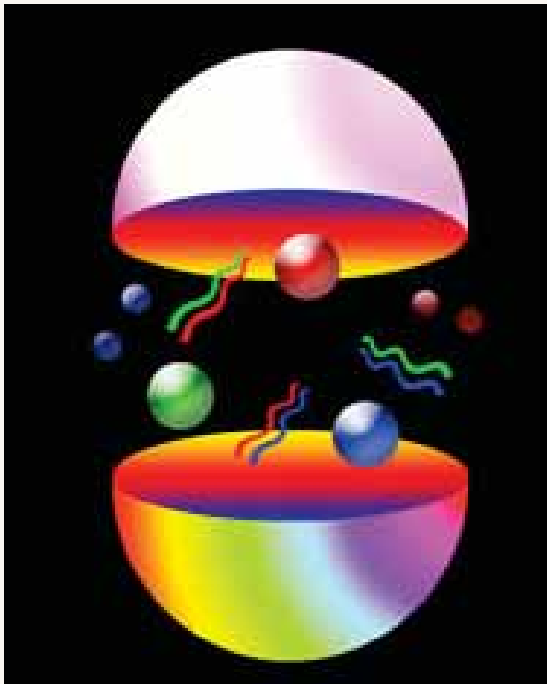
Integrable!



Hadron Spectra, Wavefunctions, Dynamics

AdS/CFT and Hadronic Physics on the Light Front

Lev Davidovich Landau



Stan Brodsky SLAC/IPPP

Landau Memorial Meeting Moscow June 20, 2008

Light-Front Holography and AdS/QCD Correspondence.

[Stanley J. Brodsky](#), [Guy F. de Teramond](#) . SLAC-PUB-13220, Apr 2008. 14pp.
e-Print: [arXiv:0804.3562](#) [hep-ph]

Light-Front Dynamics and AdS/QCD Correspondence: Gravitational Form Factors of Composite Hadrons.

[Stanley J. Brodsky](#) ([SLAC](#)) , [Guy F. de Teramond](#) ([Ecole Polytechnique, CPHT](#) & [Costa Rica U.](#)) . SLAC-PUB-13192, Apr 2008. 12pp. e-Print: [arXiv:0804.0452](#) [hep-ph]

AdS/CFT and Light-Front QCD.

[Stanley J. Brodsky](#), [Guy F. de Teramond](#) . SLAC-PUB-13107, Feb 2008. 38pp.

Invited talk at International School of Subnuclear Physics: 45th Course: Searching for the "Totally Unexpected" in the LHC Era, Erice, Sicily, Italy, 29 Aug - 7 Sep 2007.

e-Print: [arXiv:0802.0514](#) [hep-ph]

AdS/CFT and Exclusive Processes in QCD.

[Stanley J. Brodsky](#), [Guy F. de Teramond](#) . SLAC-PUB-12804, Sep 2007. 29pp. [Temporary entry](#)
e-Print: [arXiv:0709.2072](#) [hep-ph]

Light-Front Dynamics and AdS/QCD Correspondence: The Pion Form Factor in the Space- and Time-Like Regions.

[Stanley J. Brodsky](#) ([SLAC](#)) , [Guy F. de Teramond](#) ([Costa Rica U.](#) & [SLAC](#)) . SLAC-PUB-12554, SLAC-PUB-12544, Jul 2007. 20pp.

Published in **Phys.Rev.D77:056007,2008.**

e-Print: [arXiv:0707.3859](#) [hep-ph]