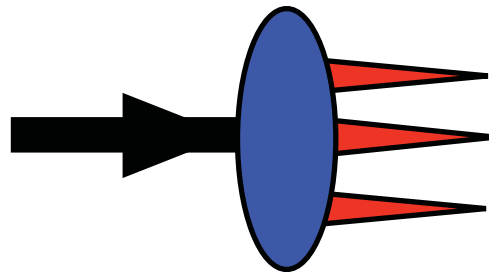
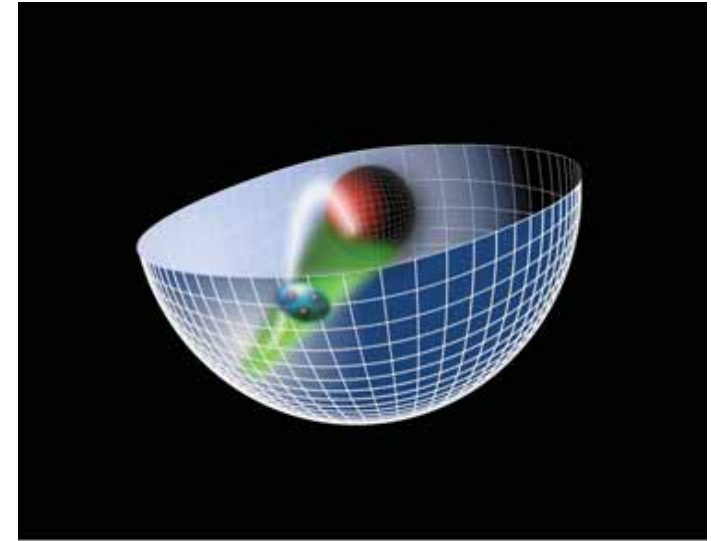


QCD at the Light-Front



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$



LC 2010
university of valencia

June 17, 2010



Stan Brodsky



& CP3 - Origins

Light Cone 2010

Relativistic Hadronic and Particle Physics

Valencia, Spain
June, 14-18, 2010

Local Organizing Committee

Chairpersons

- Joannis Papavassiliou (Universidad de Valencia-IFIC, Spain)
- Vicente Vento (Universidad de Valencia-IFIC, Spain)

Scientific Secretaries

- Arlene Cristina Aguilar (Universidade Federal ABC, Brazil)
- Aurore Courtoy (INFN Sezione di Pavia, Italy)
- Jorge Portolés (IFIC, Spain)



I·L·C·A·C, Inc.

The International Light Cone Advisory Committee, Inc.

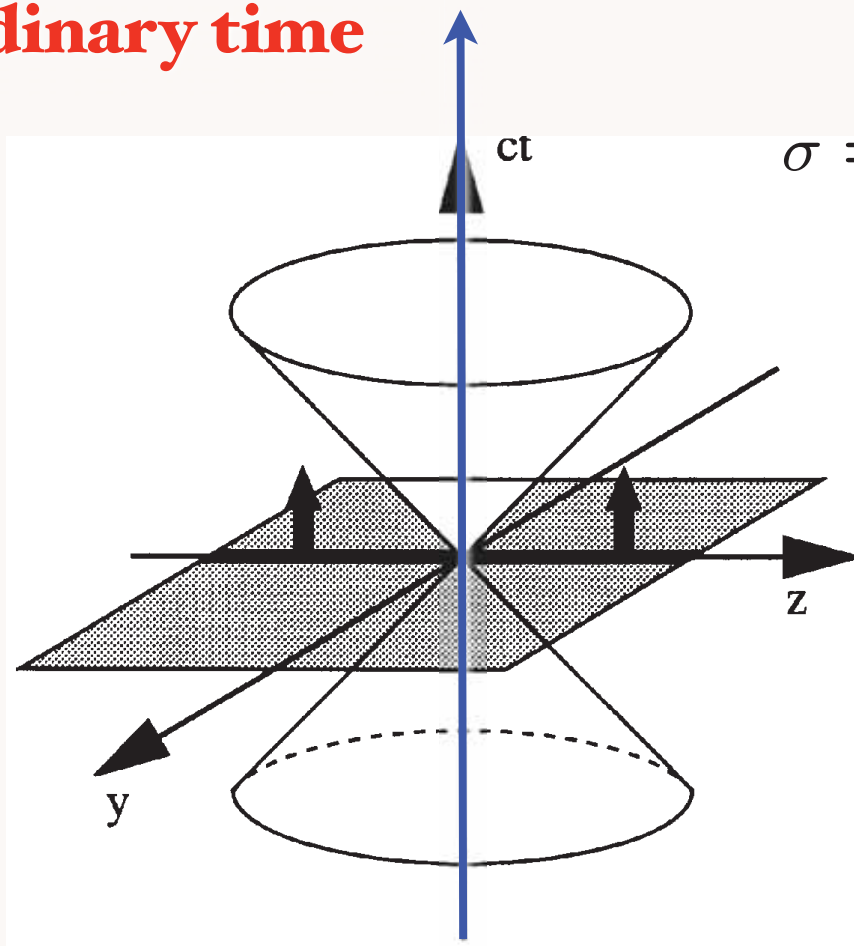
Previous Meetings:

<u>Place</u>	<u>Organizer</u>	<u>Number of Participants</u>
Light Cone 2009	Frederico, Bakker (Pictures)	
Light Cone 2008	Mathiot, Bakker & Diehl: Mulhouse Pictures 1 , 2	58
Amsterdam 2004	Bakker, Boer & Mulders	49
Durham 2003	Stirling, Dalley & Wilkinson	54
Los Alamos 2002	Johnson, Kisslinger, Burkardt & Pate	about 50
Trento2001	Bassetto, Griguolo, Nardelli & Vian	77
Heidelberg2000	Pauli, Werner & Hollenberg	57
Adelaide99	Hollenberg, McKellar & Williams	55
StPetersburg98	Lipatov	70
LesHouches97	Grange	about 80
Ames96	Vary	75
Regensburg95	Werner	about 70
Seattle94	Miller	about 80
Zürich93	Wyler	about 60
Dallas92	McCartor	about 70
Heidelberg91	Pauli	about 50

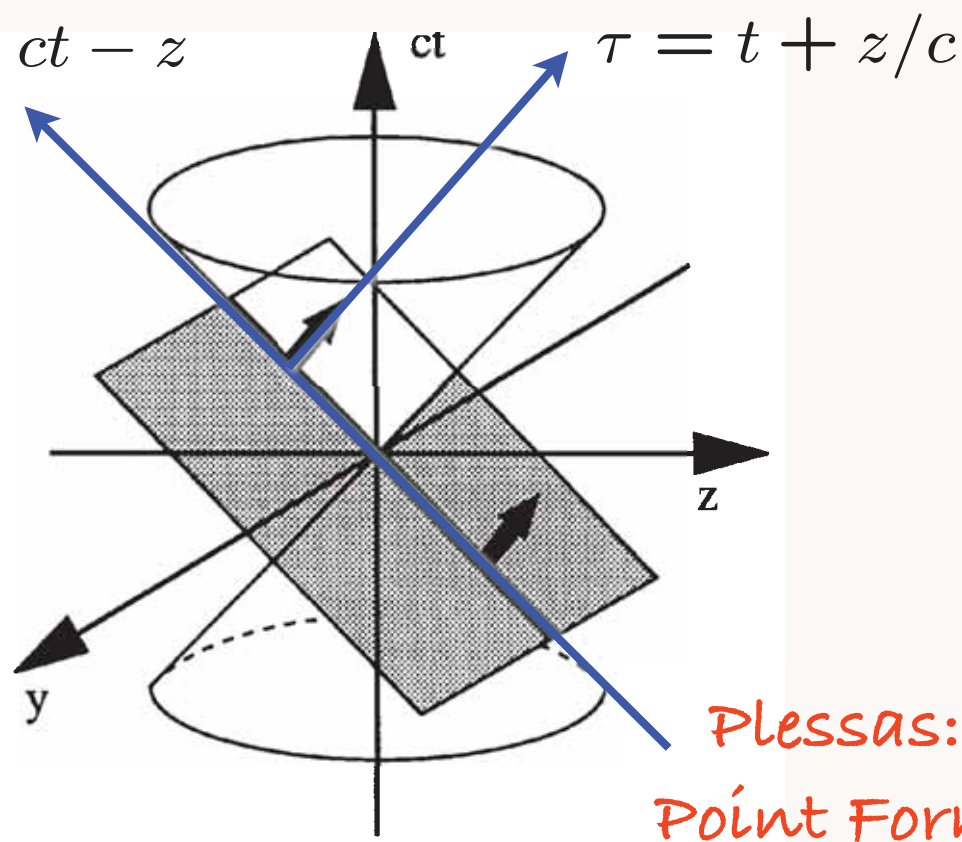
Dirac's Amazing Idea: The Front Form

Evolve in
ordinary time

Evolve in
light-front time!



$$\sigma = ct - z$$



$$\tau = t + z/c$$

Plessas:
Point Form

Instant Form

Front Form

*Each element of
flash photograph
illuminated
at same Light Front
time*

$$\tau = t + z/c$$

Evolve in LF time

$$P^- = i \frac{d}{d\tau}$$

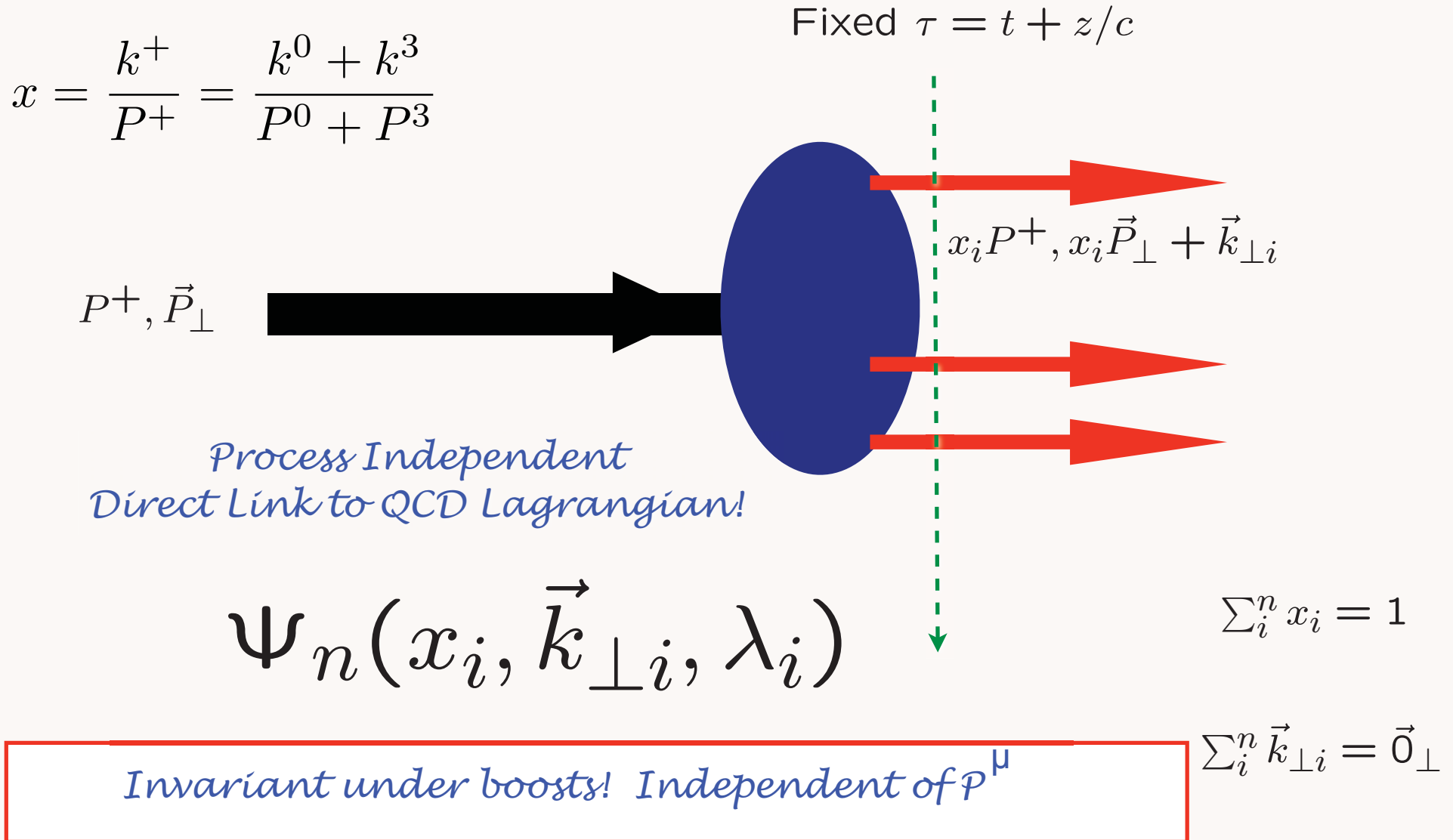
Causal, Trivial Vacuum

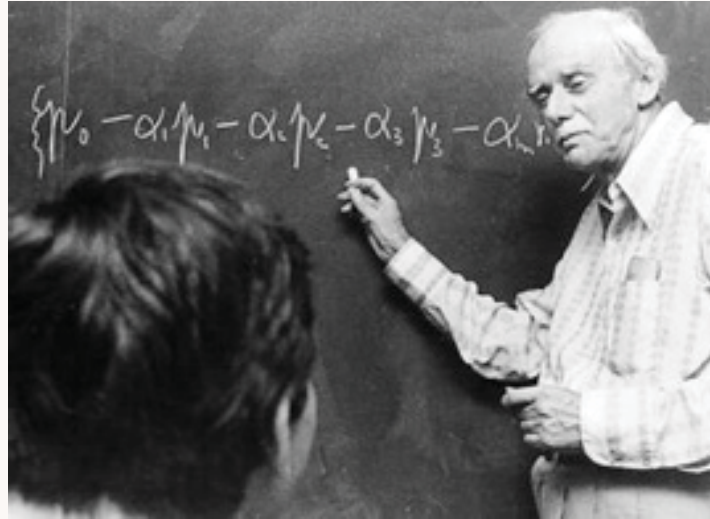
Laser Physics is
Light-Front Physics



T
e.
x
t

Light-Front Wavefunctions: rigorous representation of composite systems in quantum field theory





"Working with a front is a process that is unfamiliar to physicists. But still I feel that the mathematical simplification that it introduces is all-important.

I consider the method to be promising and have recently been making an extensive study of it.

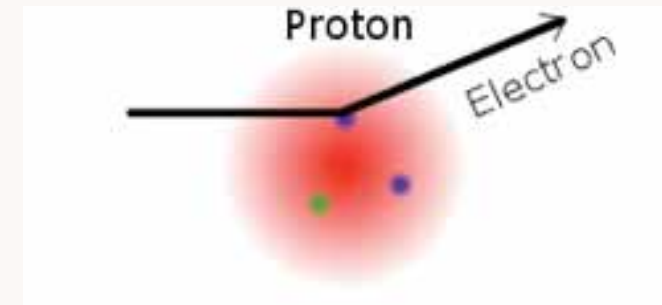
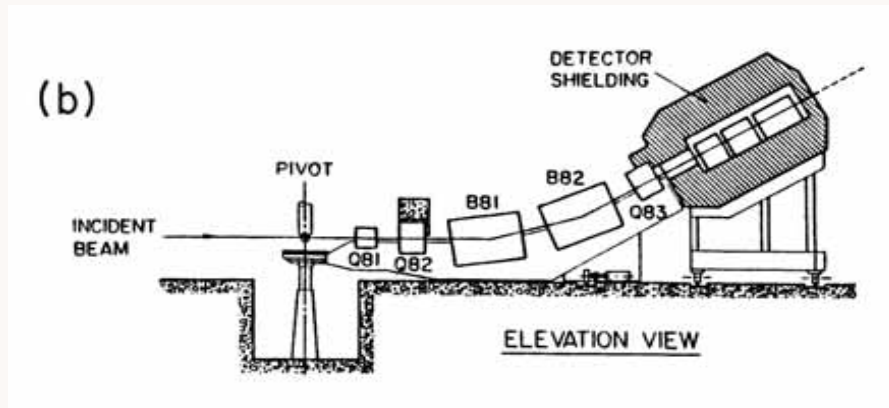
It offers new opportunities, while the familiar instant form seems to be played out." - P.A.M. Dirac (1977)

1967 SLAC Experiment:

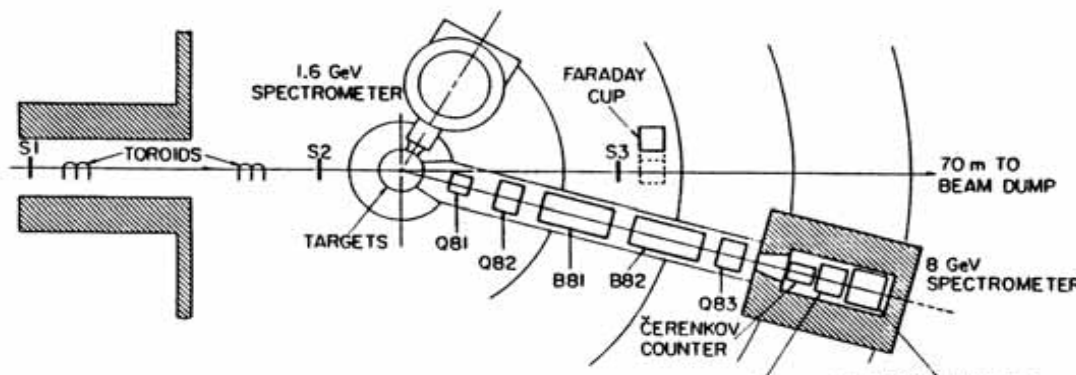
Scatter 20 GeV/c Electrons on protons
in a Hydrogen Target

Discovery of the Quark Structure of Matter

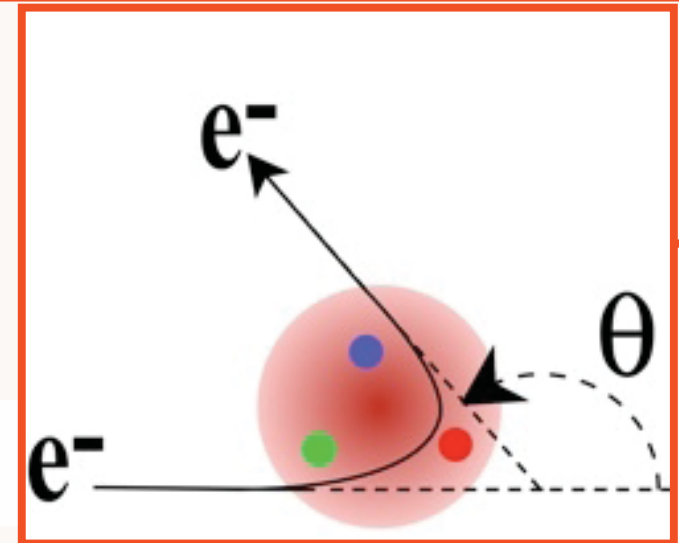
$$ep \rightarrow e'X$$



Discovery of quarks!

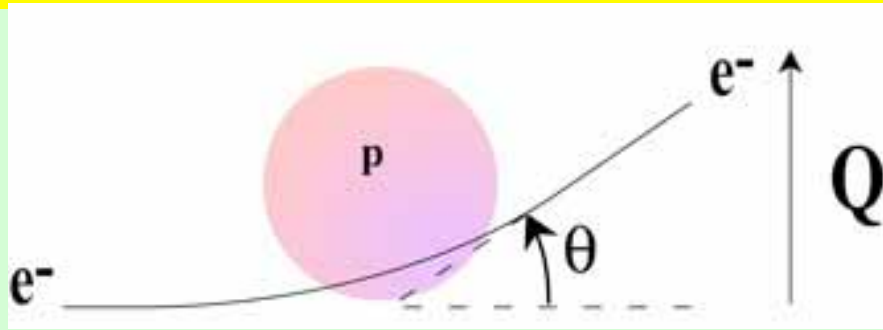


Deep inelastic scattering: Experiments on the proton and the observation of scaling*



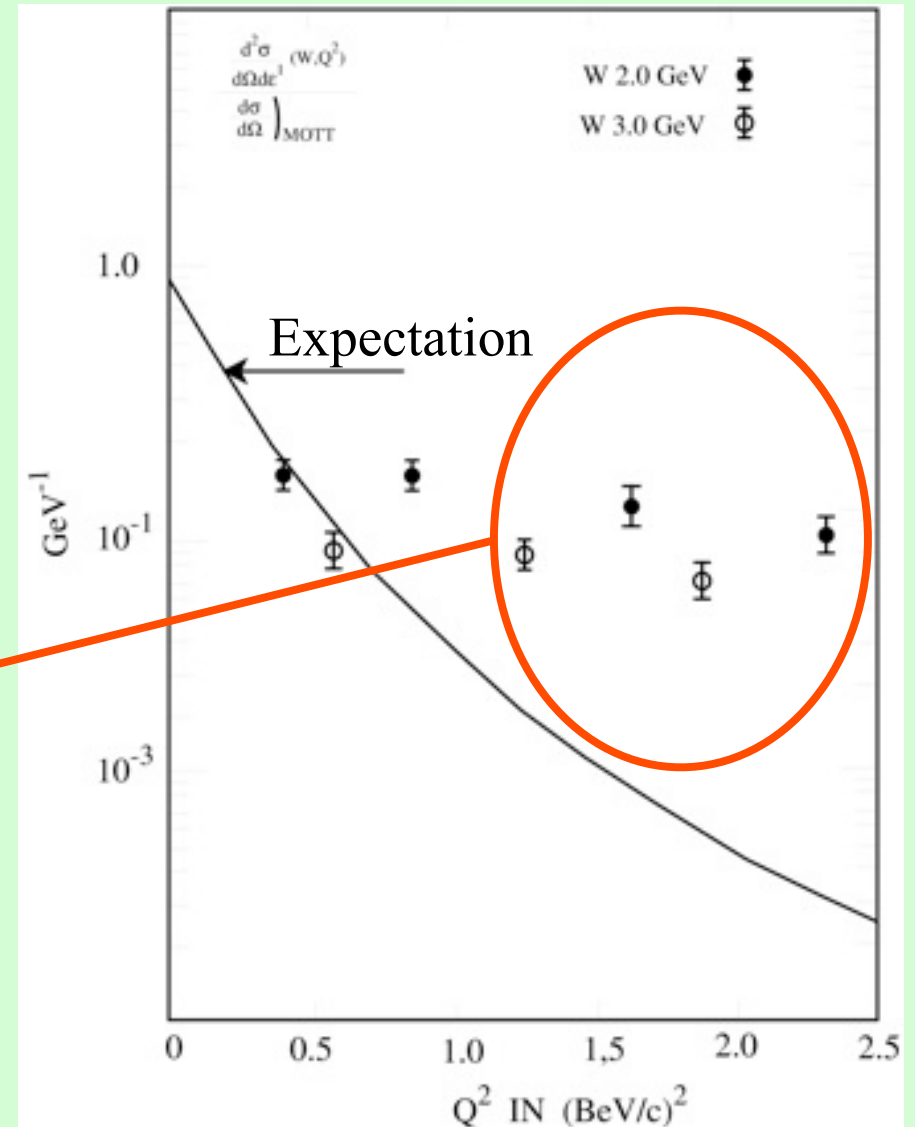
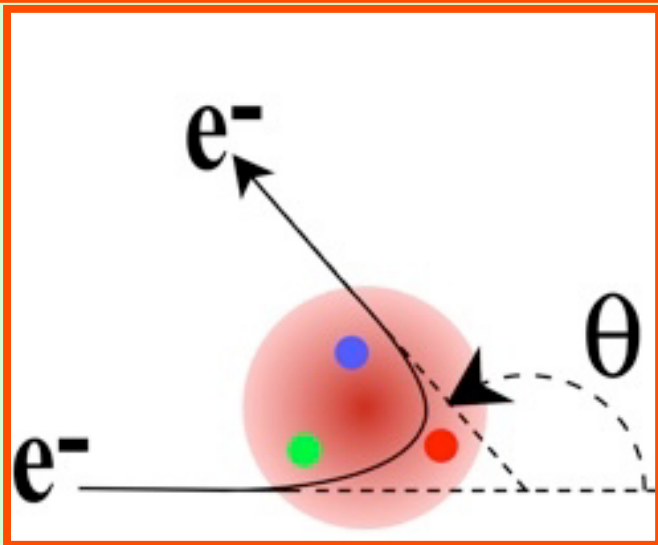
Friedman, Kendall, Taylor: Nobel Prize

Deep inelastic electron-proton scattering

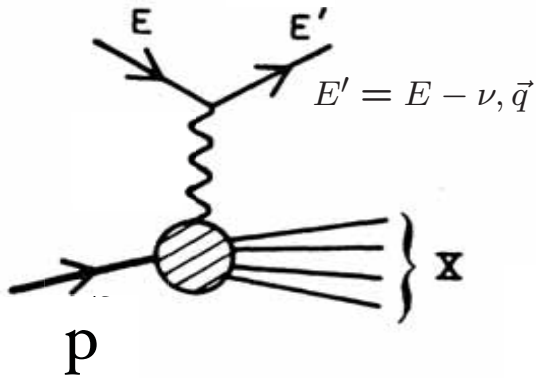


- Rutherford scattering using *very* high-energy electrons striking protons

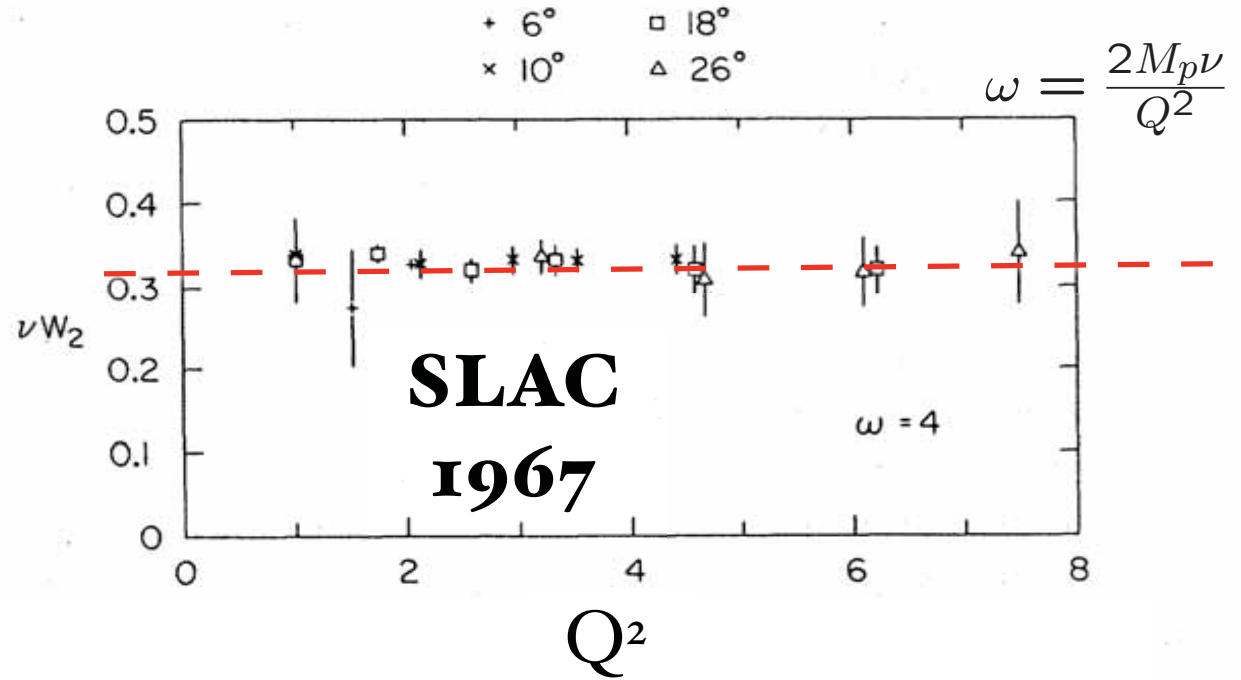
Discovery of quarks!



$$ep \rightarrow e' X$$



$$Q^2 = \vec{q}^2 - \nu^2$$



No intrinsic length scale !

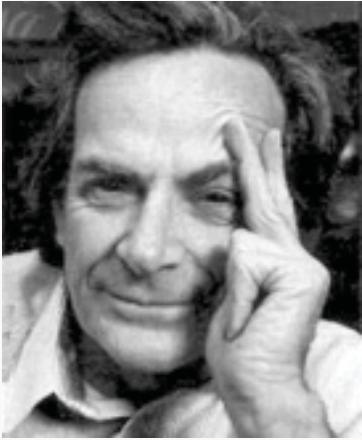
Measure rate as a function of energy loss ν and momentum transfer Q

$$x_{bj} = \frac{1}{\omega_{bj}} = \frac{Q^2}{2q \cdot p}$$

Discovery of Bjorken Scaling

Electron scatters on point-like quarks!

Quarks in the Proton



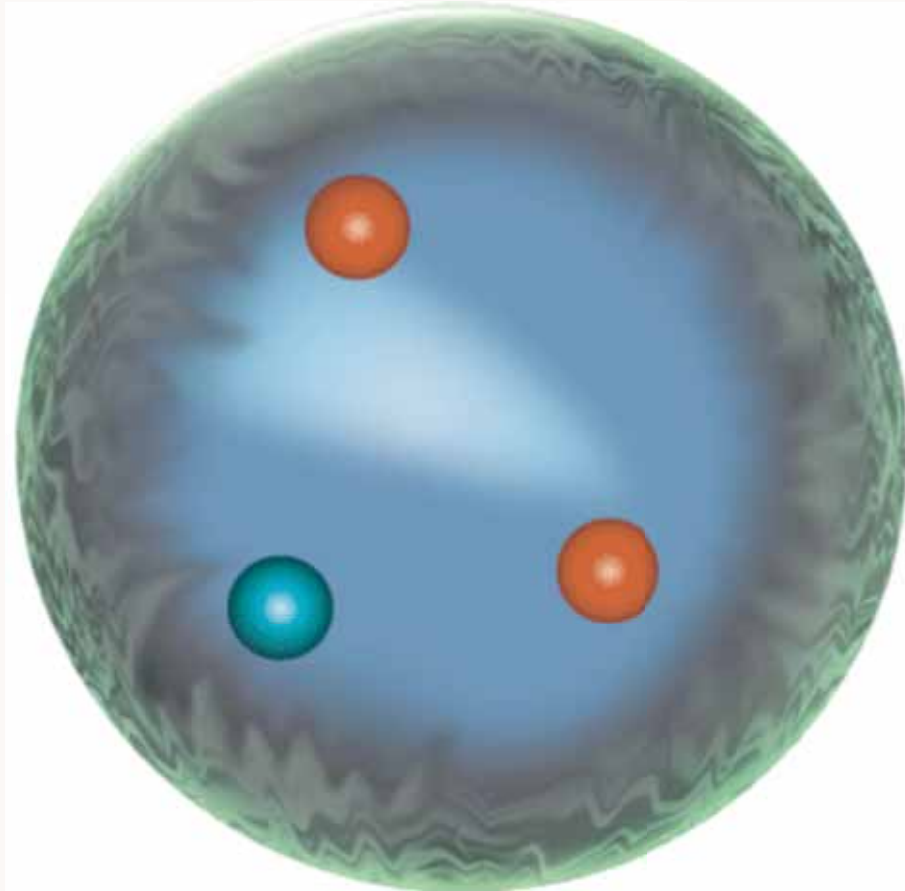
Feynman & Bjorken:
“Parton” model



Bjorken: Scaling

Valencia LC2010
June 17, 2010

$$p = (u u d)$$

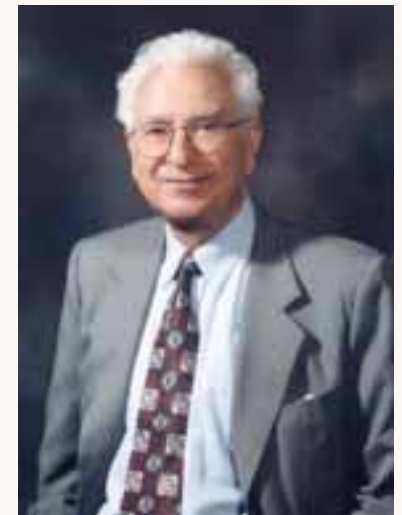


$$1 \text{ fm} \\ 10^{-15} \text{ m} = 10^{-13} \text{ cm}$$

QCD at the Light-Front

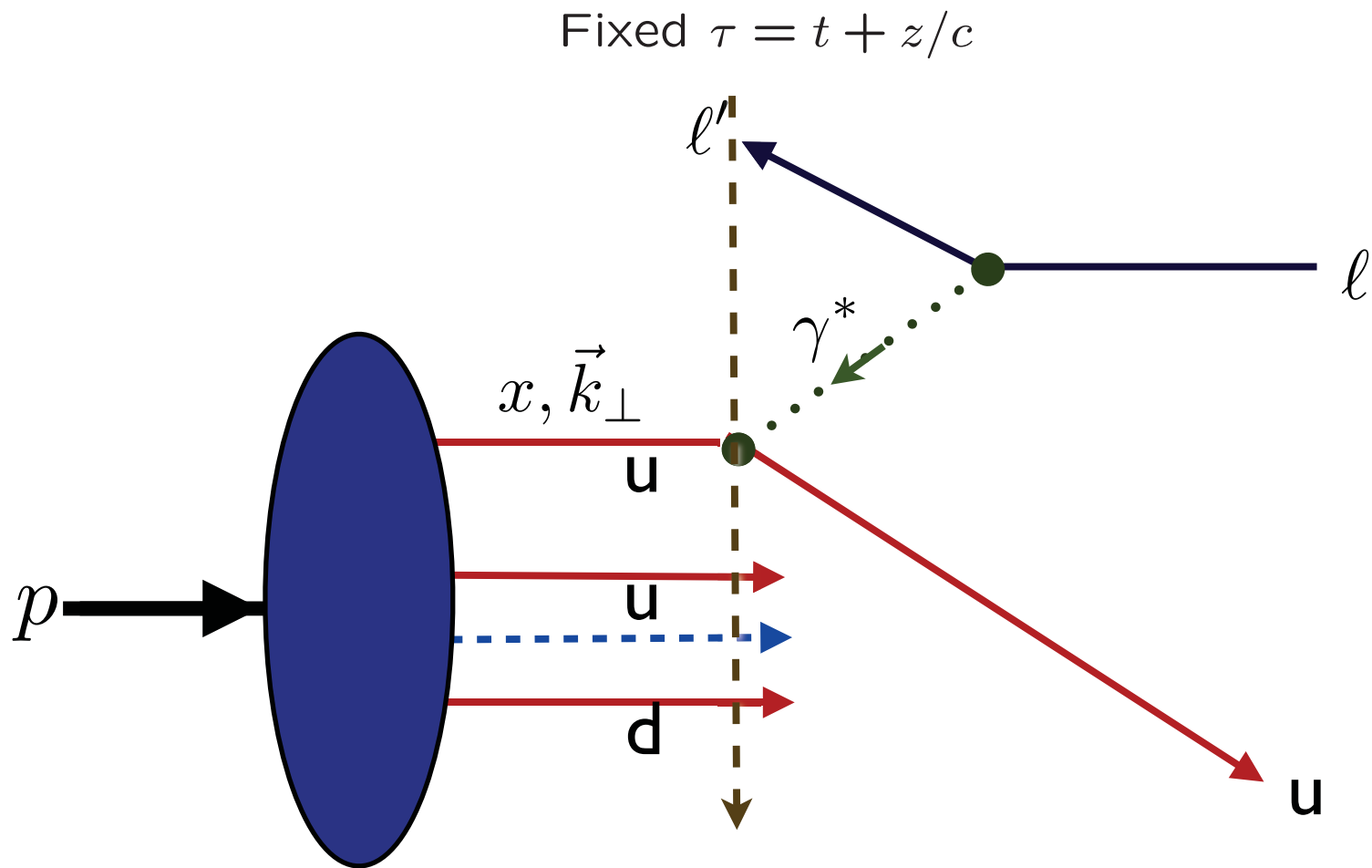


Zweig: “Aces, Deuces,
Treys”



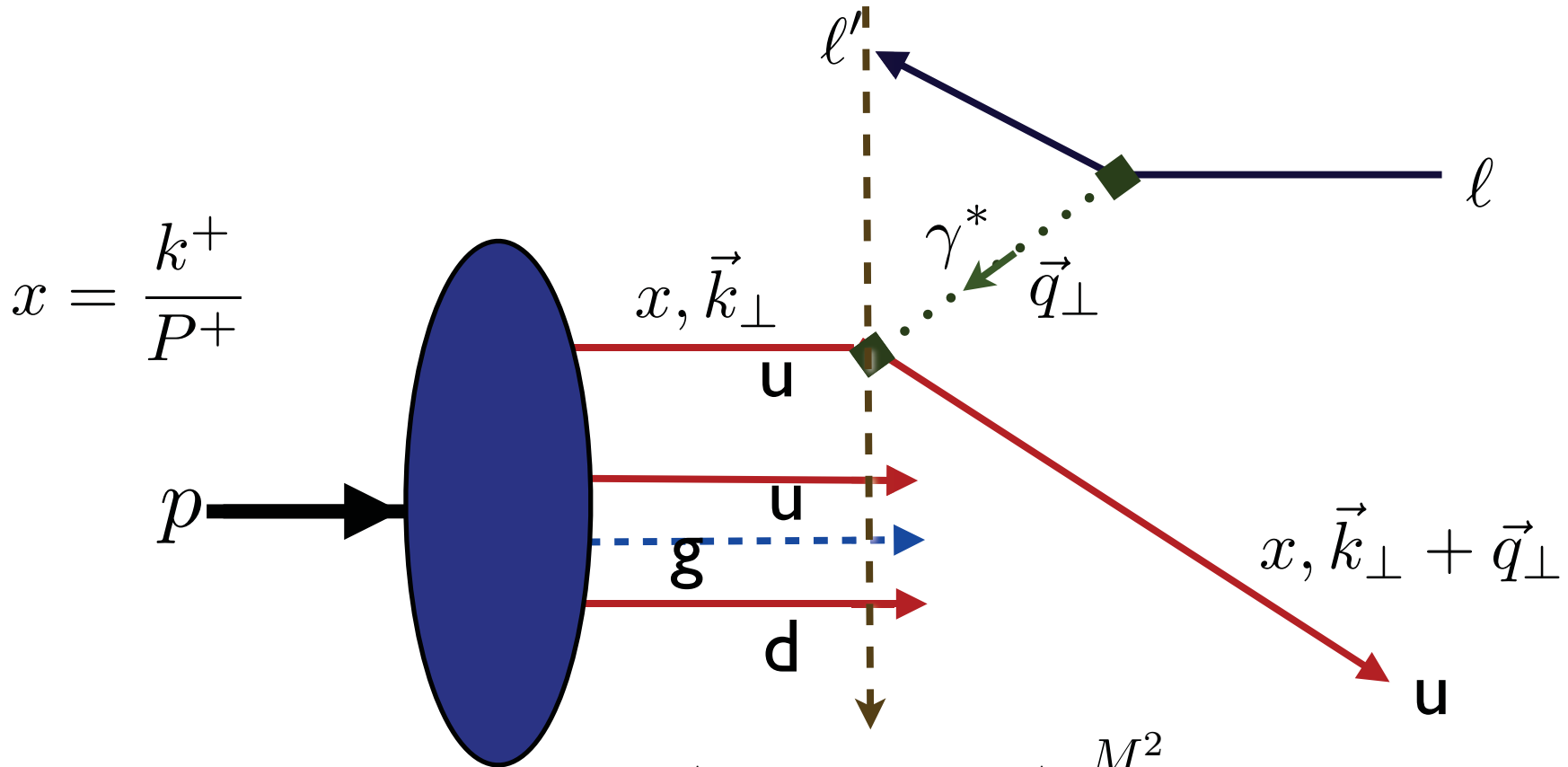
Gell Mann: “Three Quarks for
Mr. Mark”

Stan Brodsky, SLAC & CP³



$$x = \frac{k^+}{p^+} = \frac{k^0 + k^z}{p^0 + p^z} \simeq x_{bj} = \frac{Q^2}{2p \cdot q}$$

Fixed $\tau = t + z/c$



$$p^\mu = (p^+, p^-, \vec{p}_\perp) = (P^+, \frac{M^2}{P^+}, 0_\perp)$$

$$q^\mu = (q^+, q^-, \vec{q}_\perp) = (0, \frac{2q \cdot p}{P^+}, \vec{q}_\perp)$$

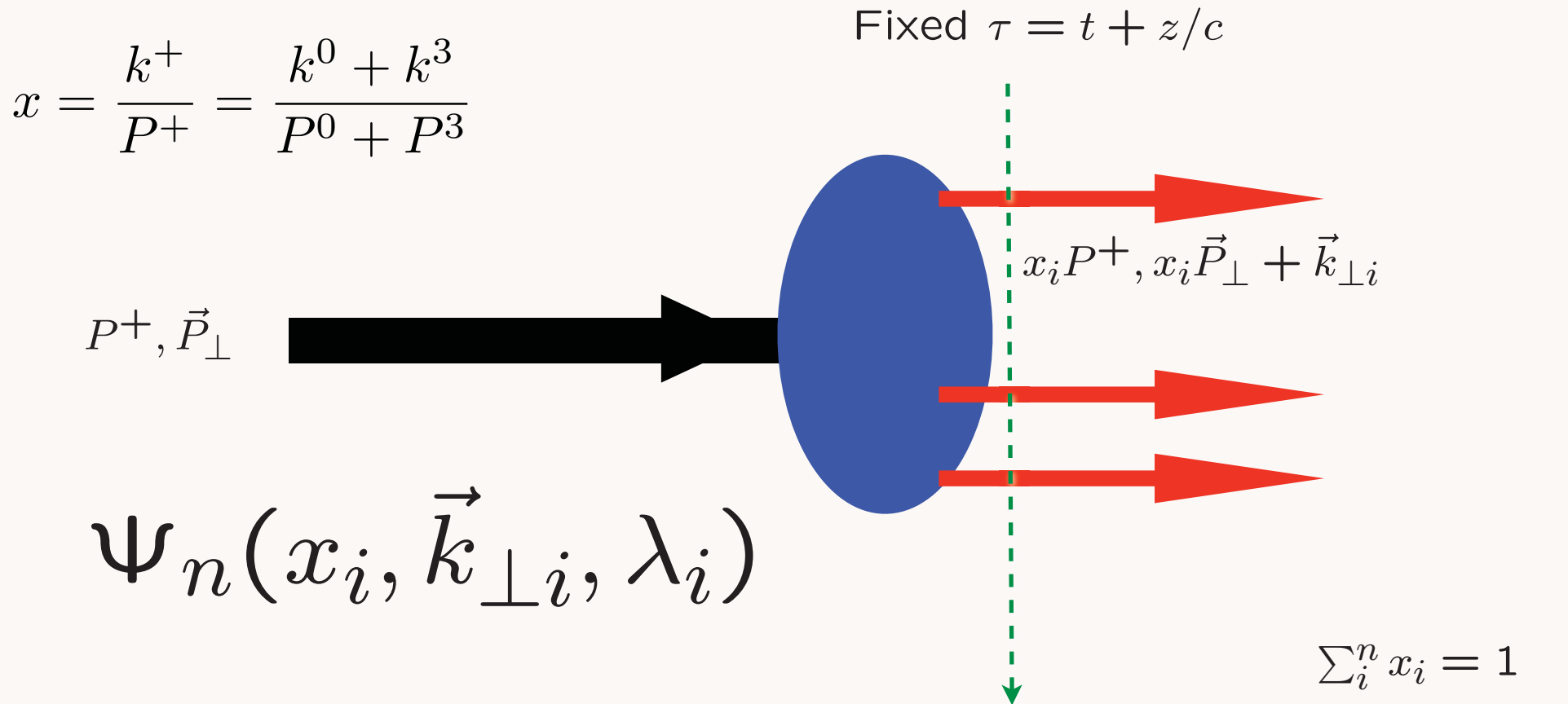
$$Q^2 = \vec{q}_\perp^2$$

P-conservation $2q \cdot p + M^2 = \frac{(\vec{k}_\perp + \vec{q}_\perp)^2 + m^2}{x} + \frac{\vec{k}_\perp^2 + M_s^2}{1-x}$

$$x = \frac{Q^2}{2q \cdot p} = x_{bj}$$

*plus mass,
transverse momentum
and final-state interaction corrections*

Light-Front Wavefunctions: rigorous representation of composite systems in quantum field theory



Structure functions and other distributions computed from the square of the LFWFs

Goal: Predict all features from first principles

QCD Lagrangian

Generalization of QED

The diagram shows the QCD Lagrangian L_{QCD} enclosed in a red box. Labels with arrows point to various parts of the equation:

- gluon dynamics** points to the first term: $-\frac{1}{4g^2} \text{Tr}(G^{\mu\nu} G_{\mu\nu})$
- quark kinetic energy + quark-gluon dynamics** points to the second term: $\sum_{f=1}^{nf} i \bar{\psi}_f D_\mu \gamma^\mu \psi_f$
- mass term** points to the third term: $\sum_{f=1}^{nf} m_f \bar{\psi}_f \psi_f$
- QCD color charge** points to the g^2 in the denominator of the first term.
- field strength tensor** points to $G_{\mu\nu}$ in the first term.
- covariant derivative** points to D_μ in the second term.
- quark field** points to ψ_f in the second term.

Yang Mills Gauge Principle:
Color Rotation and Phase
Invariance at Every Point of
Space and Time

Scale-Invariant Coupling
Renormalizable
Nearly-Conformal
Asymptotic Freedom
Color Confinement

QED Lagrangian

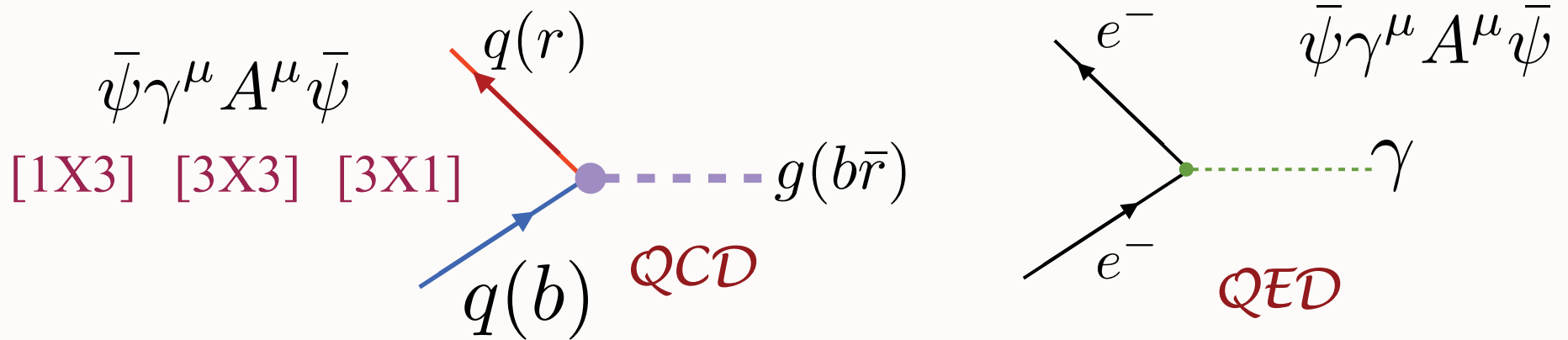
$$\mathcal{L}_{QED} = -\frac{1}{4} \text{Tr}(F^{\mu\nu} F_{\mu\nu}) + \sum_{\ell=1}^{n_\ell} i \bar{\Psi}_\ell D_\mu \gamma^\mu \Psi_\ell + \sum_{\ell=1}^{n_\ell} m_\ell \bar{\Psi}_\ell \Psi_\ell$$

$$iD^\mu = i\partial^\mu - eA^\mu \quad F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$$

Yang Mills Gauge Principle:
Phase Invariance at Every
Point of Space and Time

Scale-Invariant Coupling
Renormalizable
Nearly-Conformal
Landau Pole

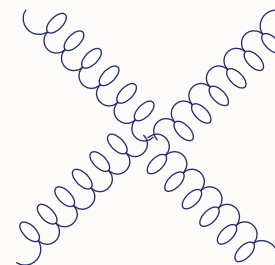
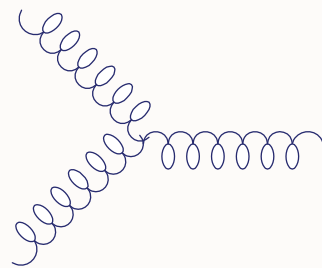
Fundamental Couplings of QED and QCD



$$\mathcal{L}_{QCD} = -\frac{1}{4} \text{Tr}(G^{\mu\nu} G_{\mu\nu}) + \sum_{f=1}^{n_f} i \bar{\Psi}_f D_\mu \gamma^\mu \Psi_f + \sum_{f=1}^{n_f} m_f \bar{\Psi}_f \Psi_f$$

$$G^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu - g[A^\mu, A^\nu]$$

Gluon vertices



$$G^{\mu\nu} G_{\mu\nu}$$

gluon self couplings

Angular Momentum on the Light-Front

LC gauge

$$J^z = \sum_{i=1}^n s_i^z + \sum_{j=1}^{n-1} l_j^z.$$

Conserved
LF Fock state by Fock State

Glueon orbital angular momentum defined in physical lc gauge

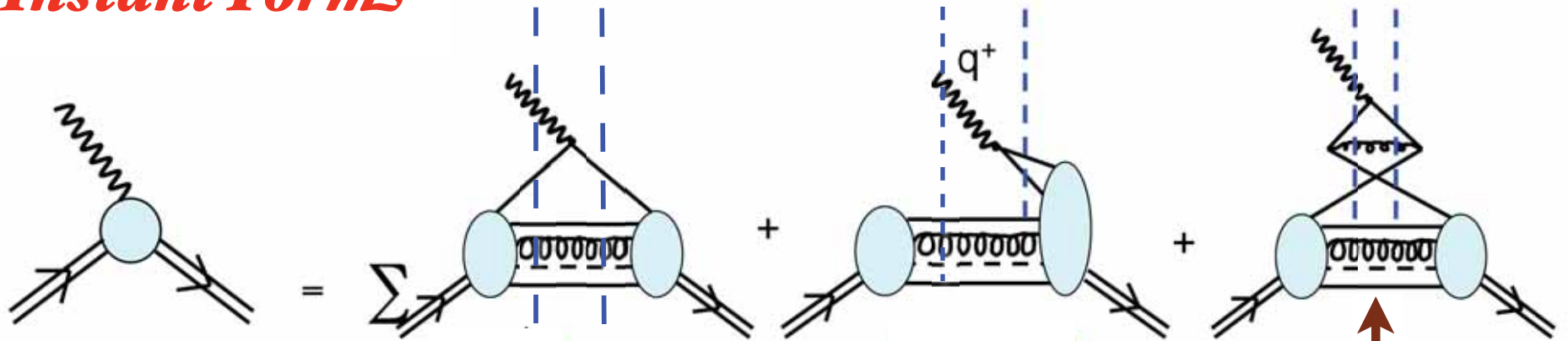
$$l_j^z = -i \left(k_j^1 \frac{\partial}{\partial k_j^2} - k_j^2 \frac{\partial}{\partial k_j^1} \right) \quad n-1 \text{ orbital angular momenta}$$

Orbital Angular Momentum is a property of LFWFS

*Nonzero Anomalous Moment -->
Nonzero quark orbital angular momentum!*

Calculation of Form Factors in Equal-Time Theory

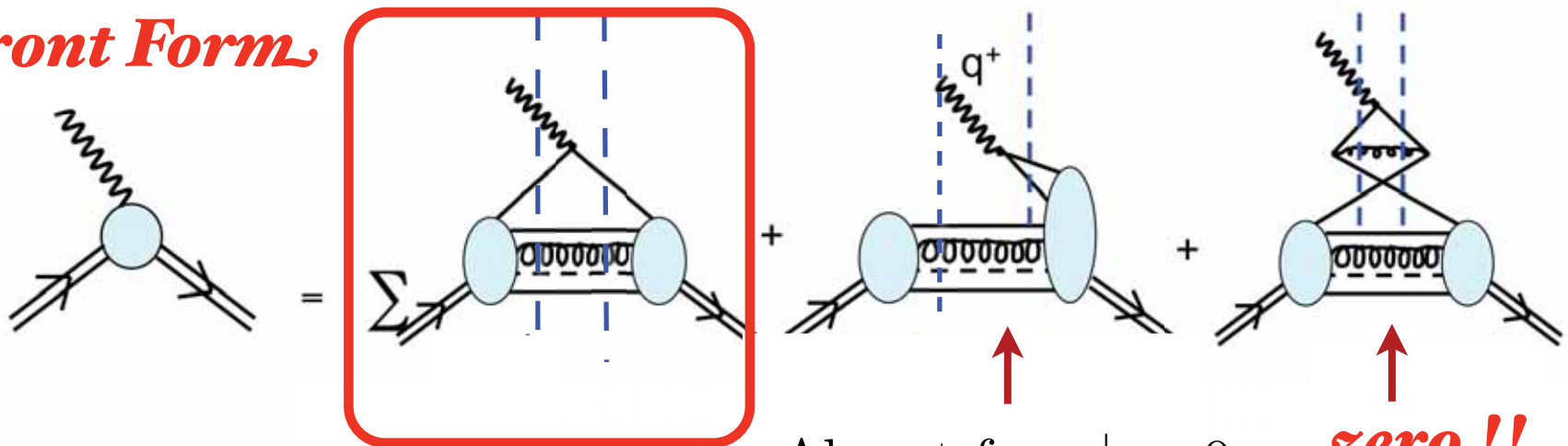
Instant Form



Need vacuum-induced currents

Calculation of Form Factors in Light-Front Theory

Front Form



Complete Answer

Absent for $q^+ = 0$

zero !!

$$\langle p + q | j^+(0) | p \rangle = 2p^+ F(q^2)$$

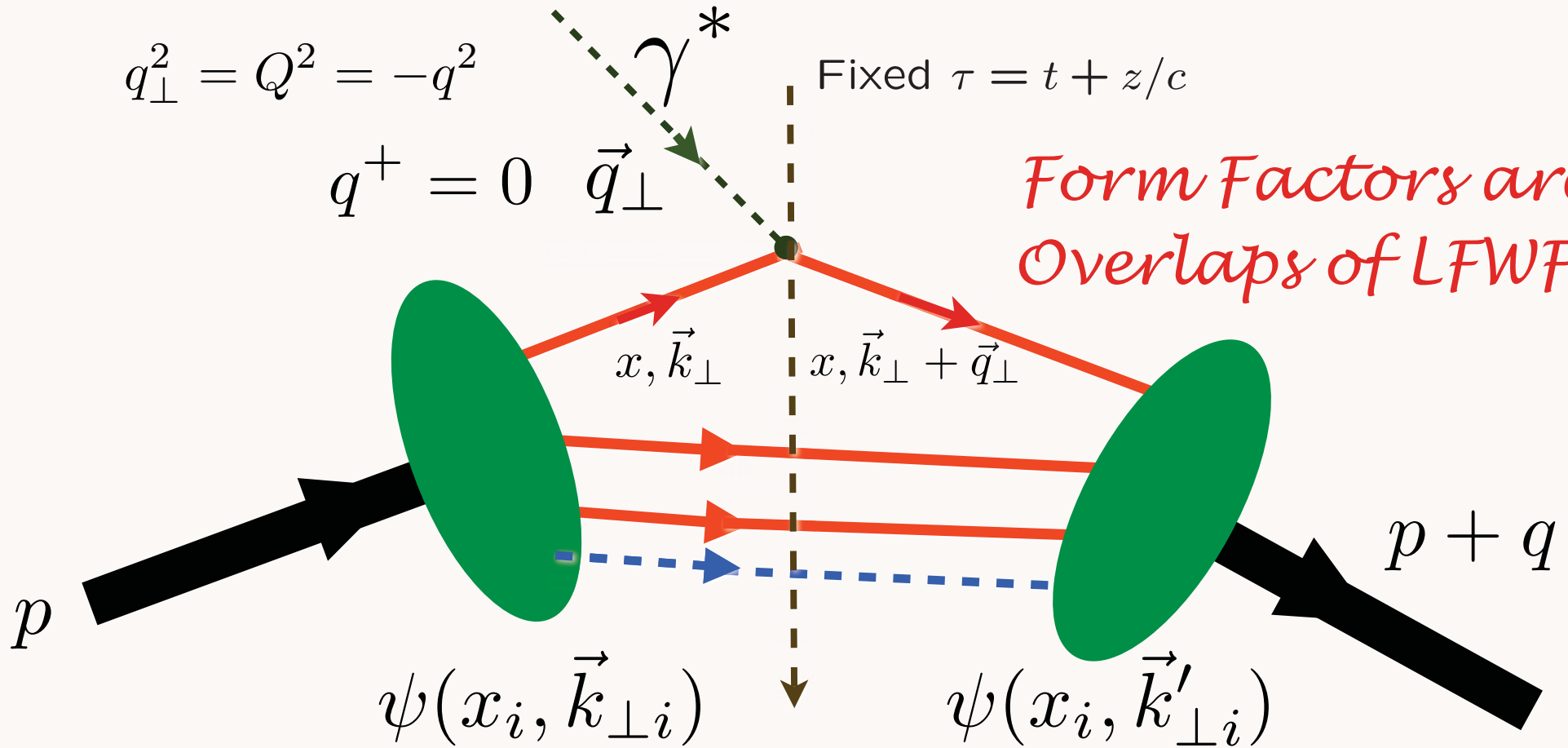
Interaction picture

$$q_{\perp}^2 = Q^2 = -q^2$$

$$q^+ = 0 \quad \vec{q}_{\perp}$$

Fixed $\tau = t + z/c$

Form Factors are Overlaps of LFWFs



Drell & Yan, West

struck $\vec{k}'_{\perp i} = \vec{k}_{\perp i} + (1 - x_i)\vec{q}_{\perp}$

spectators $\vec{k}'_{\perp i} = \vec{k}_{\perp i} - x_i\vec{q}_{\perp}$

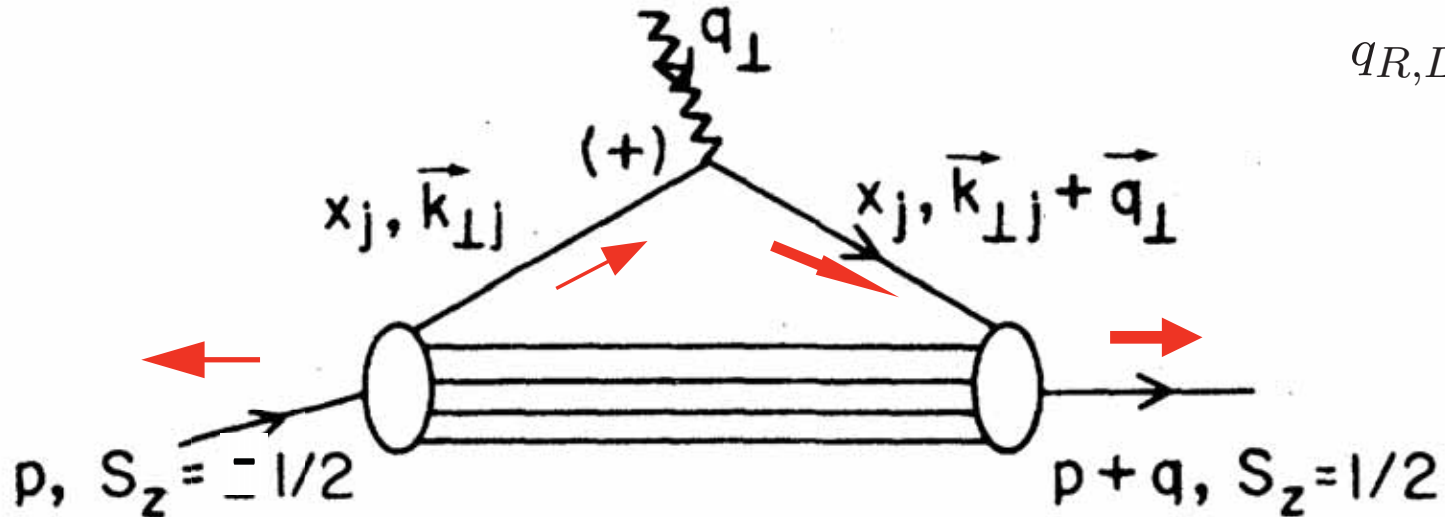
$$\frac{F_2(q^2)}{2M} = \sum_a \int [dx][d^2\mathbf{k}_\perp] \sum_j e_j \frac{1}{2} \times$$

Drell, sjb

$$\left[-\frac{1}{q^L} \psi_a^{\uparrow*}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \psi_a^\downarrow(x_i, \mathbf{k}_{\perp i}, \lambda_i) + \frac{1}{q^R} \psi_a^{\downarrow*}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \psi_a^\uparrow(x_i, \mathbf{k}_{\perp i}, \lambda_i) \right]$$

$$\mathbf{k}'_{\perp i} = \mathbf{k}_{\perp i} - x_i \mathbf{q}_\perp$$

$$\mathbf{k}'_{\perp j} = \mathbf{k}_{\perp j} + (1 - x_j) \mathbf{q}_\perp$$



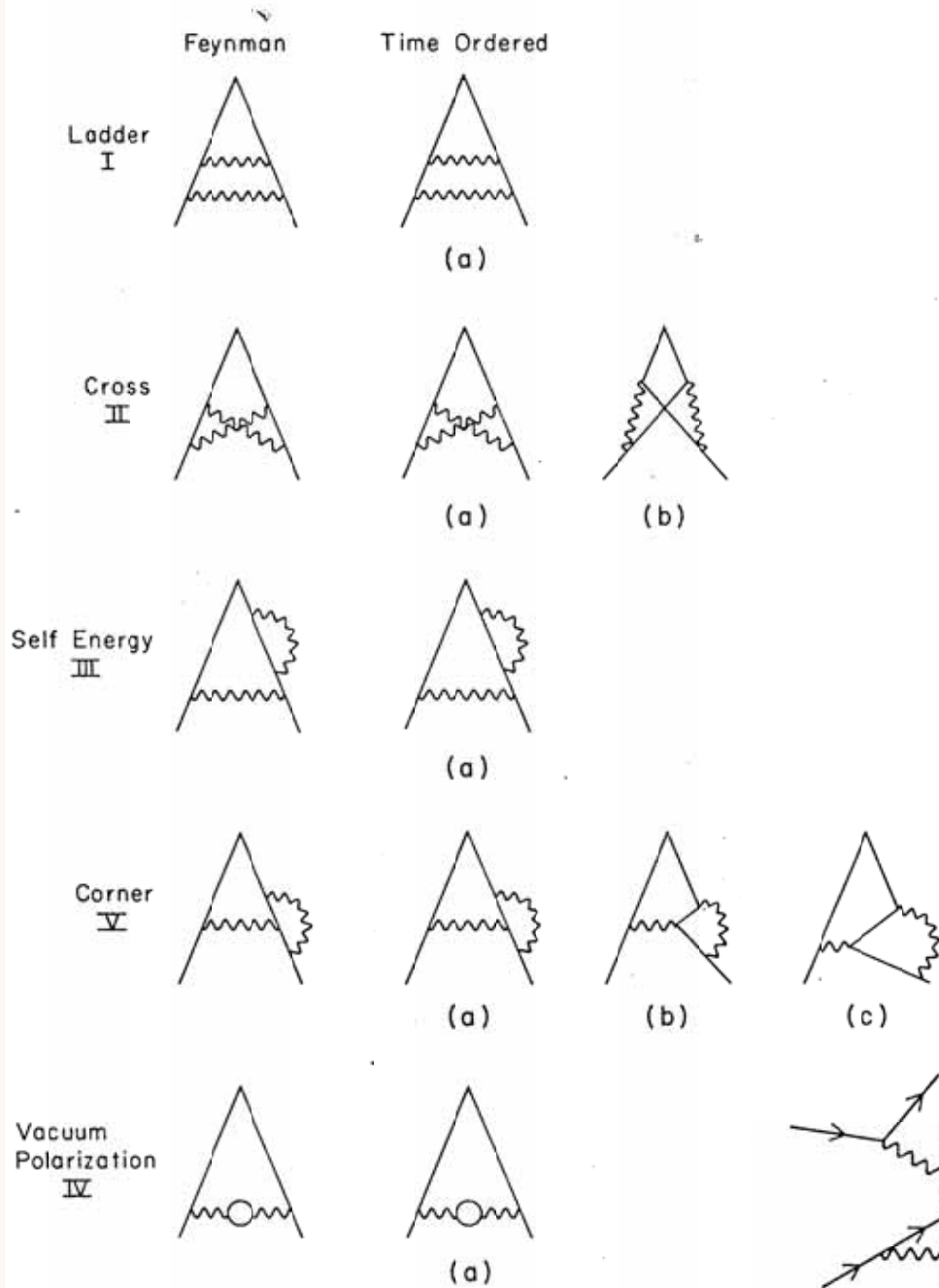
$$q_{R,L} = q^x \pm iq^y$$

Must have $\Delta l_z = \pm 1$ to have nonzero $F_2(q^2)$

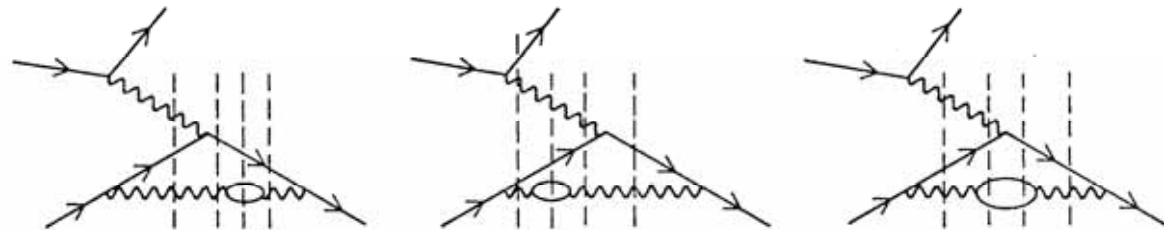
Nonzero Proton Anomalous Moment -->
Nonzero orbital quark angular momentum

QED $g-2$ in LFPth

Roskies, Suaya, and sjb

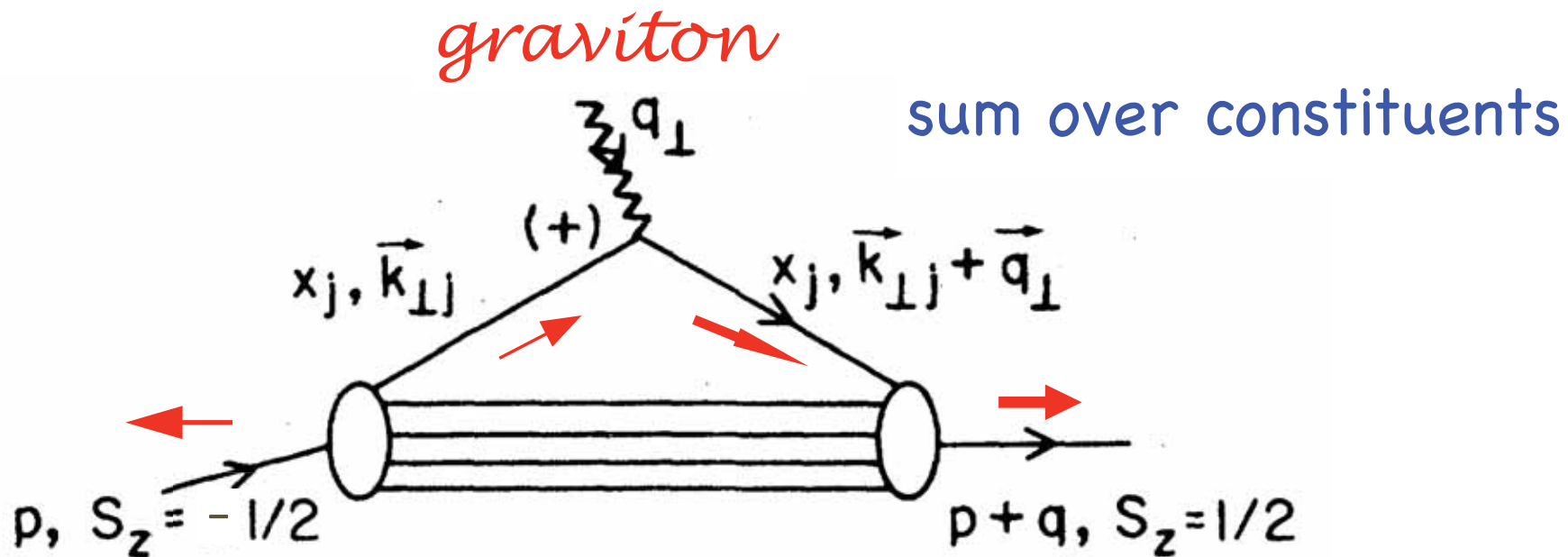


Alternate denominator renormalization



Anomalous gravitomagnetic moment $B(0)$

Terayev, Okun, et al: $B(0)$ Must vanish because of Equivalence Theorem



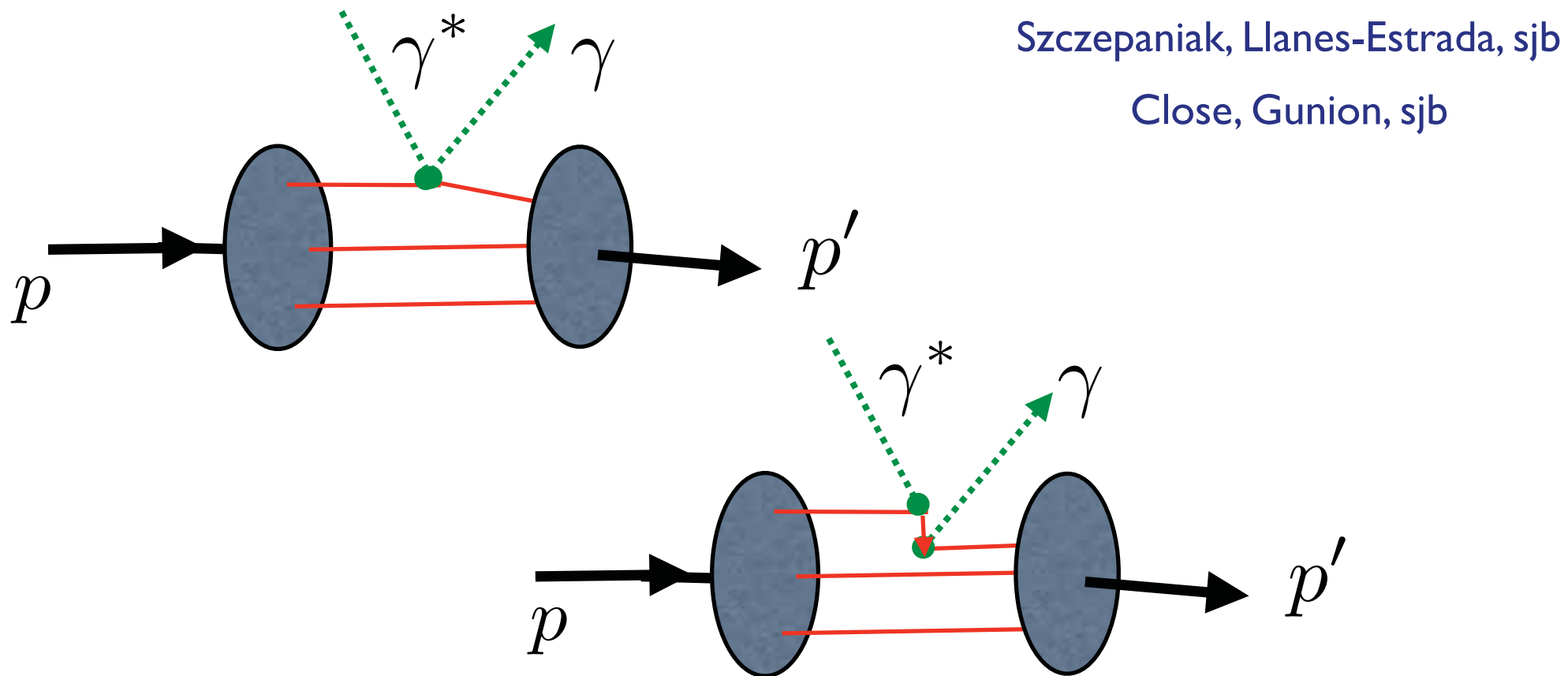
**Hwang, Schmidt, sjb;
Holstein et al**

$B(0) = 0$

Each Fock State

$J=0$ Fixed Pole Contribution to DVCS

- $J=0$ fixed pole -- direct test of QCD locality -- from seagull or instantaneous contribution to Feynman propagator



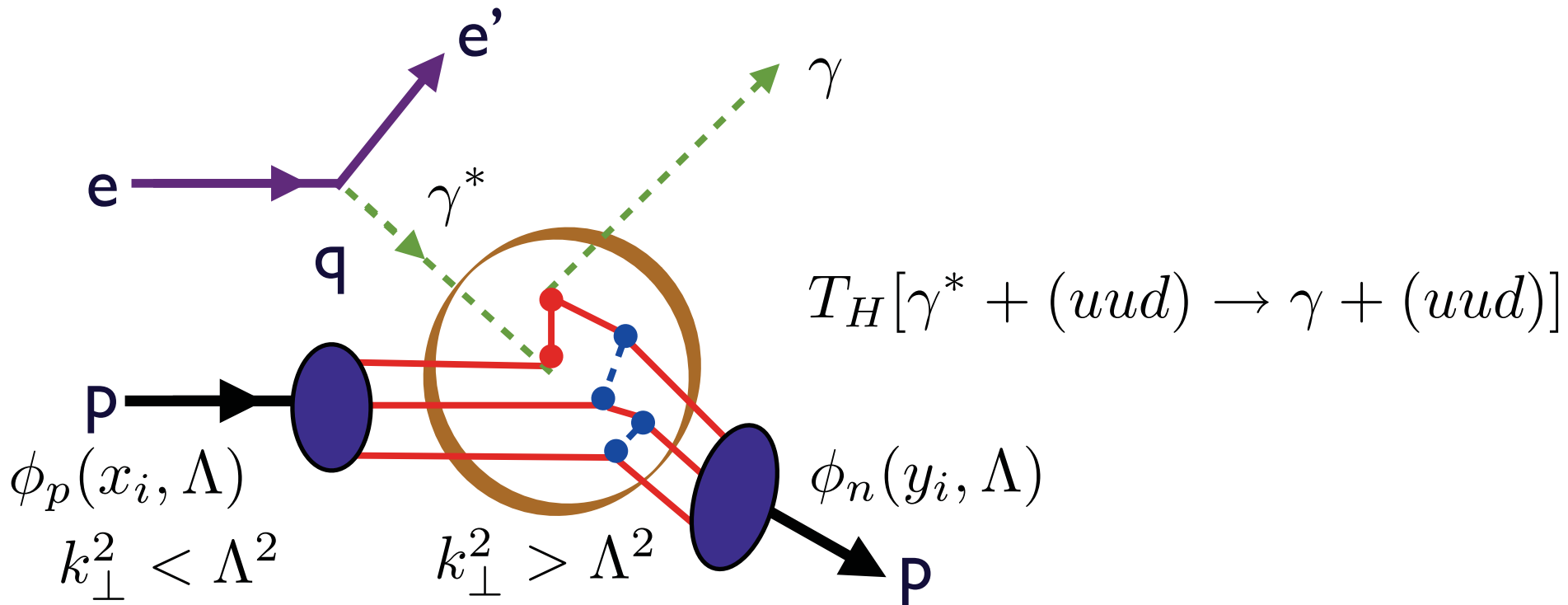
Real amplitude, independent of Q^2 at fixed t

QCD Factorization

DVCS in hard-scattering domain

Lepage, sjb

$$ep \rightarrow e' \gamma p$$

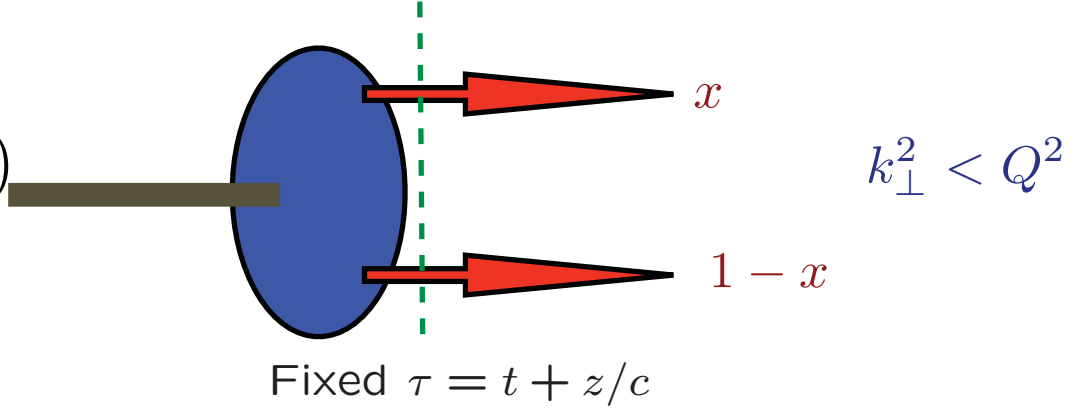


$$T = \int_0^1 dx \int_0^1 dy \int_0^1 dz \phi_p(x, \Lambda) T_H(x, y, z; Q^2, s, t; \Lambda) \phi_n(y, \Lambda) \phi_{\pi}^+(z, \Lambda)$$

Universal distribution amplitudes.
Renormalization Group Invariance:
Renormalization scale is unambiguous -- BLM

$J=0$ Fixed pole from instantaneous gluon

Hadron Distribution Amplitudes

$$\phi_M(x, Q) = \int^Q d^2 \vec{k} \psi_{q\bar{q}}(x, \vec{k}_\perp)$$


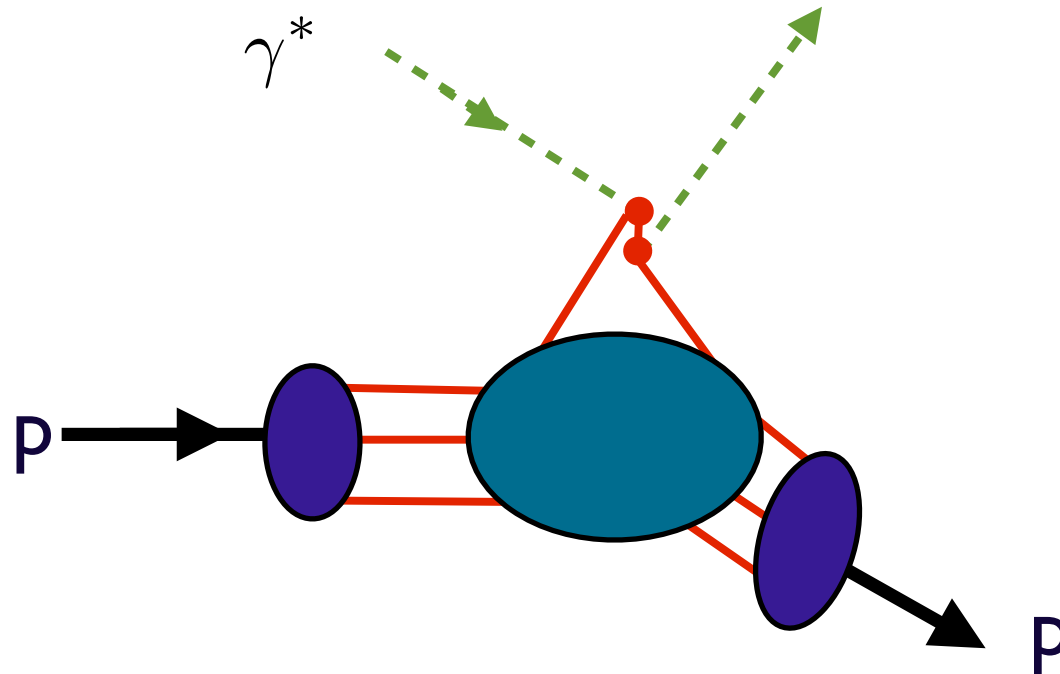
$$\sum_i x_i = 1$$

$k_\perp^2 < Q^2$

- Fundamental gauge invariant non-perturbative input to hard exclusive processes, heavy hadron decays. Defined for Mesons, Baryons *Lepage, sjb*
- Evolution Equations from PQCD, OPE *Lepage, sjb*
Efremov, Radyushkin
- Conformal Invariance *Sachrajda, Frishman Lepage, sjb*
Braun, Gardi
- Compute from valence light-front wavefunction in light-cone gauge

Deeply Virtual Compton Scattering

$$\gamma^* p \rightarrow \gamma p$$



*Seagull interaction
(instantaneous quark
exchange or Z-graph)*

$$s \gg -t, Q^2 \gg \Lambda_{QCD}^2$$

*Hard Reggeon
Domain*

$$T(\gamma^*(q)p \rightarrow \gamma(k) + p) \sim \epsilon \cdot \epsilon' \sum_R s_R^\alpha(t) \beta_R(t)$$

$$\alpha_R(t) \rightarrow 0$$

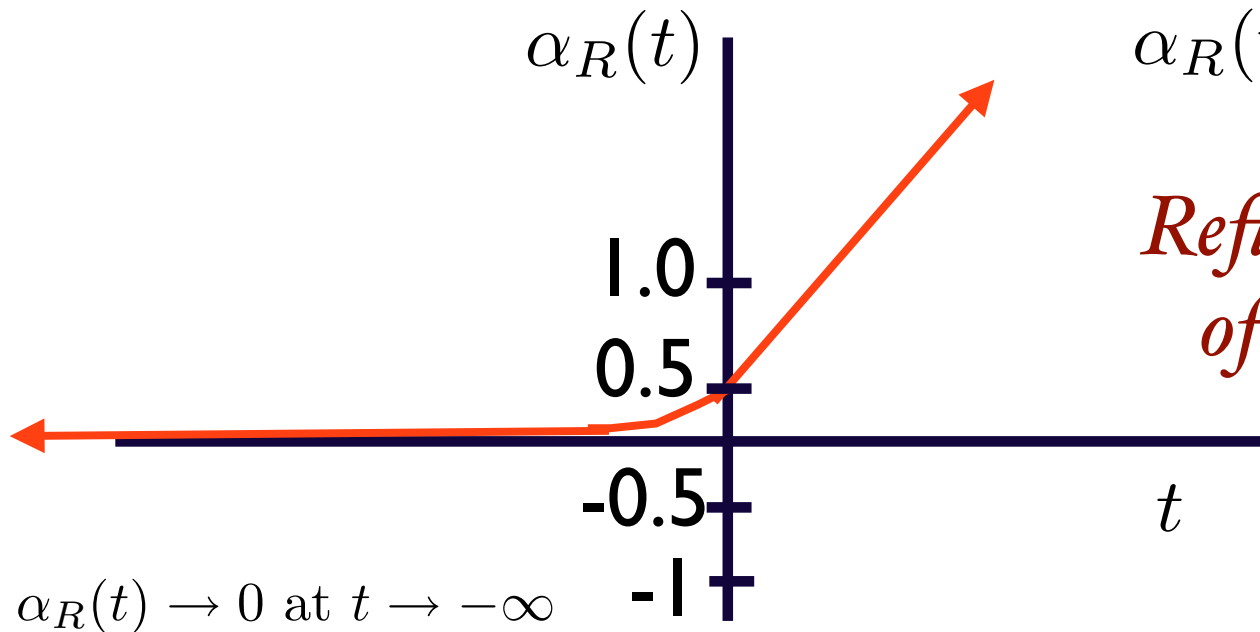
Reflects elementary coupling of two photons to quarks

$$\beta_R(t) \sim \frac{1}{t^2}$$

$$\frac{d\sigma}{dt} \sim \frac{1}{s^2} \frac{1}{t^4} \sim \frac{1}{s^6} \text{ at fixed } \frac{Q^2}{s}, \frac{t}{s}$$

Regge domain

$$T(\gamma^* p \rightarrow \pi^+ n) \sim \epsilon \cdot p_i \sum_R s_R^{\alpha_R(t)} \beta_R(t) \quad s \gg -t, Q^2$$



$$\alpha_R(t) \rightarrow 0 \text{ at } t \rightarrow -\infty$$

J=0 fixed pole

*Reflects elementary coupling
of two photons to quarks*

$$\beta_R(t) \sim \frac{1}{t^2}$$

$$\frac{d\sigma}{dt}(\gamma^* p \rightarrow \gamma p) \rightarrow \frac{1}{s^2} \beta_R^2(t) \sim \frac{1}{s^2 t^4} \sim \frac{1}{s^6} \text{ at fixed } \frac{t}{s}, \frac{Q^2}{s}$$

Fundamental test of QCD

J=0 Fixed pole in real and virtual Compton scattering

Damashek, Gilman
Close, Gunion, sjb
Llanes-Estrada,
Szczeponiak, sjb

Effective two-photon contact term

Seagull for scalar quarks

Real phase

$$M = s^0 \sum e_q^2 F_q(t)$$

Independent of Q^2 at fixed t

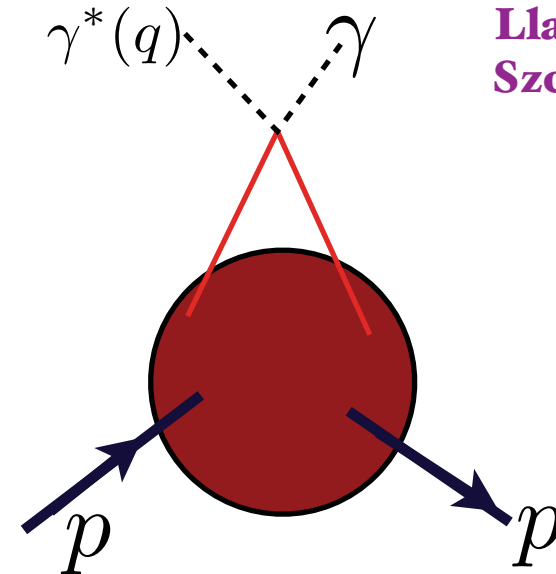
$\langle 1/x \rangle$ Moment: Related to Feynman-Hellman Theorem

Fundamental test of local gauge theory

No ambiguity in D-term

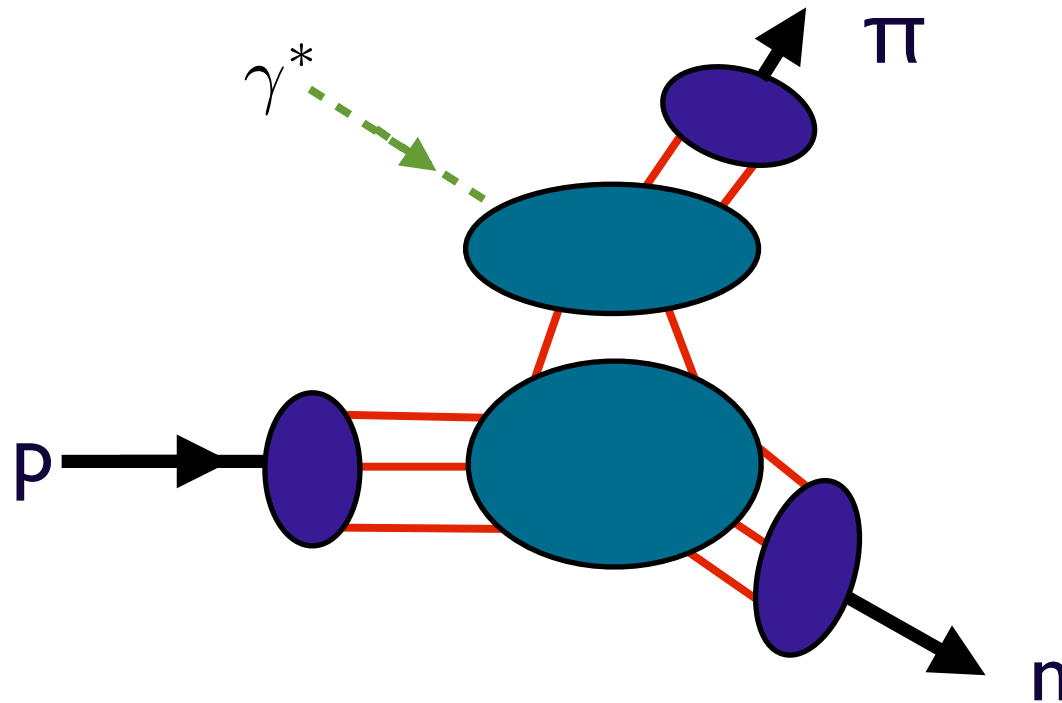
Q^2 -independent contribution to Real DVCS amplitude

$$s^2 \frac{d\sigma}{dt} (\gamma^* p \rightarrow \gamma p) = F^2(t)$$



Exclusive Electroproduction

$$ep \rightarrow e' \pi^+ n$$



*Hard Reggeon
Domain*

$$s \gg -t, Q^2 \gg \Lambda_{QCD}^2$$

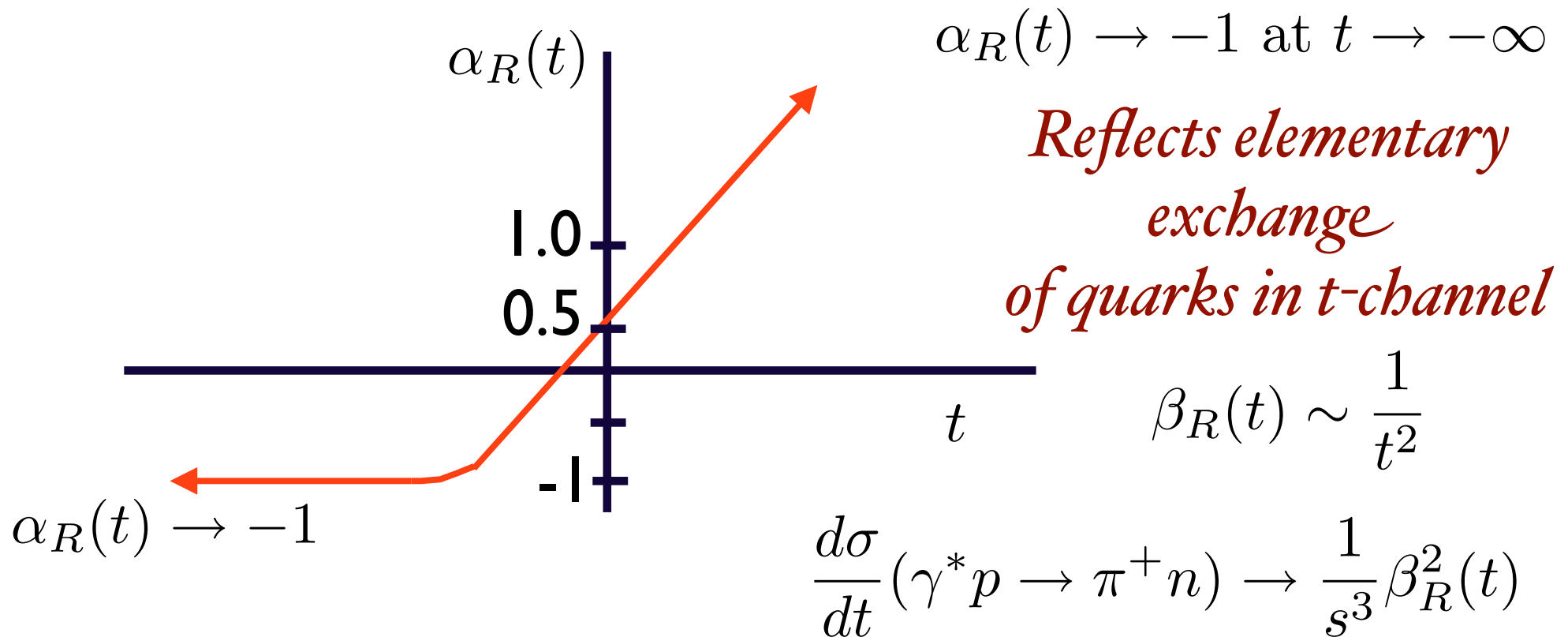
$$T(\gamma^* p \rightarrow \pi^+ n) \sim \epsilon \cdot p_i \sum_R s_R^\alpha(t) \beta_R(t)$$

$$\alpha_R(t) \rightarrow -1 \quad \text{Reflects elementary exchange of quarks in } t\text{-channel}$$

$$\beta_R(t) \sim \frac{1}{t^2} \quad \frac{d\sigma}{dt} \sim \frac{1}{s^7} \text{ at fixed } \frac{Q^2}{s}, \frac{t}{s}$$

Regge domain

$$T(\gamma^* p \rightarrow \pi^+ n) \sim \epsilon \cdot p_i \sum_R s_R^\alpha(t) \beta_R(t) \quad s \gg -t, Q^2$$



$$\frac{d\sigma}{dt} \sim \frac{1}{s^3} \frac{1}{t^4} \sim \frac{1}{s^7} \text{ at fixed } \frac{Q^2}{s}, \frac{t}{s}$$

Fundamental test of QCD

$$|p, S_z\rangle = \sum_{n=3} \Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i) |n; \vec{k}_{\perp i}, \lambda_i\rangle$$

sum over states with $n=3, 4, \dots$ constituents

The Light Front Fock State Wavefunctions

$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

are boost invariant; they are independent of the hadron's energy and momentum P^μ .

The light-cone momentum fraction

$$x_i = \frac{k_i^+}{p^+} = \frac{k_i^0 + k_i^z}{P^0 + P^z}$$

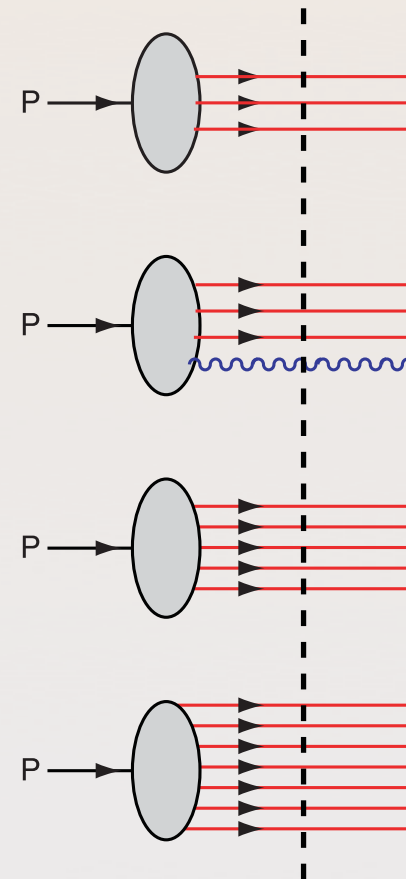
are boost invariant.

$$\sum_i^n k_i^+ = P^+, \quad \sum_i^n x_i = 1, \quad \sum_i^n \vec{k}_i^\perp = \vec{0}^\perp.$$

Intrinsic heavy quarks
 $c(x), b(x)$ at high x !

$$\bar{s}(x) \neq s(x)$$

$$\bar{u}(x) \neq \bar{d}(x)$$



Fixed LF time

Mueller: gluon Fock states \rightarrow BFKL Pomeron *Hidden Color*

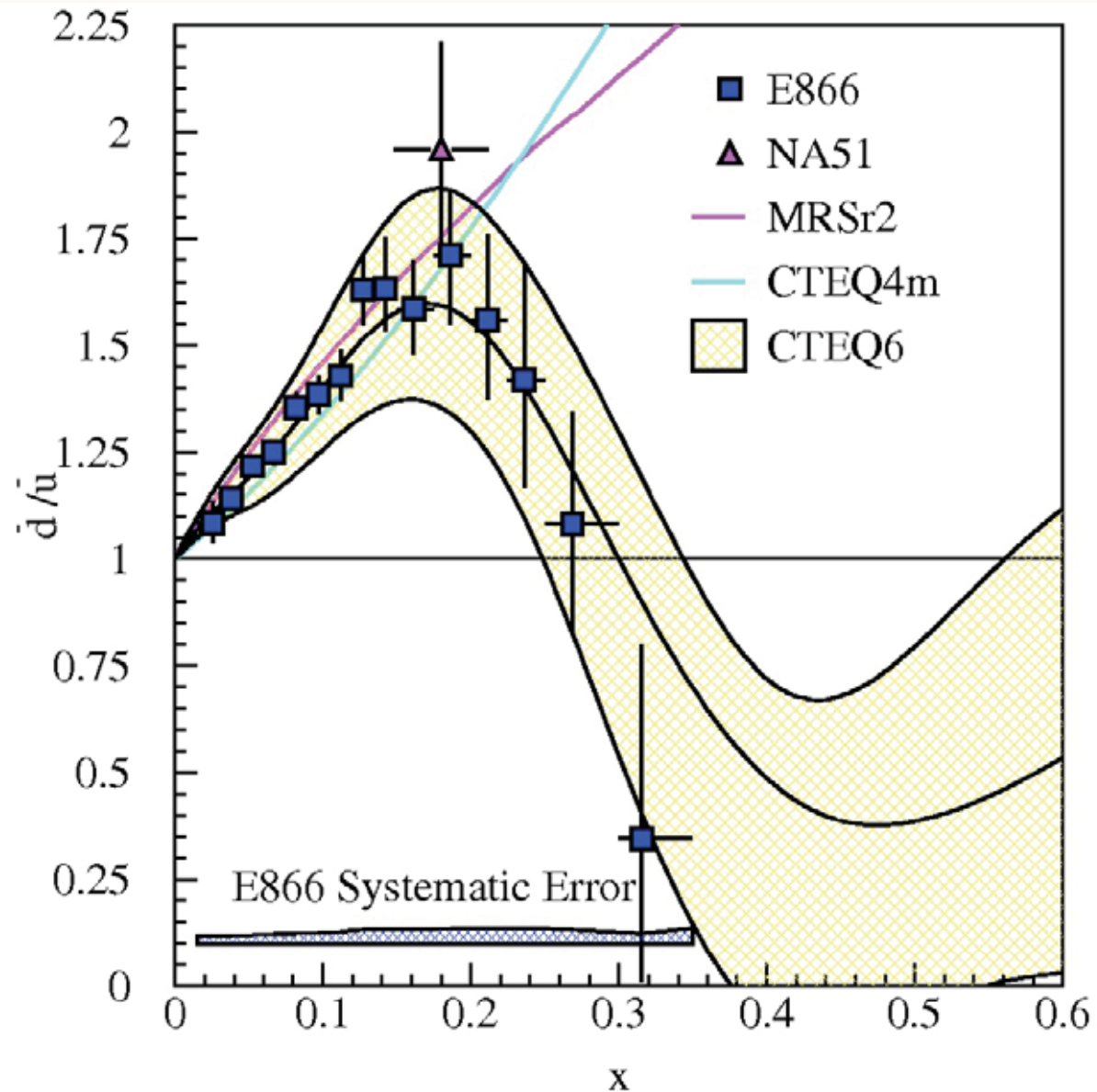
■ E866/NuSea (Drell-Yan)

$$\bar{d}(x) \neq \bar{u}(x)$$

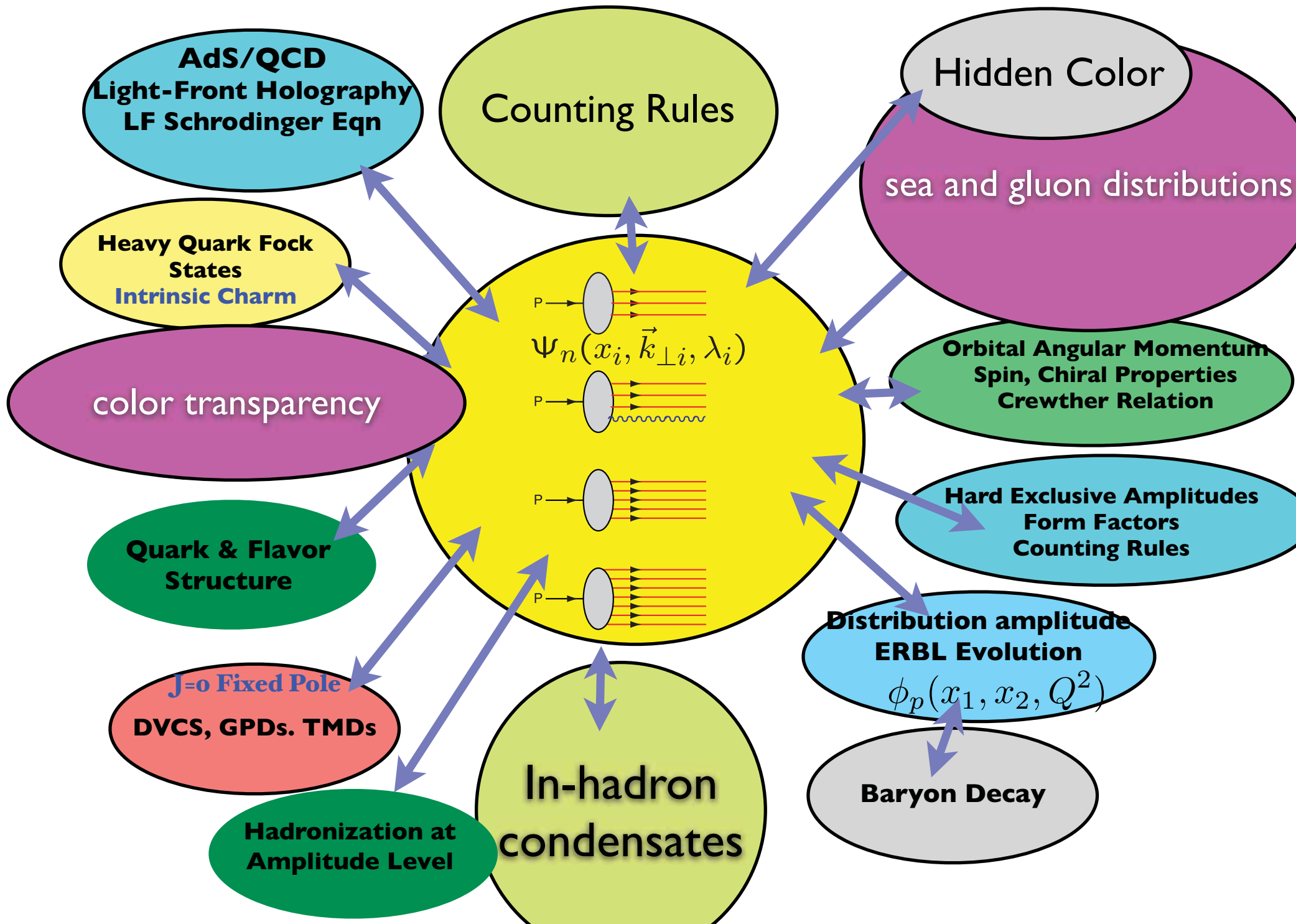
$$s(x) \neq \bar{s}(x)$$

*Intrinsic glue, sea,
heavy quarks*

$\bar{d}(x)/\bar{u}(x)$ for $0.015 \leq x \leq 0.35$



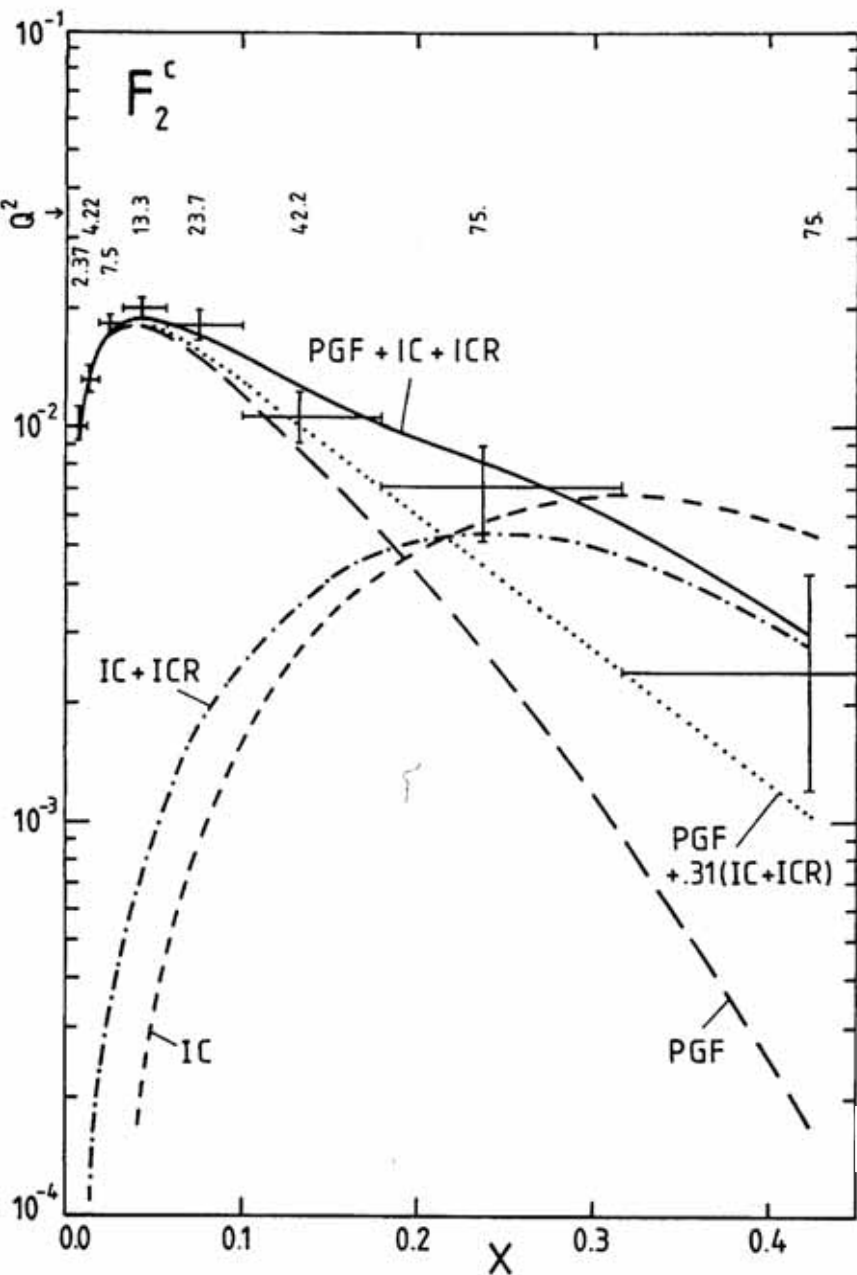
QCD and the LF Hadron Wavefunctions



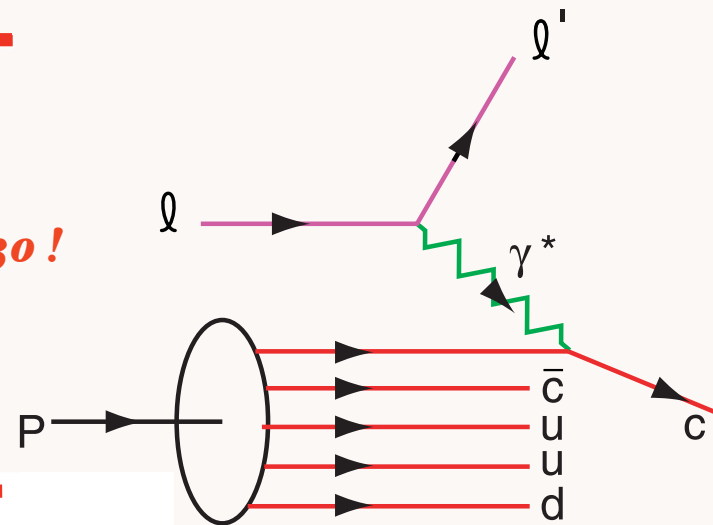
Measurement of Charm Structure Function

J. J. Aubert et al. [European Muon Collaboration], "Production Of Charmed Particles In 250-GeV Mu+ - Iron Interactions," Nucl. Phys. B 213, 31 (1983).

First Evidence for Intrinsic Charm
Never been checked!



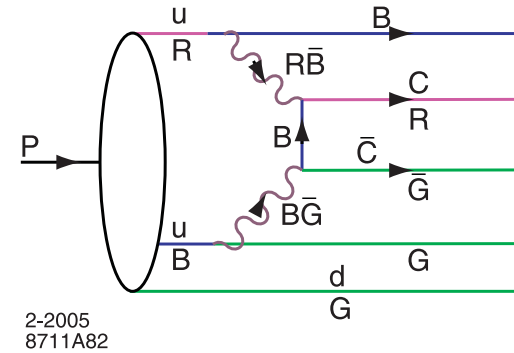
factor of 30!



DGLAP / Photon-Gluon Fusion: factor of 30 too small

Intrinsic Heavy-Quark Fock States

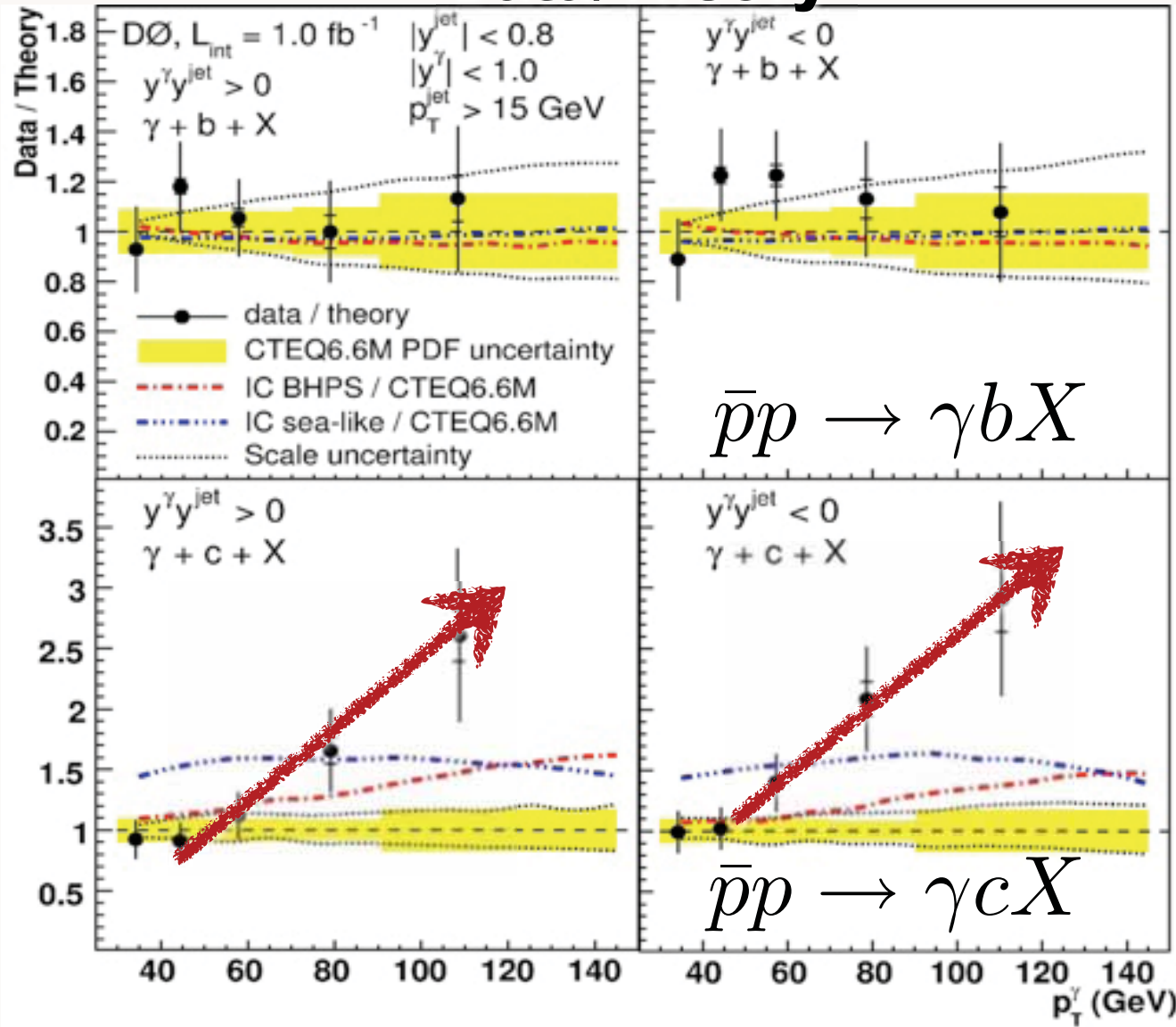
- Rigorous prediction of QCD, OPE
- Color-Octet Color-Octet Fock State!
- Probability $P_{Q\bar{Q}} \propto \frac{1}{M_Q^2}$ $P_{Q\bar{Q}Q\bar{Q}} \sim \alpha_s^2 P_{Q\bar{Q}}$ $P_{c\bar{c}/p} \simeq 1\%$
- Large Effect at high x
- Greatly increases kinematics of colliders such as Higgs production (Kopeliovich, Schmidt, Soffer, sjb)
- Severely underestimated in conventional parameterizations of heavy quark distributions (Pumplin, Tung)
- Many empirical tests



M. Polyakov et al.
OPE

Measurement of $\gamma + b + X$ and $\gamma + c + X$ Production Cross Sections
in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ TeV

Data/Theory



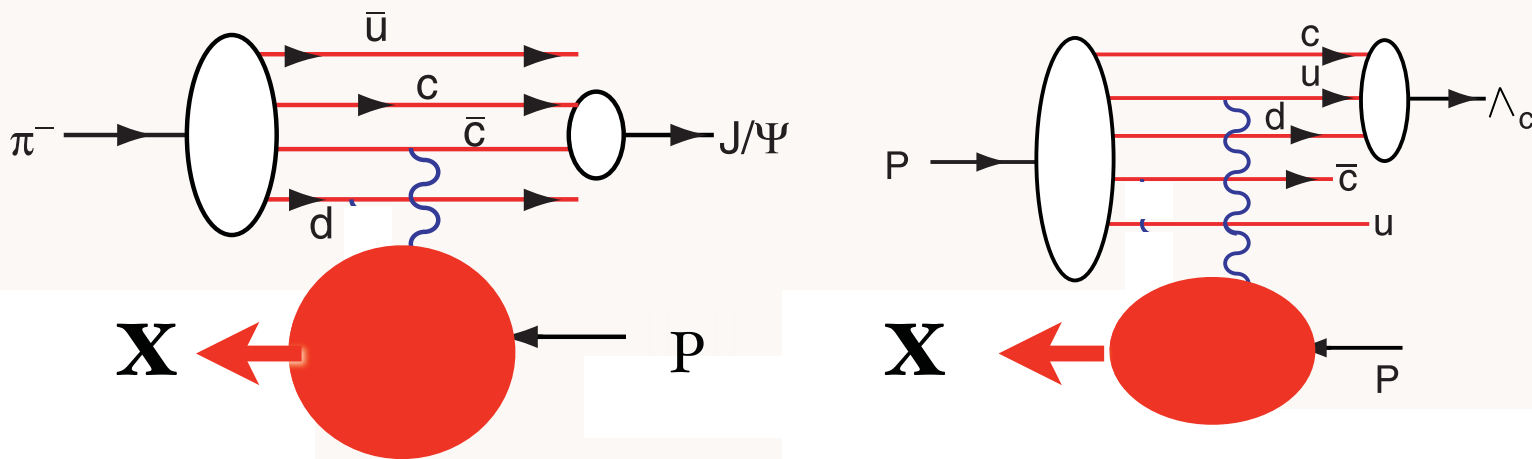
$$\frac{\Delta\sigma(\bar{p}p \rightarrow \gamma c X)}{\Delta\sigma(\bar{p}p \rightarrow \gamma b X)}$$

**Ratio
insensitive to
gluon PDF,
scales**

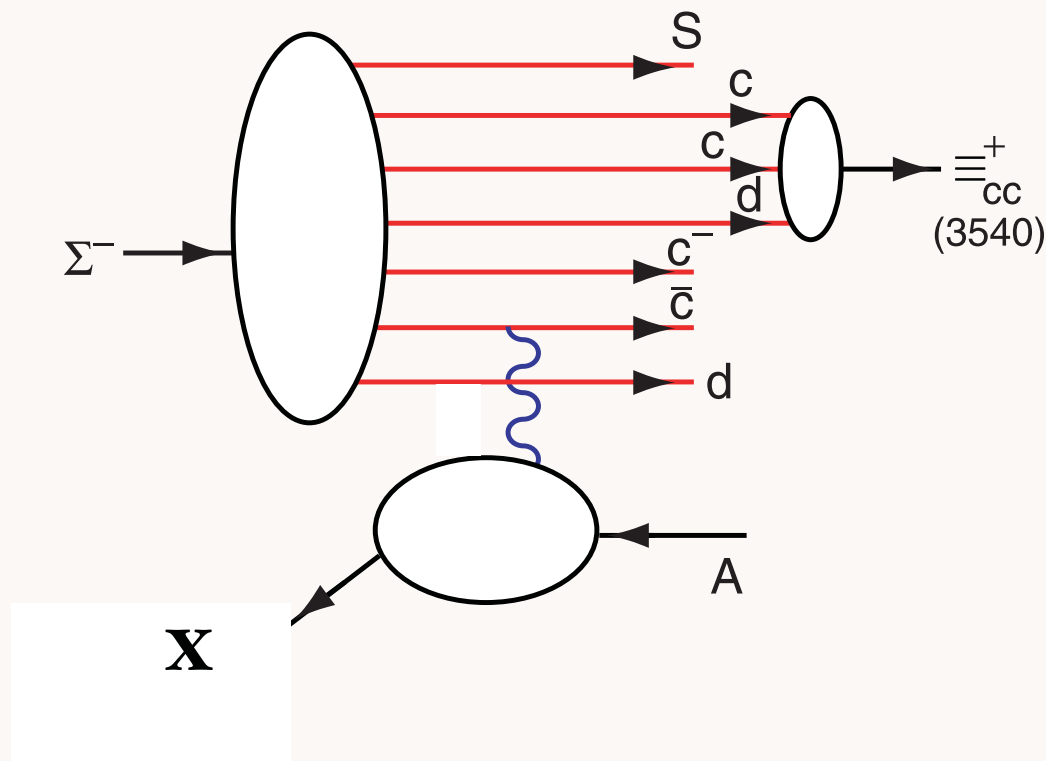
**Signal for
significant IC
at $x > 0.1$**

- EMC data: $c(x, Q^2) > 30 \times \text{DGLAP}$
 $Q^2 = 75 \text{ GeV}^2, x = 0.42$
- High x_F $pp \rightarrow J/\psi X$
- High x_F $pp \rightarrow J/\psi J/\psi X$
- High x_F $pp \rightarrow \Lambda_c X$ ISR
- High x_F $pp \rightarrow \Lambda_b X$ ISR
- High x_F $pp \rightarrow \Xi(ccd) X$ (SELEX)

Leading Hadron Production from Intrinsic Charm



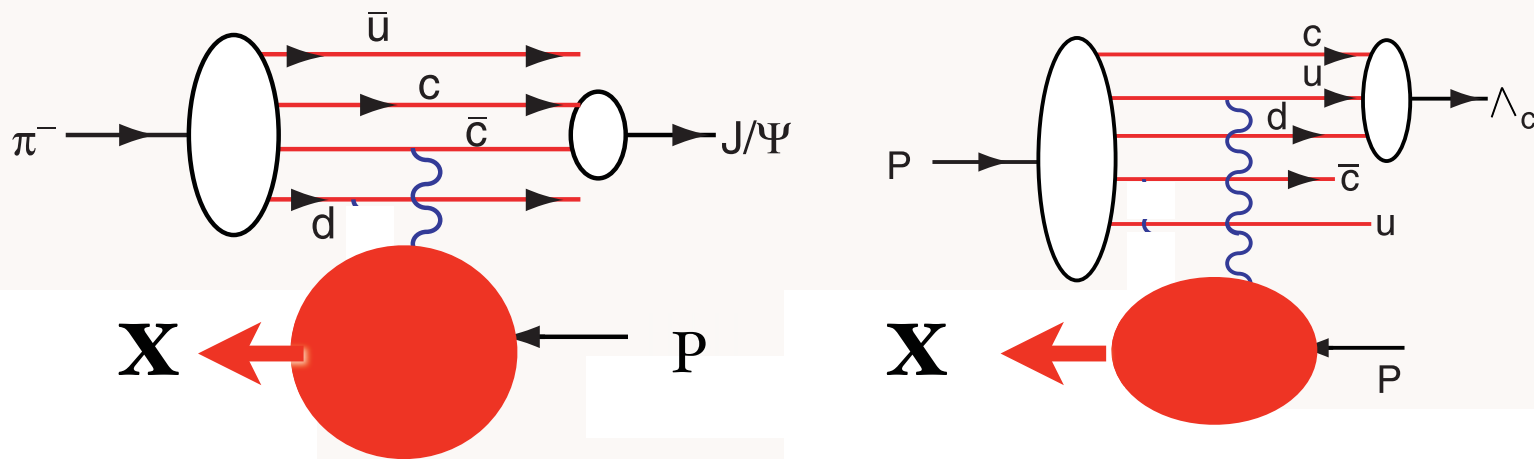
Coalescence of Comoving Charm and Valence Quarks
Produce J/ψ , Λ_c and other Charm Hadrons at High x_F



Production of a Double-Charm Baryon

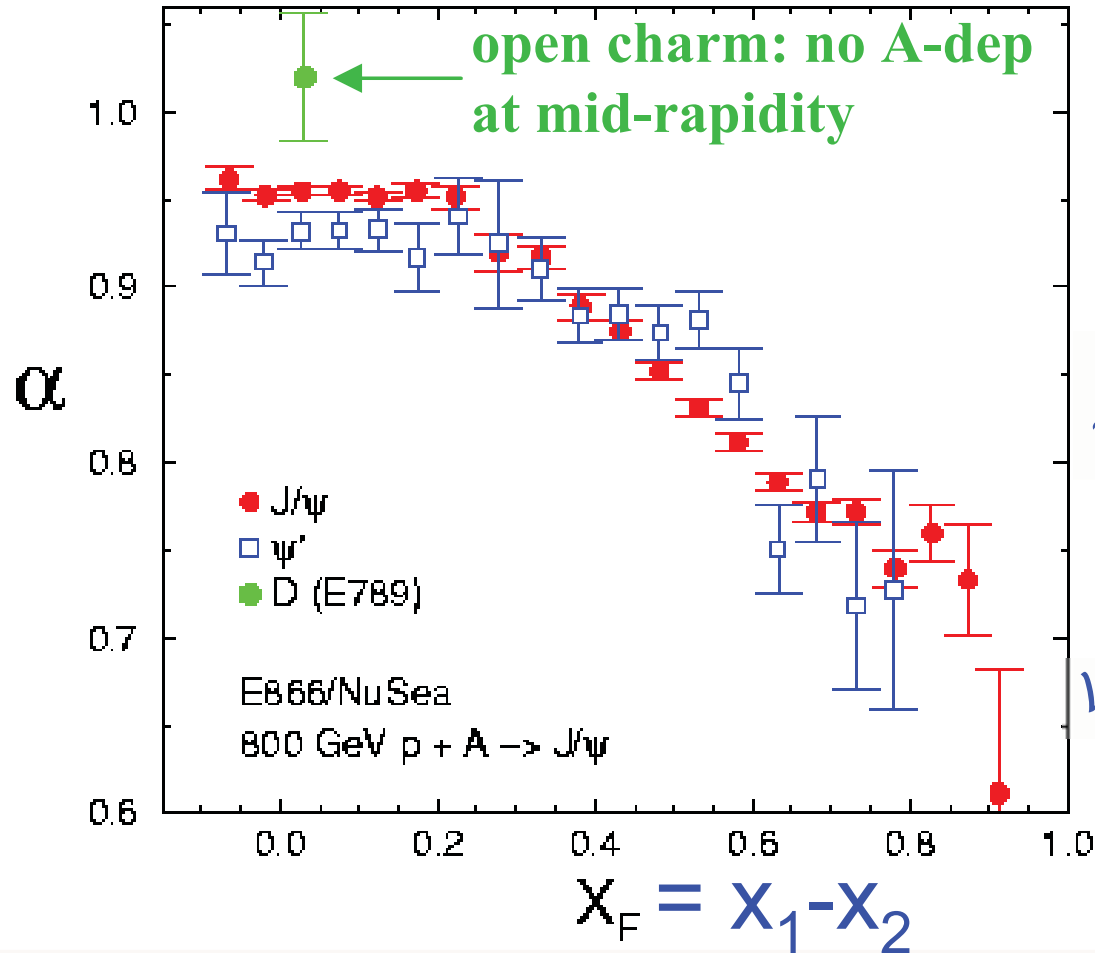
SELEX high x_F $\langle x_F \rangle = 0.33$

Leading Hadron Production from Intrinsic Charm



Coalescence of Comoving Charm and Valence Quarks
Produce J/ψ , Λ_c and other Charm Hadrons at High x_F

800 GeV p-A (FNAL) $\sigma_A = \sigma_p * A^\alpha$
PRL 84, 3256 (2000); PRL 72, 2542 (1994)



$$\frac{d\sigma}{dx_F} (pA \rightarrow J/\psi X)$$

Remarkably Strong Nuclear Dependence for Fast Charmonium

Violation of PQCD Factorization!

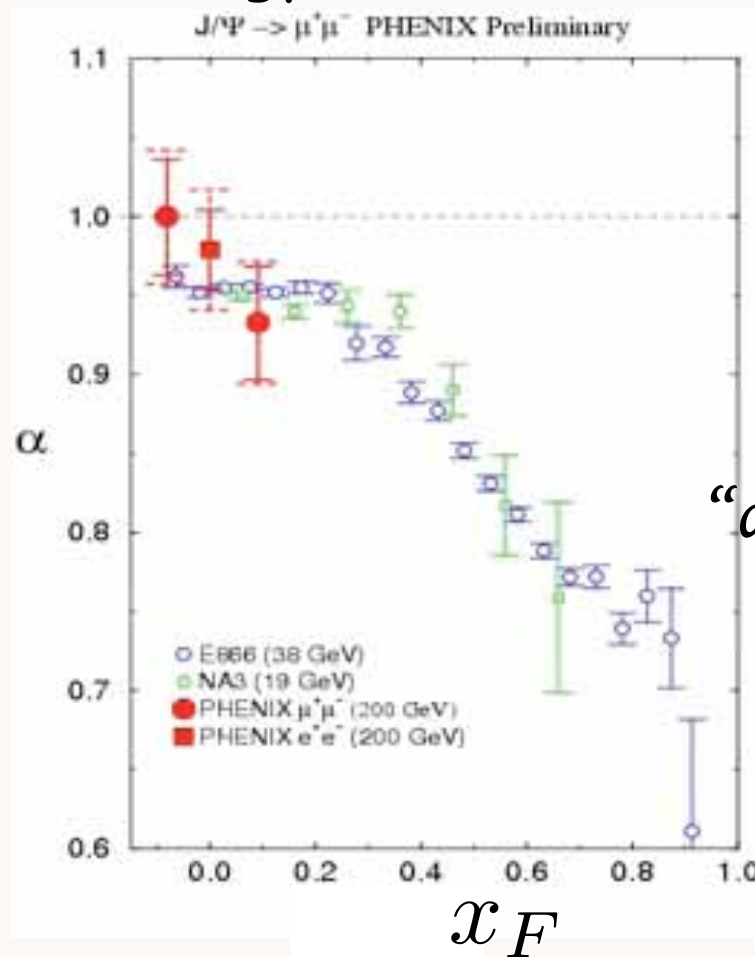
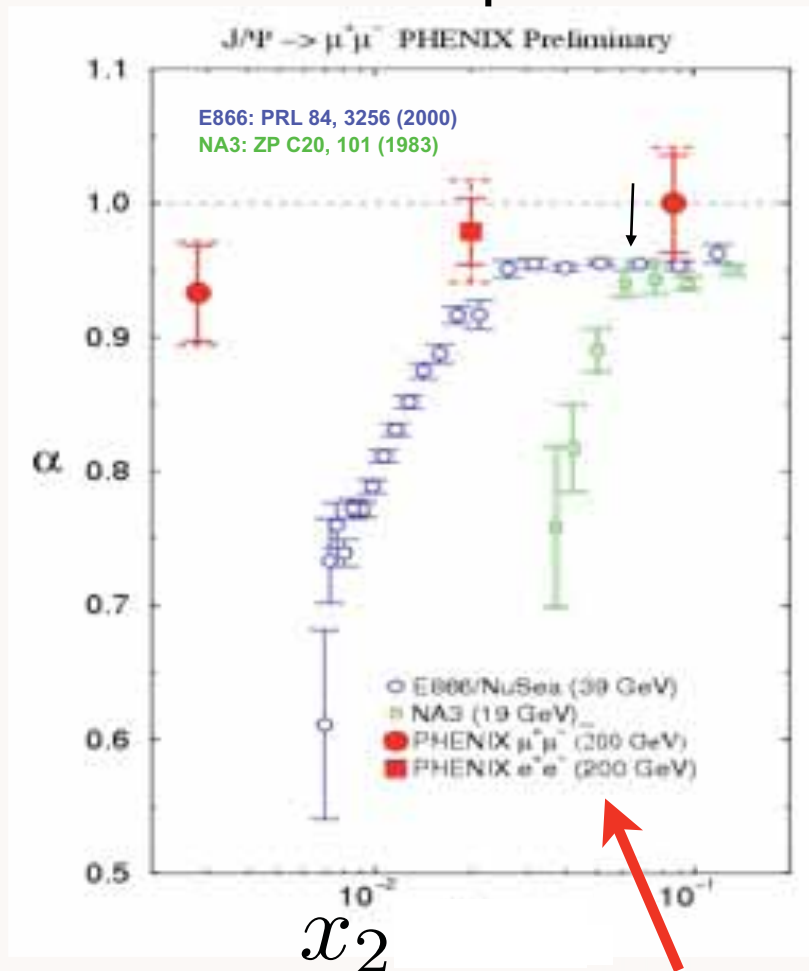
Violation of factorization in charm hadroproduction.

[P. Hoyer](#), [M. Vanttinen](#) ([Helsinki U.](#)), [U. Sukhatme](#) ([Illinois U., Chicago](#)) . HU-TFT-90-14, May 1990. 7pp.
 Published in Phys.Lett.B246:217-220,1990

J/ψ nuclear dependence vrs rapidity, x_{Au} , x_F

M.Leitch

PHENIX compared to lower energy measurements



Huge
"absorption"
effect

Klein, Vogt, PRL 91:142301, 2003
Kopeliovich, NP A696:669, 2001

*Violates PQCD
factorization!*

$$\frac{d\sigma}{dx_F}(pA \rightarrow J/\psi X)$$

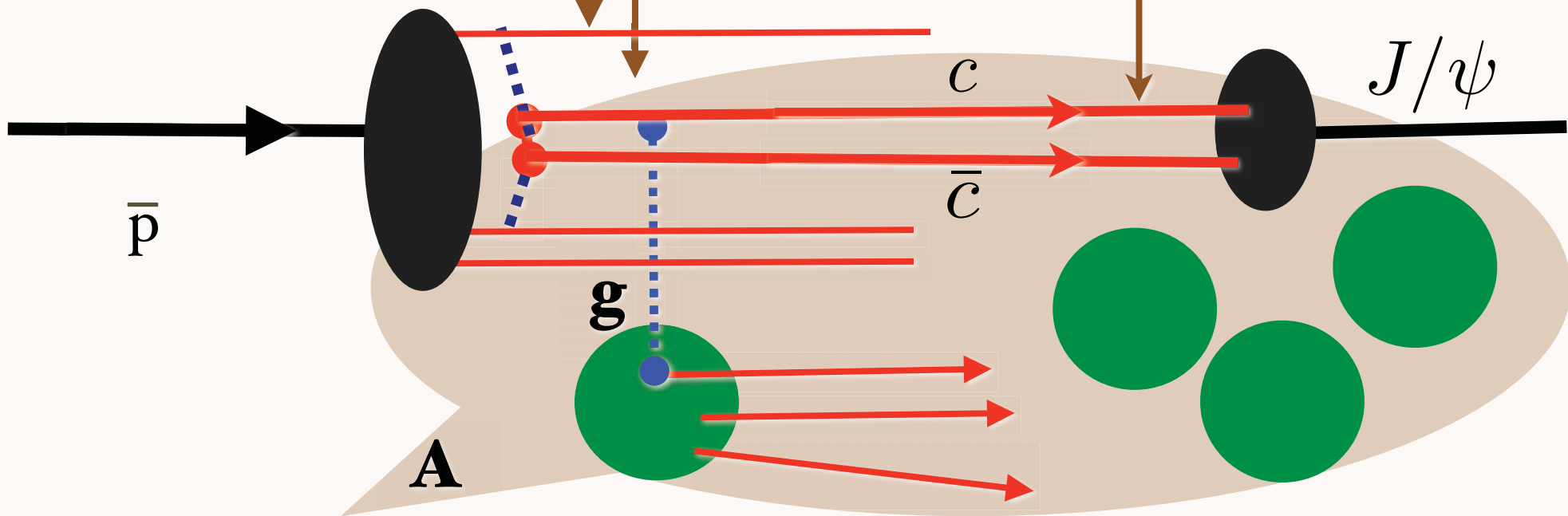
Hoyer, Sukhatme, Vanttinen

*Color-Opaque IC Fock state
interacts on nuclear front surface*

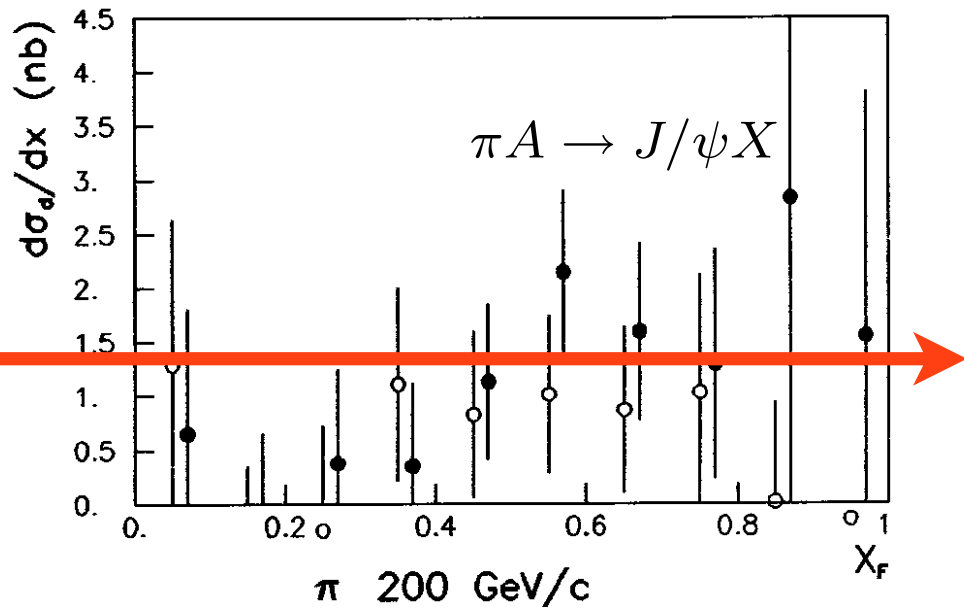
Scattering on front-face nucleon produces color-singlet $c\bar{c}$ pair

Octet-Octet IC Fock State

*No absorption of
small color-singlet*

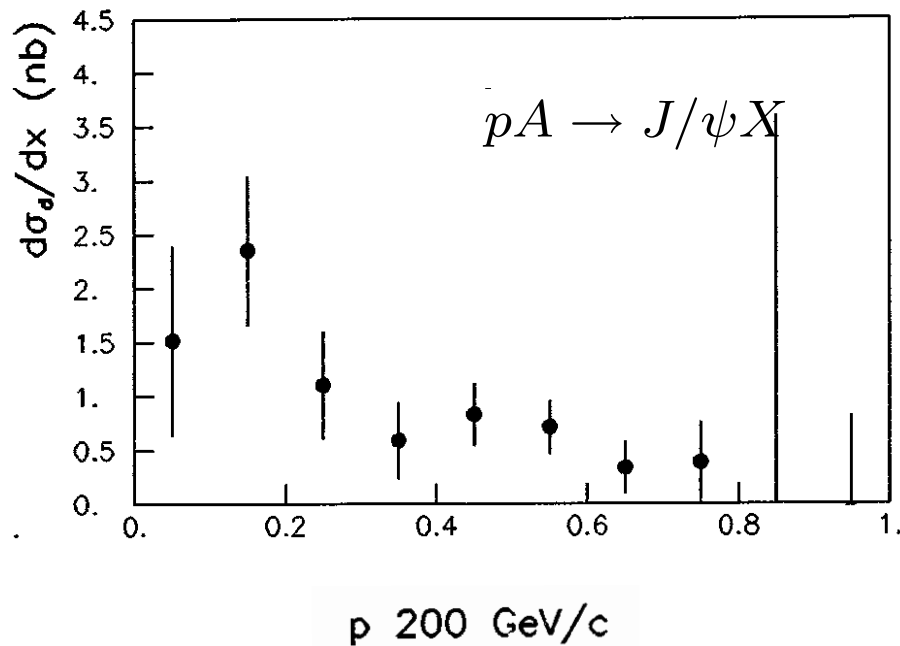


$$\frac{d\sigma}{dx_F}(pA \rightarrow J/\psi X) = A^{2/3} \times \frac{d\sigma}{dx_F}(pN \rightarrow J/\psi X)$$



$A^{2/3}$ component

J. Badier et al, NA3



$$\frac{d\sigma}{dx_F}(pA \rightarrow J/\psi X) = A^1 \frac{d\sigma_1}{dx_F} + A^{2/3} \frac{d\sigma_{2/3}}{dx_F}$$

**Excess beyond conventional PQCD
subprocesses**

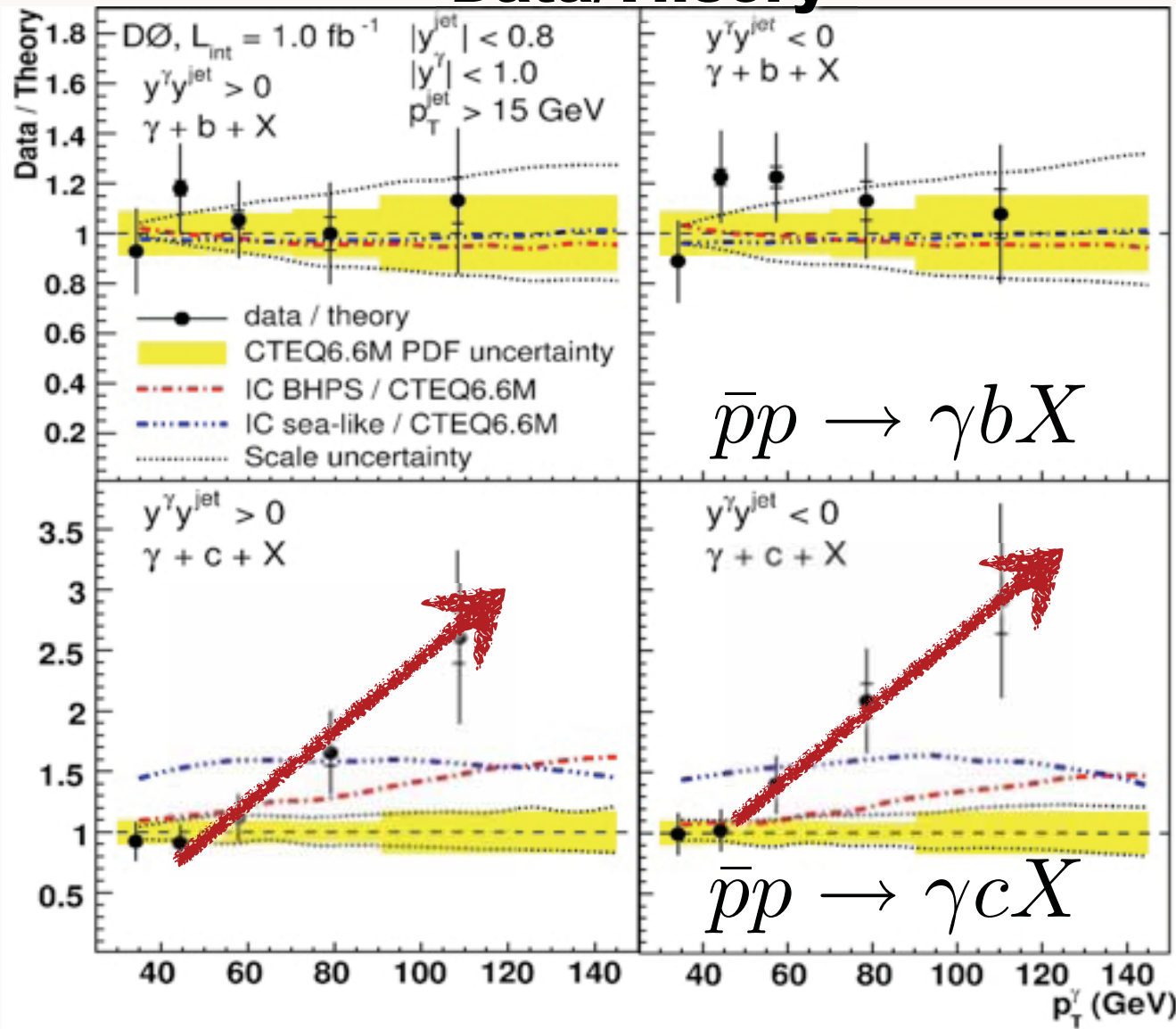
- IC Explains Anomalous $\alpha(x_F)$ not $\alpha(x_2)$ dependence of $pA \rightarrow J/\psi X$
(Mueller, Gunion, Tang, SJB)
- Color Octet IC Explains $A^{2/3}$ behavior at high x_F (NA3, Fermilab) *Color Opacity*
(Kopeliovitch, Schmidt, Soffer, SJB)
- IC Explains $J/\psi \rightarrow \rho\pi$ puzzle
(Karliner, SJB)
- IC leads to new effects in B decay
(Gardner, SJB)

Higgs production at $x_F = 0.8$!

Goldhaber, Kopeliovich,
Schmidt, Soffer, sjb

Measurement of $\gamma + b + X$ and $\gamma + c + X$ Production Cross Sections
in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ TeV

Data/Theory

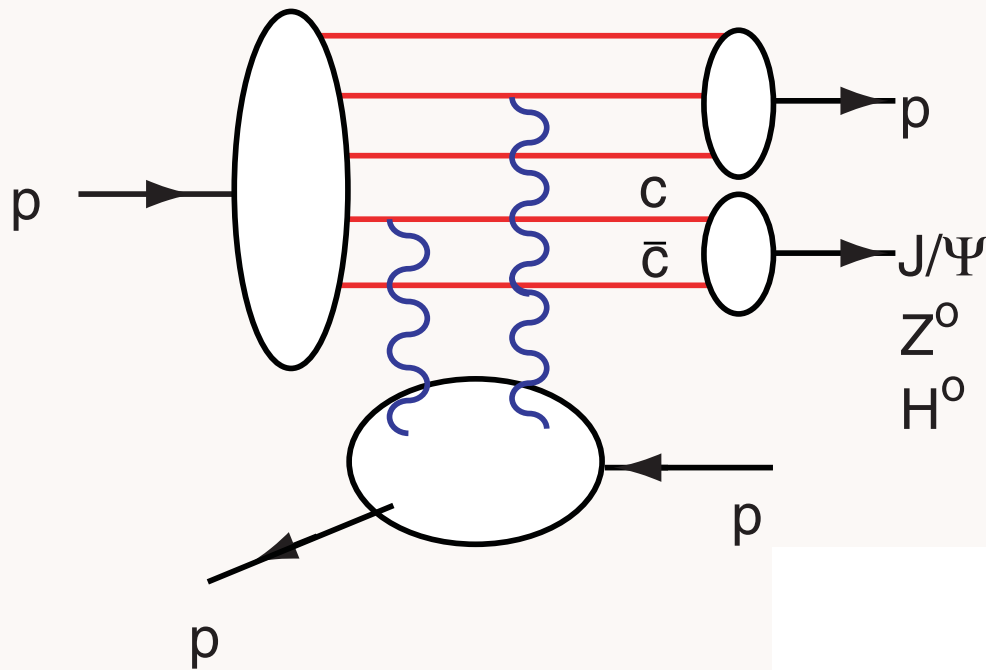


$$\frac{\Delta\sigma(\bar{p}p \rightarrow \gamma c X)}{\Delta\sigma(\bar{p}p \rightarrow \gamma b X)}$$

**Ratio
insensitive to
gluon PDF,
scales**

**Signal for
significant IC
at $x > 0.1$**

Intrinsic Charm Mechanism for Exclusive Diffraction Production



$$p p \rightarrow J/\psi p p$$

$$x_{J/\psi} = x_c + x_{\bar{c}}$$

Inclusive and Diffractive High- x_F Higgs Production!

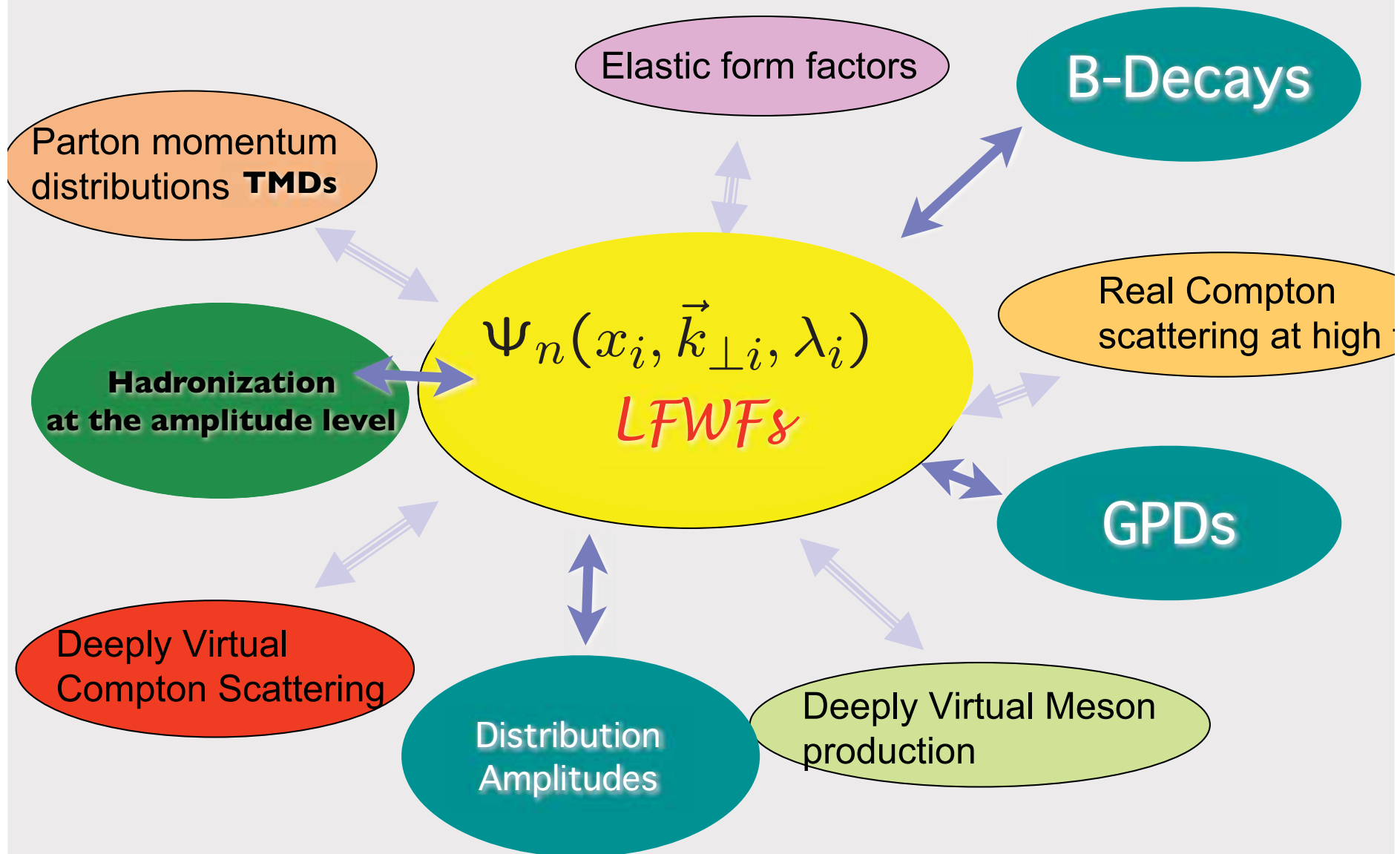
Kopeliovich, Schmidt,
Soffer, sjb

Intrinsic $c\bar{c}$ pair formed in color octet 8_C in proton wavefunction Large Color Dipole

Collision produces color-singlet J/ψ through color exchange

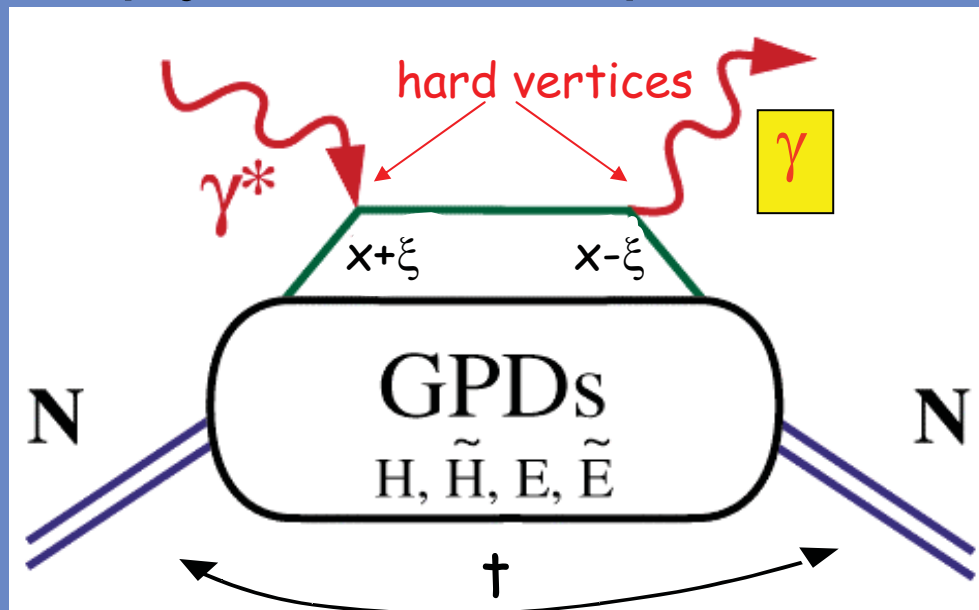
RHIC Experiment

A Unified Description of Hadron Structure



GPDs & Deeply Virtual Exclusive Processes - New Insight into Nucleon Structure

Deeply Virtual Compton Scattering (DVCS)



x - quark momentum fraction

ξ - longitudinal momentum transfer

$\sqrt{-t}$ - Fourier conjugate to transverse impact parameter

$H(x, \xi, t), E(x, \xi, t), \dots$ "Generalized Parton Distributions"

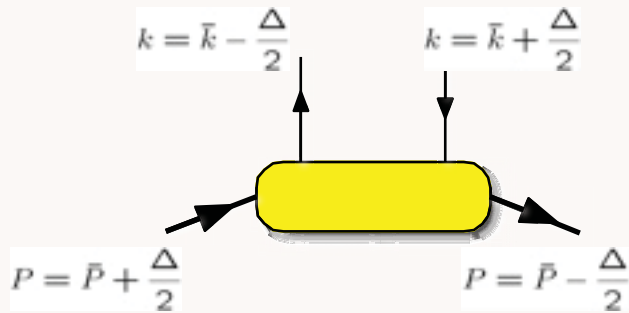
Timelike DVCS: **Mukherjee, Afanasev, Carlson, sjb**

Light-Front Wave Function Overlap Representation

DVCS/GPD

Diehl, Hwang, sjb, NPB596, 2001

See also: Diehl, Feldmann, Jakob, Kroll

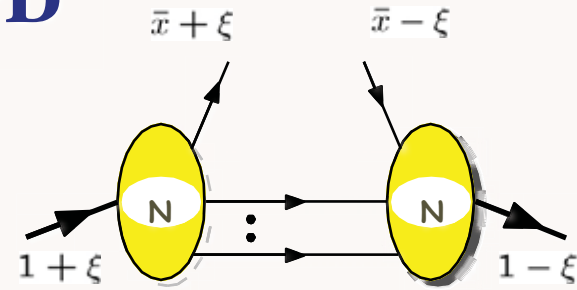


$\xi < \bar{x} < 1$

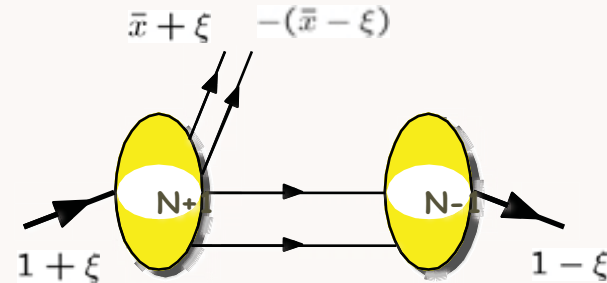
$-\xi < \bar{x} < \xi$

$-1 < \bar{x} < -\xi$

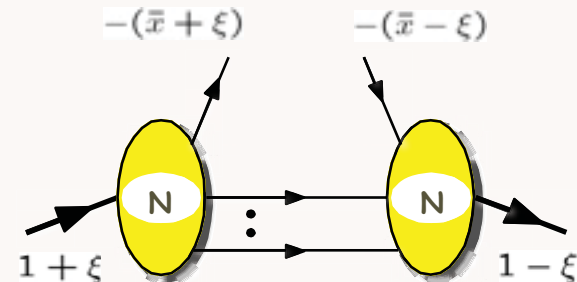
$$\sum_N$$



$$\sum_N$$



$$\sum_N$$



DGLAP
region

ERBL
region

DGLAP
region

Bakker & Ji
Lorce

Example of LFWF representation of GPDs ($n+1 \Rightarrow n-1$)

Diehl, Hwang, sjb

$$\begin{aligned}
 & \frac{1}{\sqrt{1-\zeta}} \frac{\Delta^1 - i\Delta^2}{2M} E_{(n+1 \rightarrow n-1)}(x, \zeta, t) \\
 &= (\sqrt{1-\zeta})^{3-n} \sum_{n, \lambda_i} \int \prod_{i=1}^{n+1} \frac{dx_i d^2\vec{k}_{\perp i}}{16\pi^3} 16\pi^3 \delta\left(1 - \sum_{j=1}^{n+1} x_j\right) \delta^{(2)}\left(\sum_{j=1}^{n+1} \vec{k}_{\perp j}\right) \\
 & \quad \times 16\pi^3 \delta(x_{n+1} + x_1 - \zeta) \delta^{(2)}(\vec{k}_{\perp n+1} + \vec{k}_{\perp 1} - \vec{\Delta}_{\perp}) \\
 & \quad \times \delta(x - x_1) \psi_{(n-1)}^{\uparrow*}(x'_i, \vec{k}'_{\perp i}, \lambda_i) \psi_{(n+1)}^{\downarrow}(x_i, \vec{k}_{\perp i}, \lambda_i) \delta_{\lambda_1 - \lambda_{n+1}},
 \end{aligned}$$

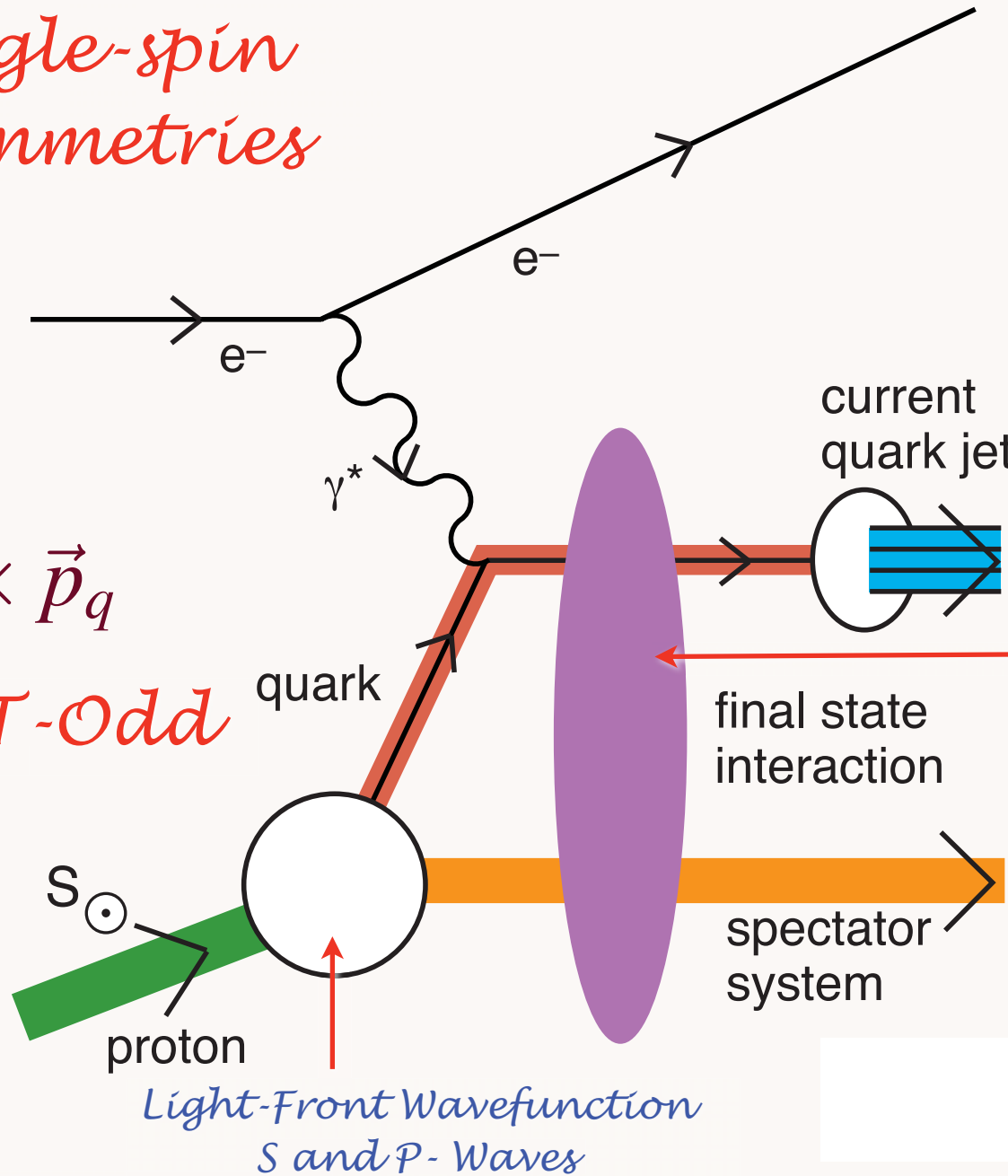
where $i = 2, \dots, n$ label the $n - 1$ spectator partons which appear in the final-state hadron wavefunction with

$$x'_i = \frac{x_i}{1-\zeta}, \quad \vec{k}'_{\perp i} = \vec{k}_{\perp i} + \frac{x_i}{1-\zeta} \vec{\Delta}_{\perp}.$$

Single-spin asymmetries

$$i \vec{S}_p \cdot \vec{q} \times \vec{p}_q$$

Pseudo-T-Odd



Leading Twist Sivers Effect

Hwang, Schmidt, sjb

Collins, **Burkardt**, Ji, Yuan

QCD S- and P-Coulomb Phases --Wilson Line

Leading-Twist Rescattering Violates pQCD Factorization!

Sign reversal in DY!

Final State Interactions Produce T-Odd (Sivers Effect)

- Bjorken Scaling!
- Arises from Interference of Final-State Coulomb Phases in S and P waves
- Relate to the quark contribution to the target proton anomalous magnetic moment
- Sum of Sivers Functions for all quarks and gluons vanishes. (Zero anomalous gavitomagnetic moment)

$$\vec{S} \cdot \vec{p}_{jet} \times \vec{q}$$

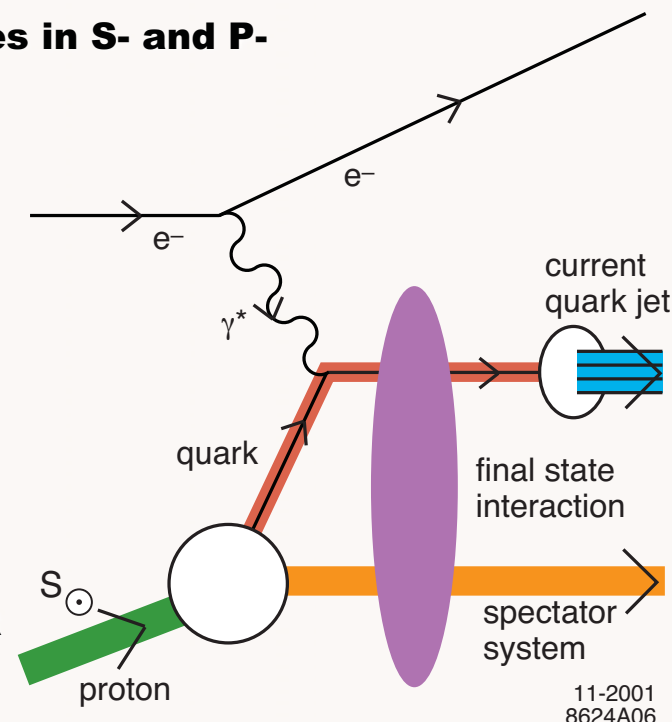
Hwang, Schmidt.
sjb; Burkardt

Final-State Interactions Produce Pseudo-T-Odd (Sivers Effect)

Hwang, Schmidt, sjb
Collins

- **Leading-Twist Bjorken Scaling!**
- **Requires nonzero orbital angular momentum of quark**
- **Arises from the interference of Final-State QCD Coulomb phases in S- and P-waves;**
- **Wilson line effect -- lc gauge prescription** *Stefanis*
- **Relate to the quark contribution to the target proton anomalous magnetic moment and final-state QCD phases**
- **QCD phase at soft scale!**
- **New window to QCD coupling and running gluon mass in the IR**
- **QED S and P Coulomb phases infinite -- difference of phases finite!**
- **Alternate: Retarded and Advanced Gauge: Augmented LFWFs**

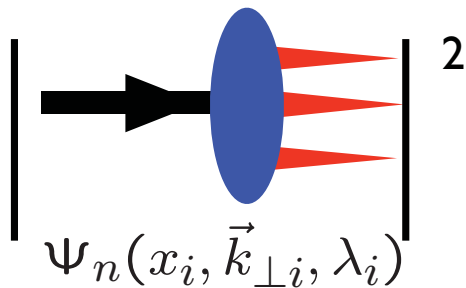
$$\mathbf{i} \vec{S} \cdot \vec{p}_{jet} \times \vec{q}$$



Pasquini, Xiao, Yuan, sjb
Mulders, Boer Qiu, Sterman

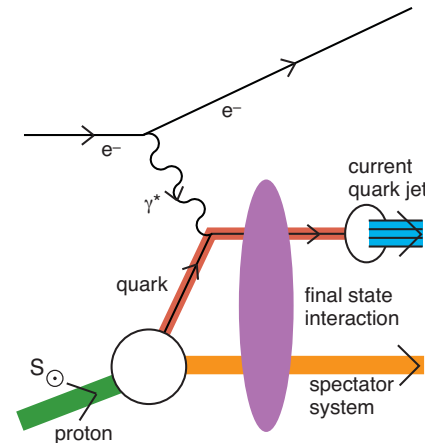
Static

- Square of Target LFWFs
- No Wilson Line
- Probability Distributions
- Process-Independent
- T-even Observables
- No Shadowing, Anti-Shadowing
- Sum Rules: Momentum and J^z
- DGLAP Evolution; mod. at large x
- No Diffractive DIS



Dynamic

- Modified by Rescattering: ISI & FSI
- Contains Wilson Line, Phases
- No Probabilistic Interpretation
- Process-Dependent - From Collision
- T-Odd (Sivers, Boer-Mulders, etc.)
- Shadowing, Anti-Shadowing, Saturation
- Sum Rules Not Proven
- DGLAP Evolution
- Hard Pomeron and Odderon Diffractive DIS



**Hwang,
Schmidt, sjb,
Mulders, Boer
Qiu, Sterman
Collins, Qiu
Pasquini, Xiao,
Yuan, sjb**

QCD Lagrangian

gluon dynamics quark kinetic energy + quark-gluon dynamics mass term

$$\mathcal{L}_{QCD} = -\frac{1}{4} \text{Tr}(G^{\mu\nu} G_{\mu\nu}) + \sum_{f=1}^{n_f} i \bar{\Psi}_f D_\mu \gamma^\mu \Psi_f + \sum_{f=1}^{n_f} m_f \bar{\Psi}_f \Psi_f$$
$$iD^\mu = i\partial^\mu - gA^\mu \qquad [D^\mu, D^\nu] = igG^{\mu\nu}$$

$\lim N_C \rightarrow 0$ at fixed $\alpha = C_F \alpha_s, n_\ell = n_F / C_F$

Analytic limit of QCD: Abelian Gauge Theory

QCD  **QED**

P. Huet, sjb

QED ($N_c=0$): Underlies Atomic Physics, Molecular Physics, Chemistry, Electromagnetic Interactions ...

QCD: Underlies Hadron Physics, Nuclear Physics, Strong Interactions, Jets

Theoretical Tools

- Feynman diagrams and perturbation theory
- Bethe Salpeter Equation, Dyson-Schwinger Equations
- Lattice Gauge Theory, Hägler, Lepage
- Light-Front Methods: Discretized Light-Front Quantization, Transverse Lattice
- AdS/CFT !

LF Quantization

Bjorken, Kogut, Soper, Susskind

LFWFs and Exclusive QCD:

Lepage and SJB, Efremov, Radyushkin

RGE and LF Hamiltonians:

Glazek & Wilson

DLCQ:

Hornbostel, Pauli, & SJB

Pinsky, Hiller

Renormalization of H_{LF}

Hiller, Chabysheva, Pauli, Pinsky, McCartor, Suaya, sjb

Rotation Invariance, Regularization

Karmanov, Mathiot

Zero-Modes: Standard Model

Srivastava, sjb

Light-Front formalism links dynamics to spectroscopy

Physical gauge: $A^+ = 0$

$$L^{QCD} \rightarrow H_{LF}^{QCD}$$

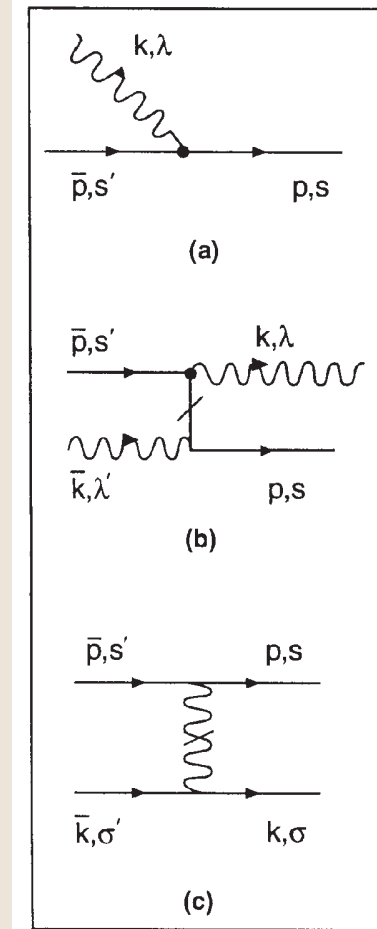
Heisenberg Matrix Formulation

$$H_{LF}^{QCD} = \sum_i \left[\frac{m^2 + k_{\perp}^2}{x} \right]_i + H_{LF}^{int}$$

H_{LF}^{int} : Matrix in Fock Space

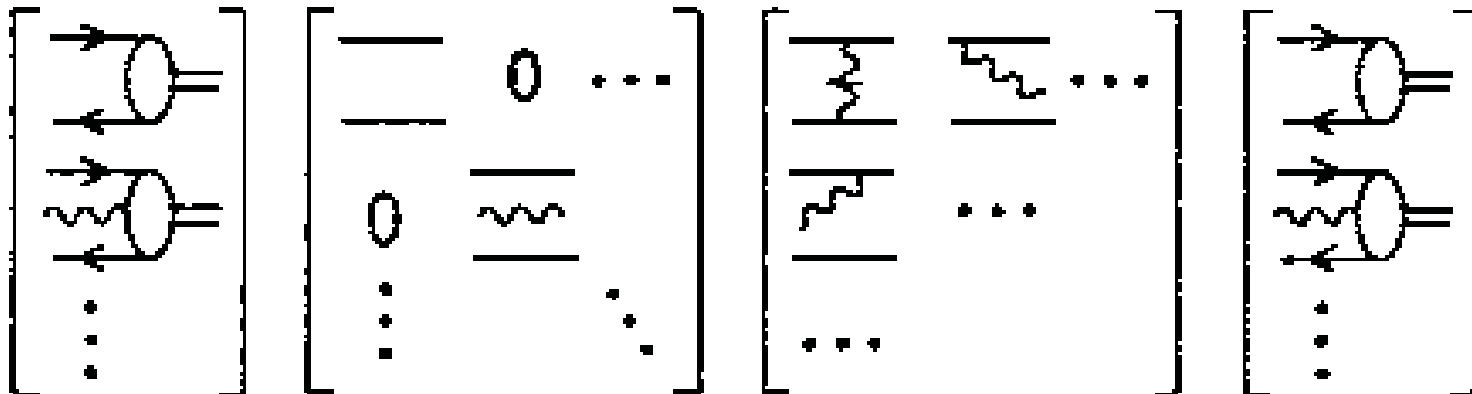
$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

Eigenvalues and Eigensolutions give Hadron Spectrum and Light-Front wavefunctions



LIGHT-FRONT SCHRÖDINGER EQUATION

$$\left(M_\pi^2 - \sum_i \frac{\vec{k}_{\perp i}^2 + m_i^2}{x_i} \right) \begin{bmatrix} \psi_{q\bar{q}/\pi} \\ \psi_{q\bar{q}g/\pi} \\ \vdots \end{bmatrix} = \begin{bmatrix} \langle q\bar{q} | V | q\bar{q} \rangle & \langle q\bar{q} | V | q\bar{q}g \rangle & \cdots \\ \langle q\bar{q}g | V | q\bar{q} \rangle & \langle q\bar{q}g | V | q\bar{q}g \rangle & \cdots \\ \vdots & \vdots & \ddots \end{bmatrix} \begin{bmatrix} \psi_{q\bar{q}/\pi} \\ \psi_{q\bar{q}g/\pi} \\ \vdots \end{bmatrix}$$



$$A^+ = 0$$

G.P. Lepage, sjb

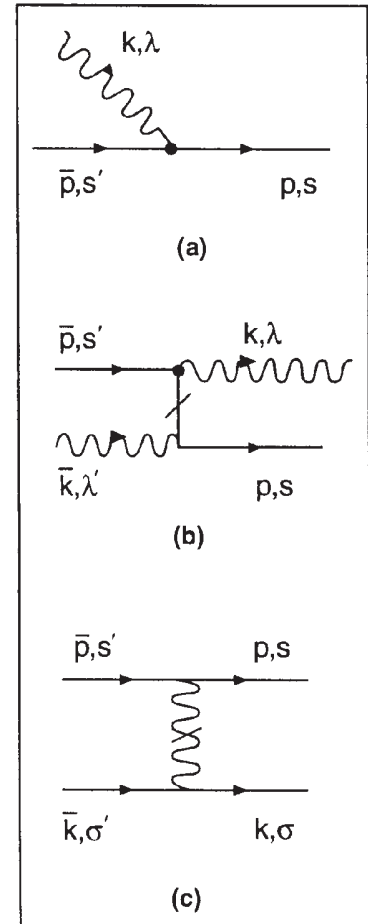
Light-Front QCD

$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

H.C. Pauli & sjb

Heisenberg Matrix Formulation

Discretized Light-Cone Quantization

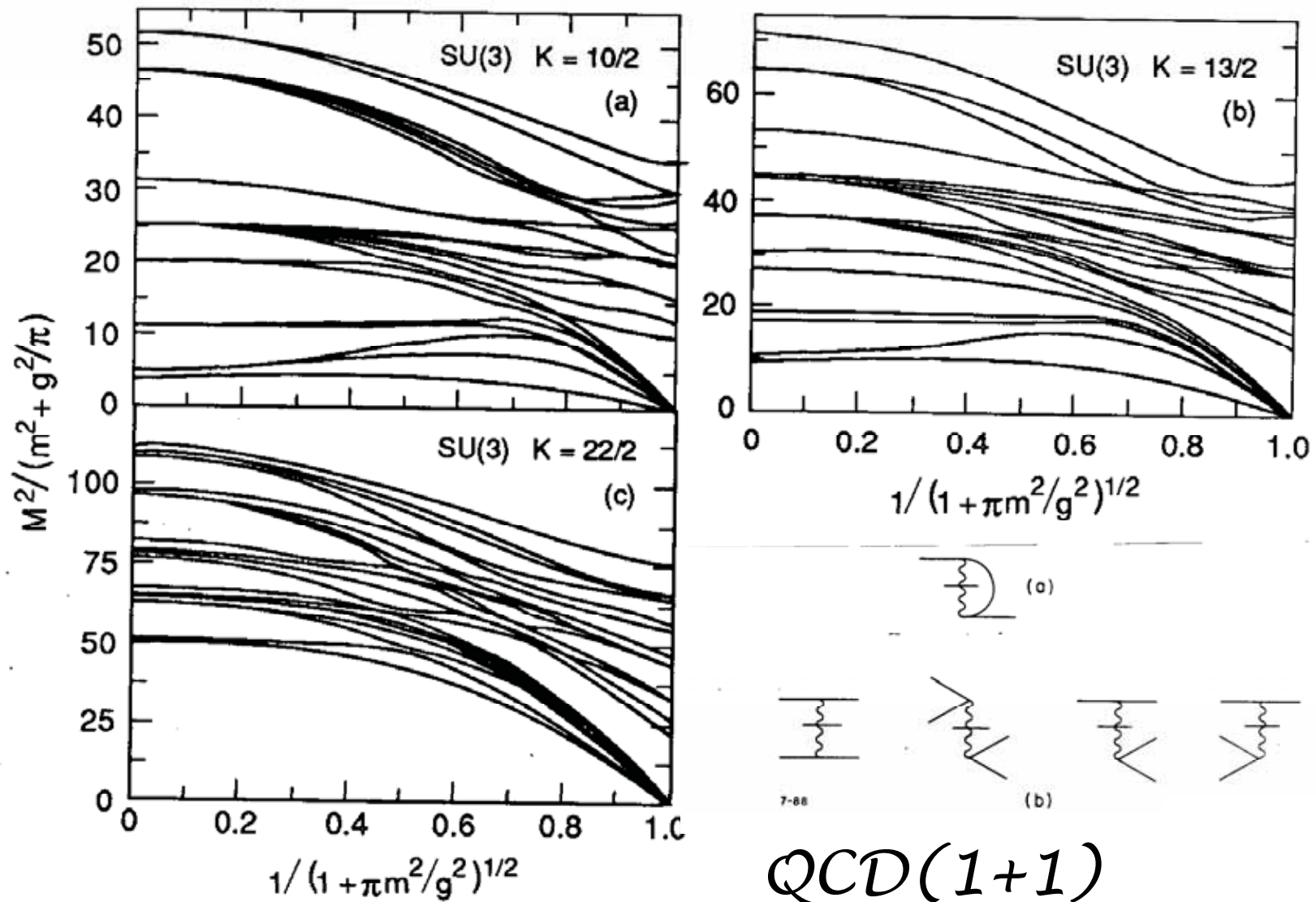


n	Sector	1 q \bar{q}	2 gg	3 q \bar{q} g	4 q \bar{q} q \bar{q}	5 ggg	6 q \bar{q} gg	7 q \bar{q} q \bar{q} g	8 q \bar{q} q \bar{q} q \bar{q}	9 gggg	10 q \bar{q} ggg	11 q \bar{q} q \bar{q} gg	12 q \bar{q} q \bar{q} q \bar{q} g	13 q \bar{q} q \bar{q} q \bar{q} q \bar{q}
1	q \bar{q}				
2	gg			
3	q \bar{q} g							
4	q \bar{q} q \bar{q}	
5	ggg
6	q \bar{q} gg								.				.	.
7	q \bar{q} q \bar{q} g
8	q \bar{q} q \bar{q} q \bar{q}				
9	gggg
10	q \bar{q} ggg
11	q \bar{q} q \bar{q} gg
12	q \bar{q} q \bar{q} q \bar{q} g				
13	q \bar{q} q \bar{q} q \bar{q} q \bar{q}				

Eigenvalues and Eigensolutions give Hadron Spectrum and Light-Front wavefunctions

DLCQ: Frame-independent, No fermion doubling; Minkowski Space

DLCQ: Periodic BC in x^- . Discrete k^+ ; frame-independent truncation



Spectra for $N = 3$, baryon number $B = 0, 1$ and 2 as a function of g/m ; K fixed.

Light Cone Quantized QCD in (1+1)-Dimensions.

[Kent Hornbostel](#), [Stanley J. Brodsky](#), (SLAC), [Hans Christian Pauli](#), (Heidelberg, Max Planck Inst.) . SLAC-PUB-4678, Phys.Rev.D41:3814,1990.



ELSEVIER

PHYSICS REPORTS

Physics Reports 301 (1998) 299–486

Quantum chromodynamics and other field theories on the light cone

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^c *Ohio State University, Columbus, OH 43210, USA*

QUANTUM CHROMODYNAMICS AND OTHER FIELD THEORIES ON THE LIGHT CONE

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Goal: an analytic first approximation to QCD

- **As Simple as Schrödinger Theory in Atomic Physics**
- **Relativistic, Frame-Independent, Color-Confining**
- **QCD Coupling at all scales**
- **Hadron Spectroscopy**
- **Light-Front Wavefunctions**
- **Form Factors, Hadronic Observables, Constituent Counting Rules**
- **Insight into QCD Condensates**
- **Systematically improvable**

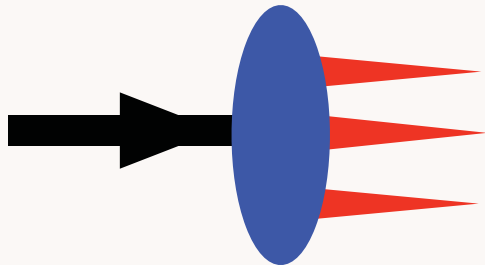
de Teramond, Deur, Shrock, Roberts, Tandy

Light-Front Holography and Non-Perturbative QCD

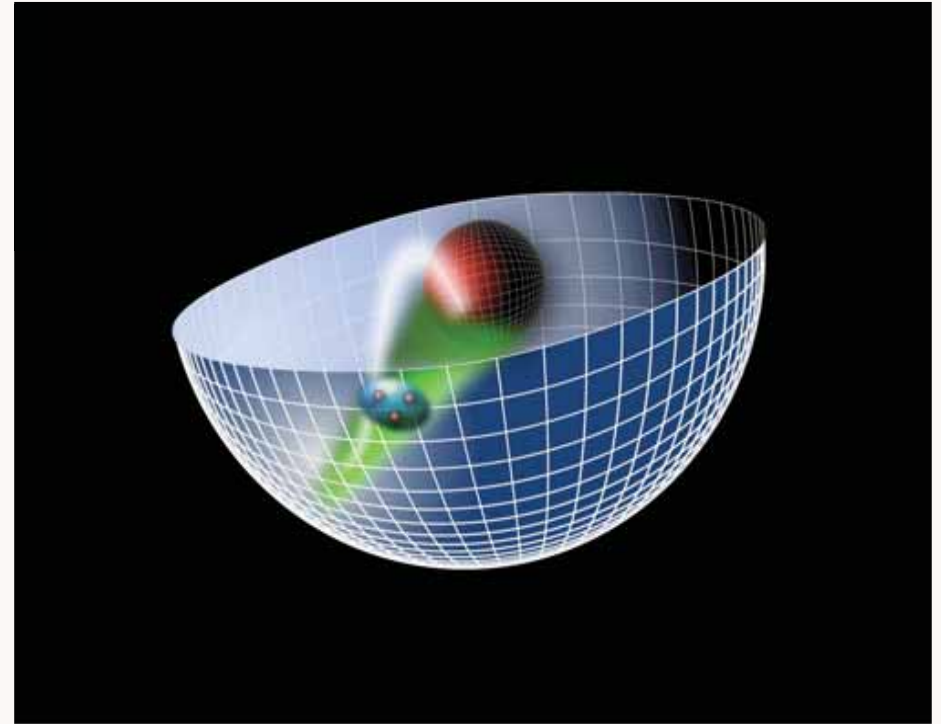
Goal:

**Use AdS/QCD duality to construct
a first approximation to QCD**

*Hadron Spectrum
Light-Front Wavefunctions,
Running coupling in IR*



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$



**in collaboration with
Guy de Teramond and Alexandre Deur**

Central problem for strongly-coupled gauge theories

Light-Front Wavefunctions

Dirac's Front Form: Fixed $\tau = t + z/c$

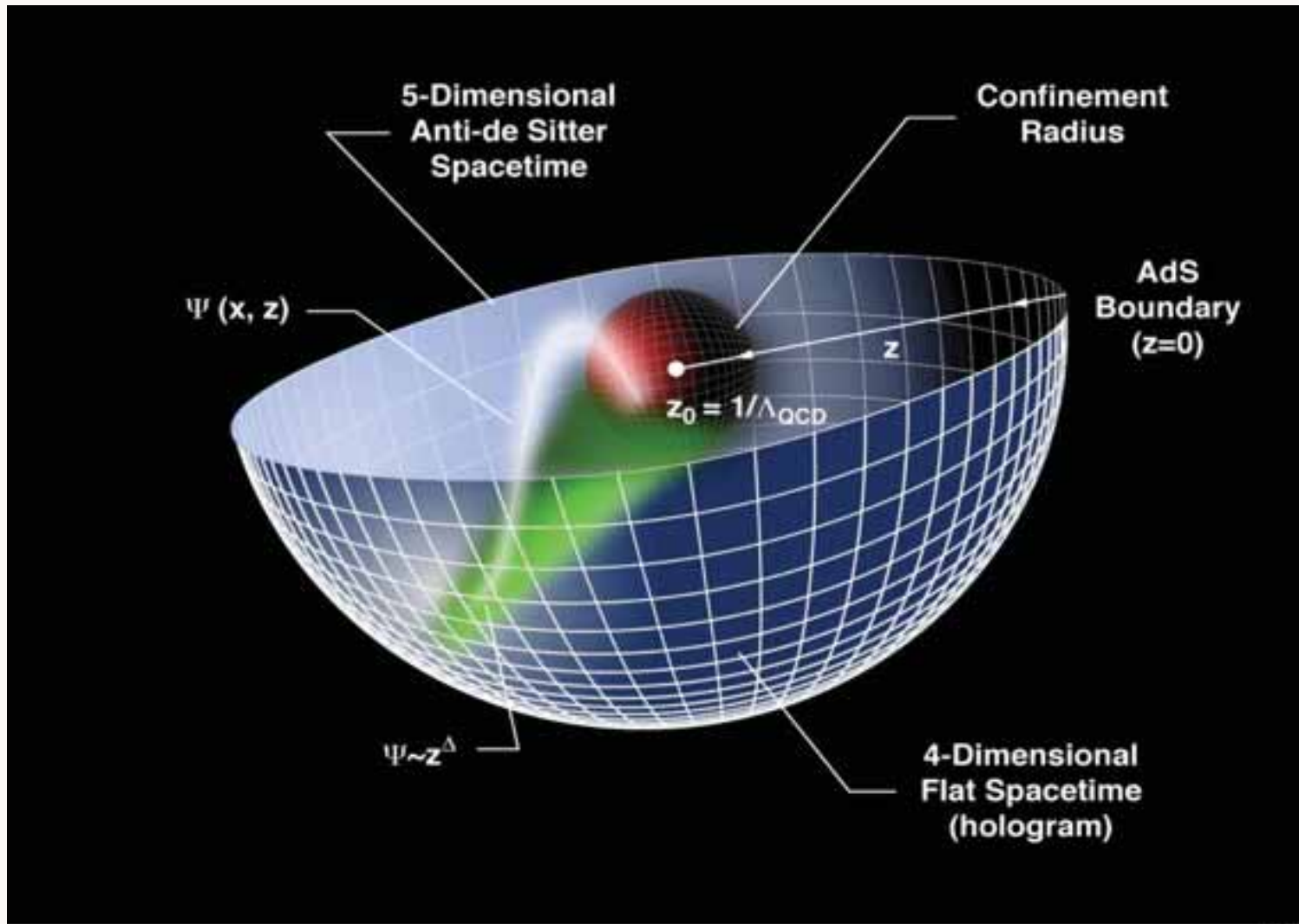
$$\Psi(x, k_{\perp}) \quad x_i = \frac{k_i^+}{P^+}$$

Invariant under boosts. Independent of P^{μ}

$$H_{LF}^{QCD} |\psi\rangle = M^2 |\psi\rangle$$

Remarkable new insights from AdS/CFT, the duality between conformal field theory and Anti-de Sitter Space

Applications of AdS/CFT to QCD



Changes in physical length scale mapped to evolution in the 5th dimension z

in collaboration with Guy de Teramond

Conformal Theories are invariant under the Poincare and conformal transformations with

$$M^{\mu\nu}, P^\mu, D, K^\mu,$$

the generators of $SO(4,2)$

$SO(4,2)$ has a mathematical representation on AdS_5

AdS/CFT: Anti-de Sitter Space / Conformal Field Theory

Maldacena:

Map $AdS_5 \times S^5$ to conformal $N=4$ SUSY

- **QCD is not conformal**; however, it has manifestations of a scale-invariant theory: Bjorken scaling, dimensional counting for hard exclusive processes
- **Conformal window**: $\alpha_s(Q^2) \simeq \text{const}$ at small Q^2
- Use mathematical mapping of the conformal group $SO(4,2)$ to AdS_5 space

