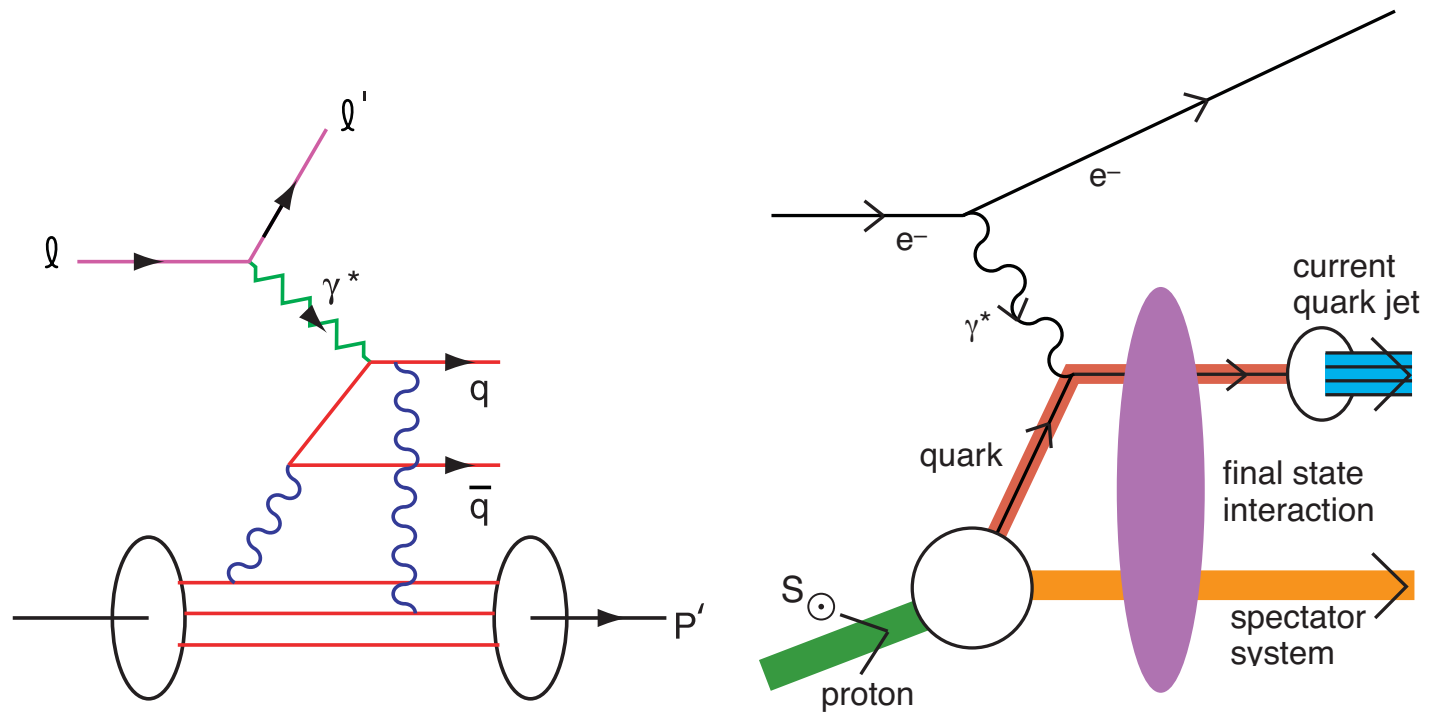


Novel QCD Phenomenology

BBC
NEWS

Stan Brodsky, SLAC/IPPP



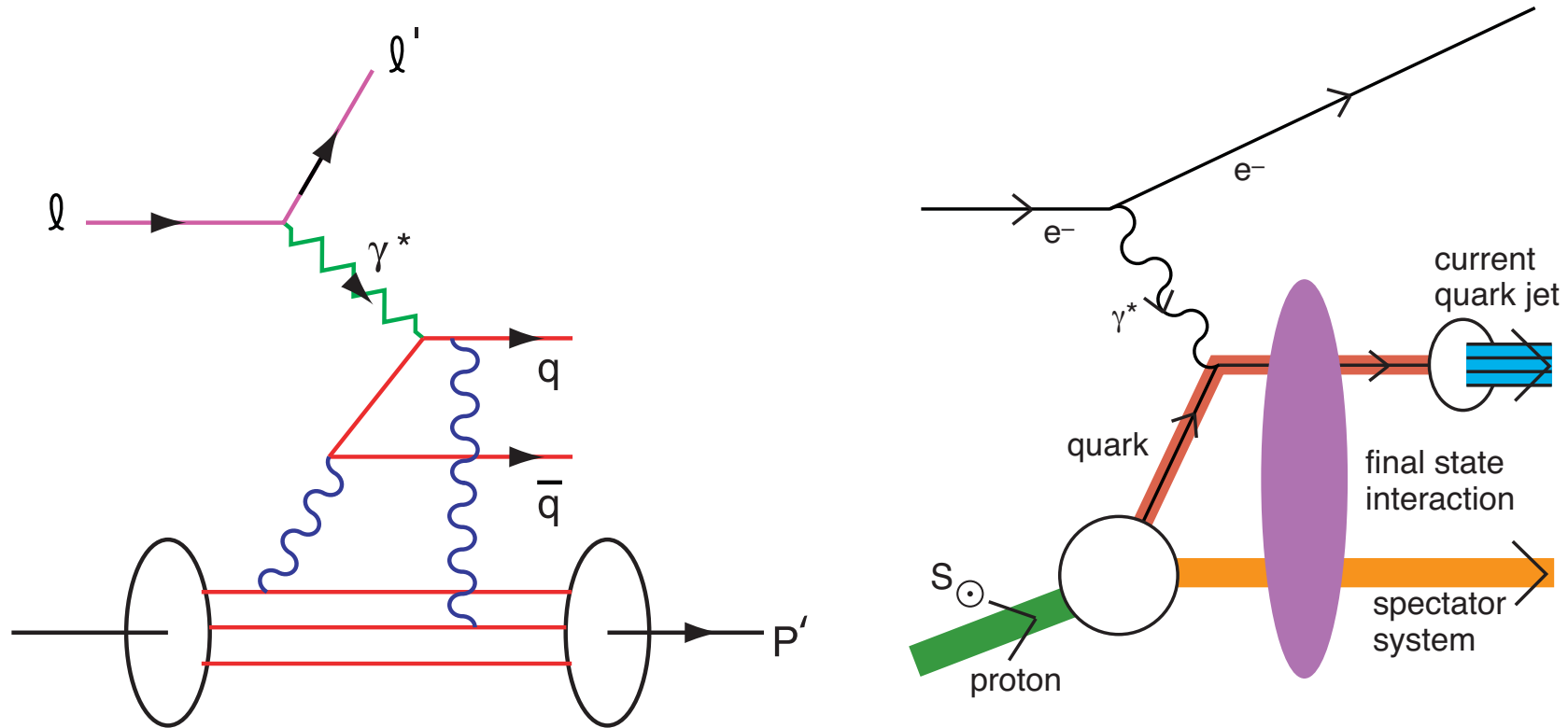
Imperial College, London

16 Sep 2008

Queen guitarist Brian May has handed in his astronomy PhD thesis - 36 years after abandoning it to join the band.

Novel QCD Phenomenology

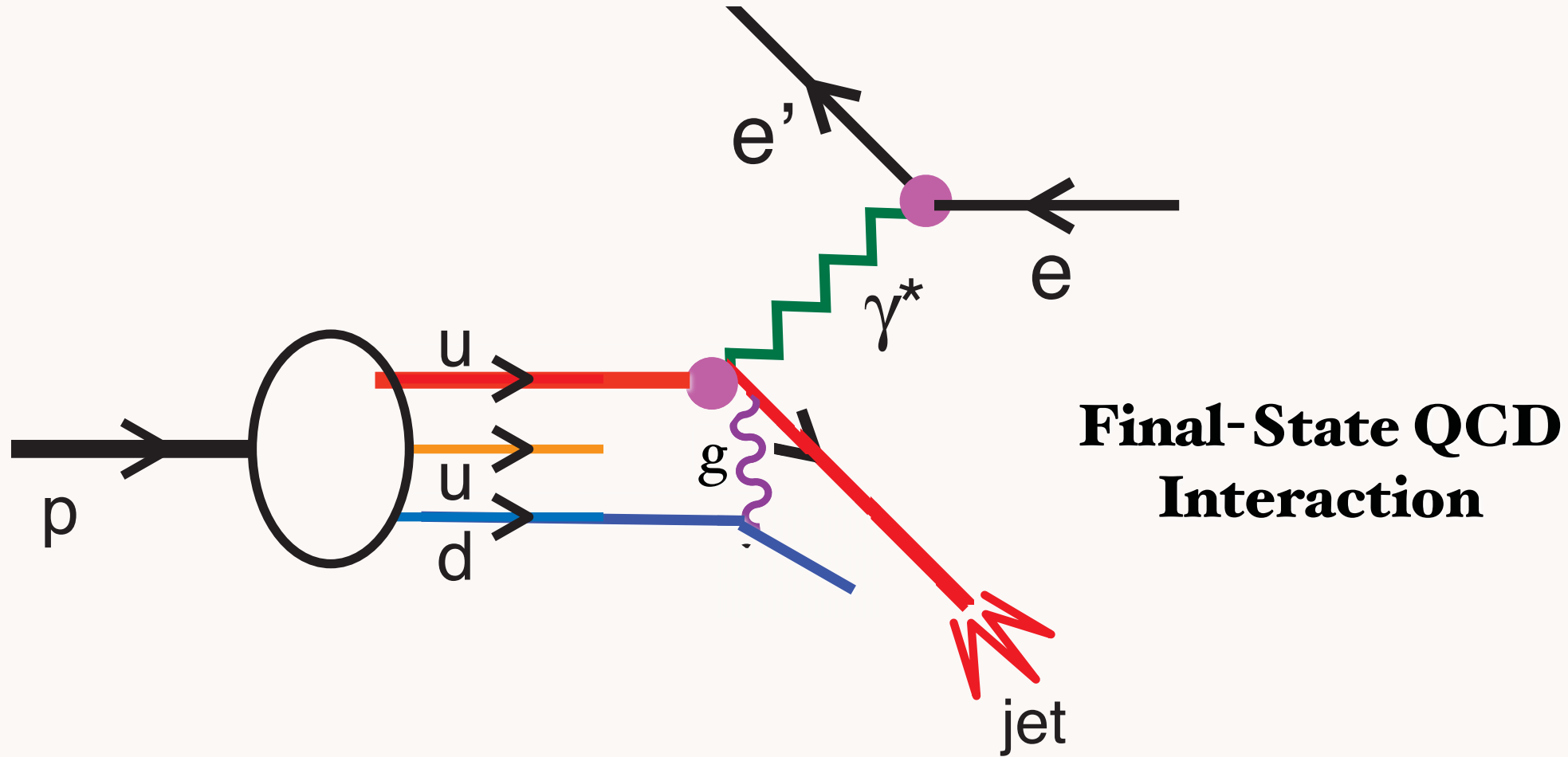
Stan Brodsky, SLAC/IPPP



Imperial College, London

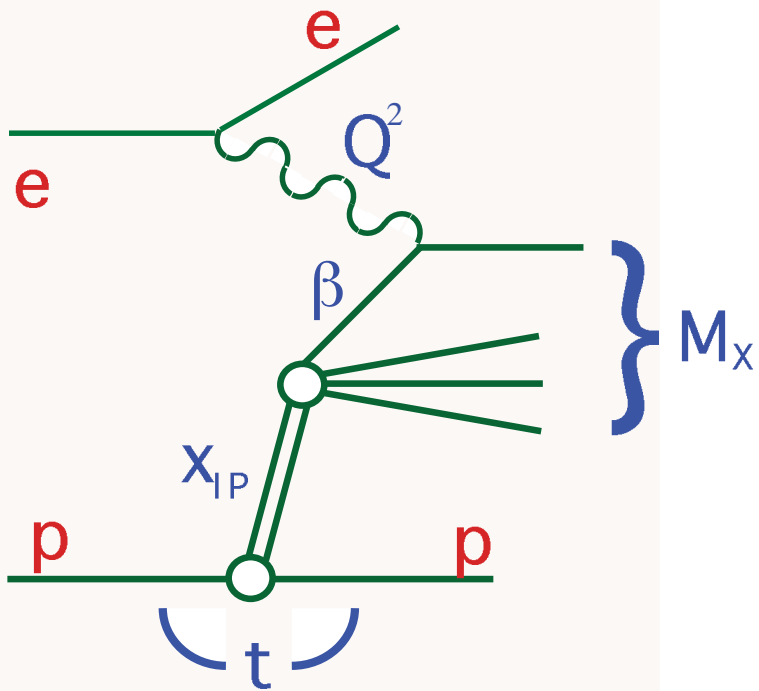
16 Sep 2008

Deep Inelastic Electron-Proton Scattering

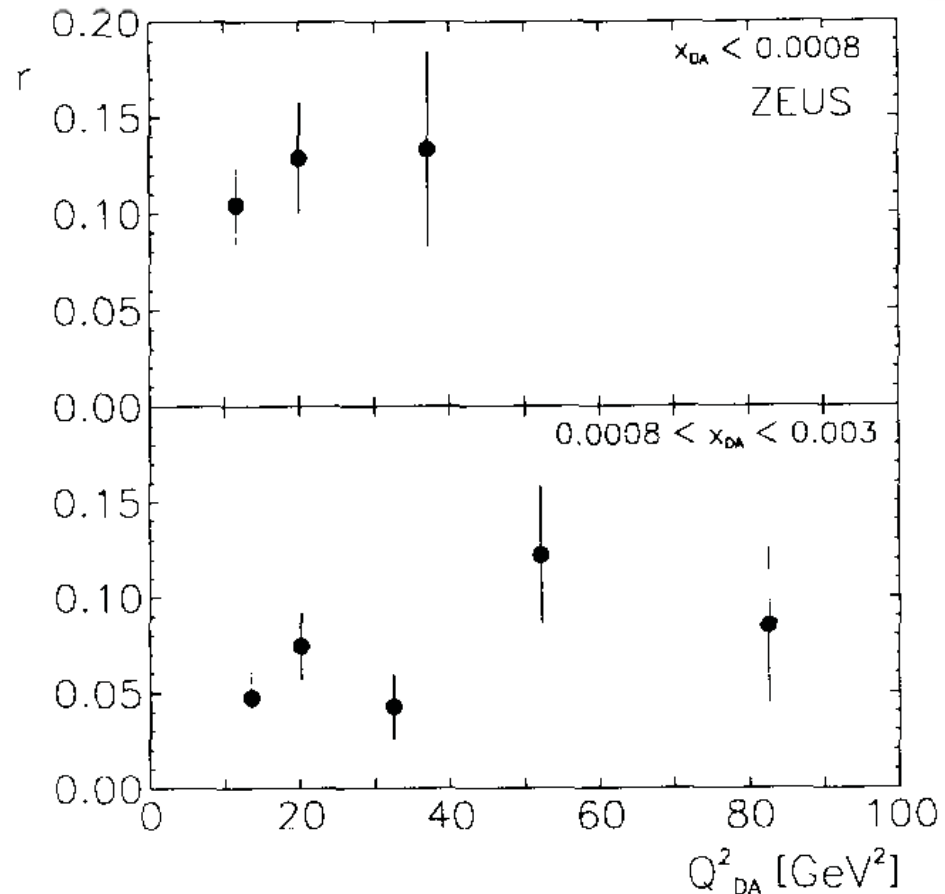


*Conventional wisdom:
Final-state interactions of struck quark can be neglected*

Remarkable observation at HERA



*10% to 15%
of DIS events
are
diffractive!*

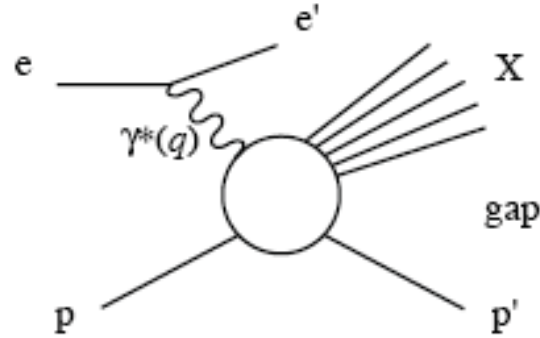


Fraction r of events with a large rapidity gap, $\eta_{\max} < 1.5$, as a function of Q^2_{DA} for two ranges of x_{DA} . No acceptance corrections have been applied.

M. Derrick et al. [ZEUS Collaboration], Phys. Lett. B 315, 481 (1993).

DDIS

Diffractive Deep Inelastic Lepton-Proton Scattering



- In a large fraction ($\sim 10\text{--}15\%$) of DIS events, the proton escapes intact, keeping a large fraction of its initial momentum
- This leaves a large *rapidity gap* between the proton and the produced particles
- The t -channel exchange must be *color singlet* \rightarrow a *pomeron*

Profound effect: target stays intact despite production of a massive system X

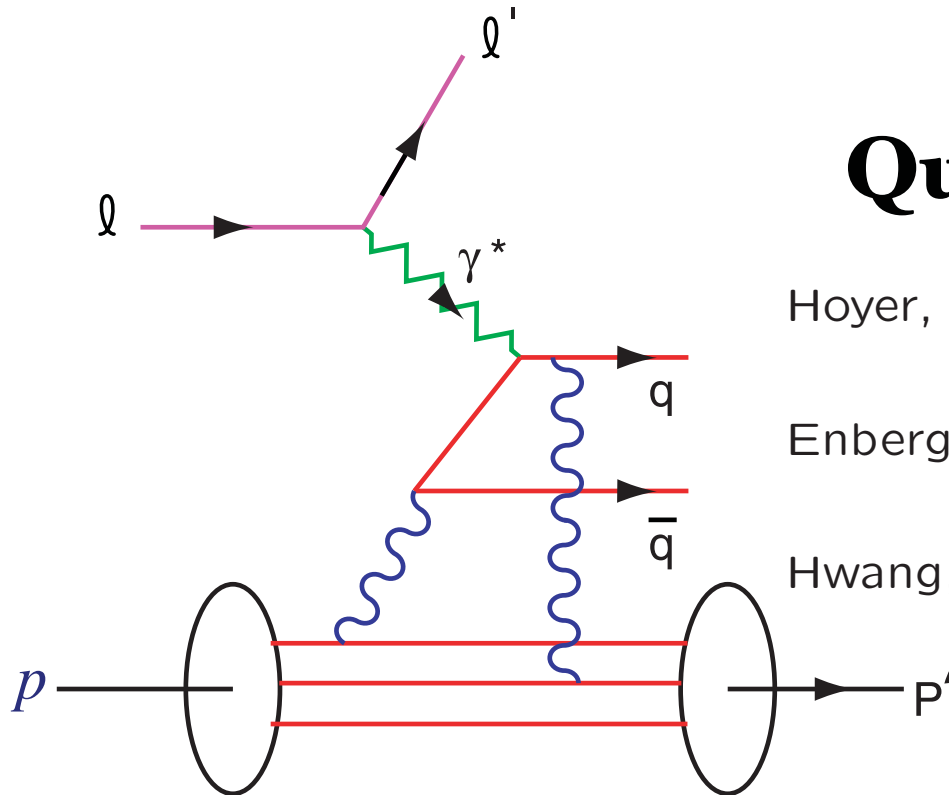
Final-State QCD Interaction Produces Diffractive DIS

Quark Rescattering

Hoyer, Marchal, Peigne, Sannino, SJB (BHMPS)

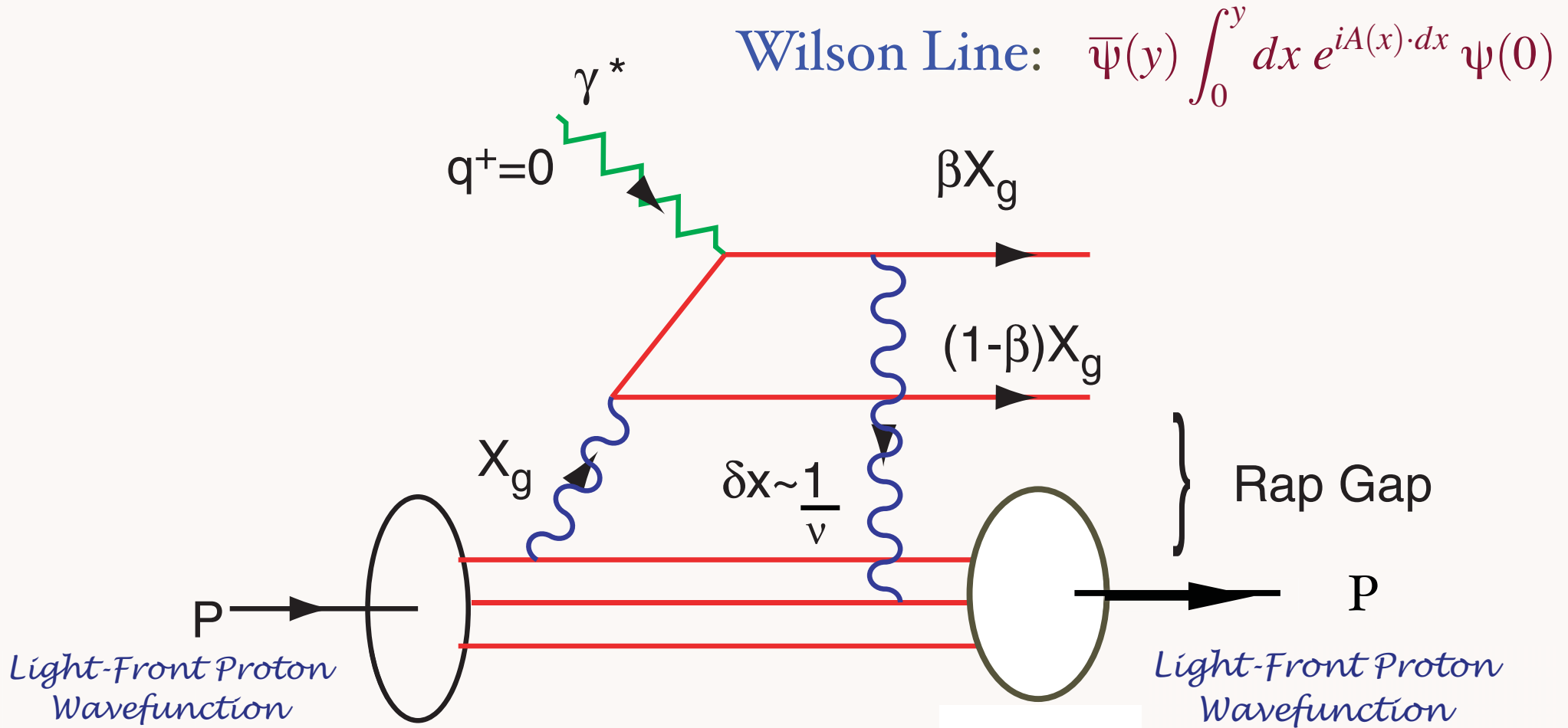
Enberg, Hoyer, Ingelman, SJB

Hwang, Schmidt, SJB



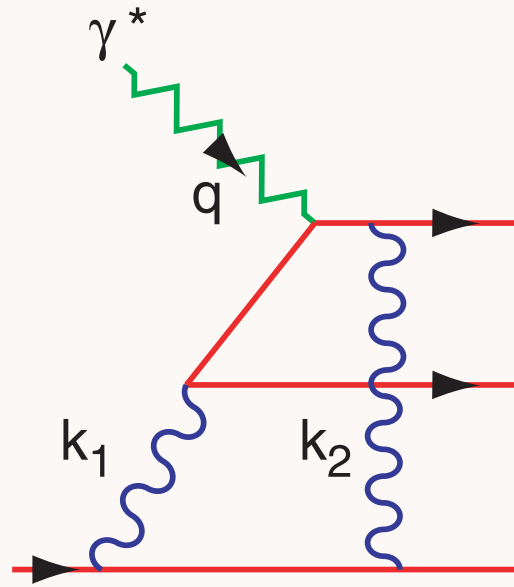
Low-Nussinov model of Pomeron

QCD Mechanism for Rapidity Gaps

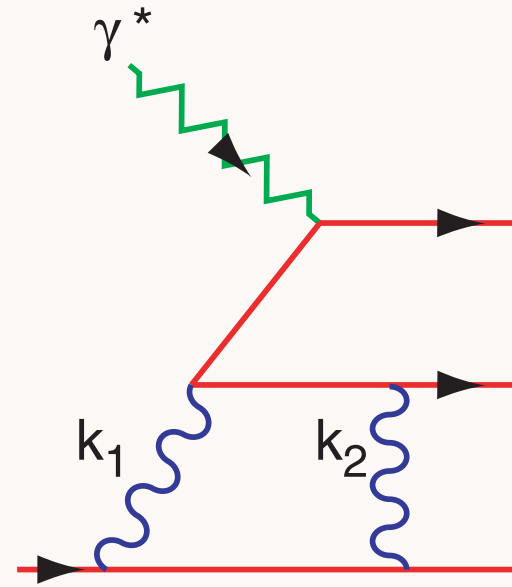


Reproduces lab-frame color dipole approach

Final State Interactions in QCD

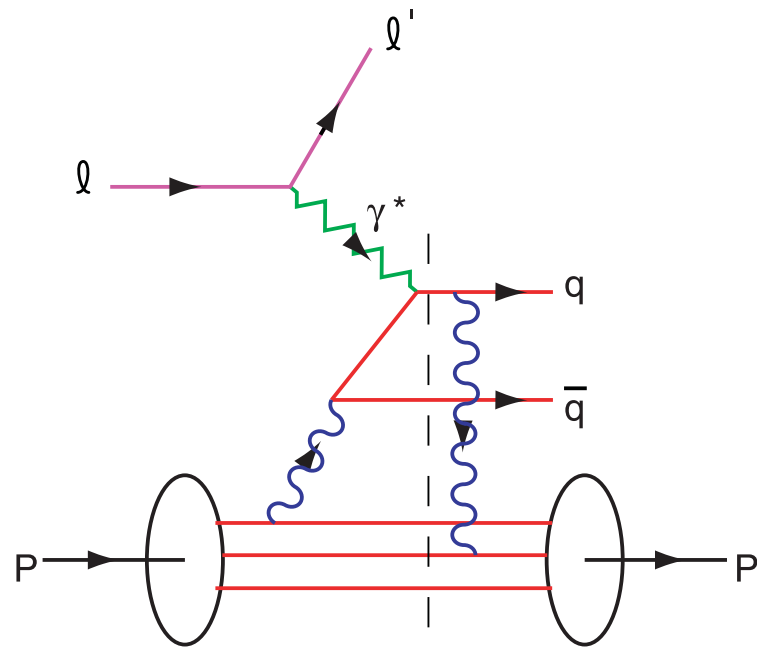


Feynman Gauge



Light-Cone Gauge

Result is Gauge Independent

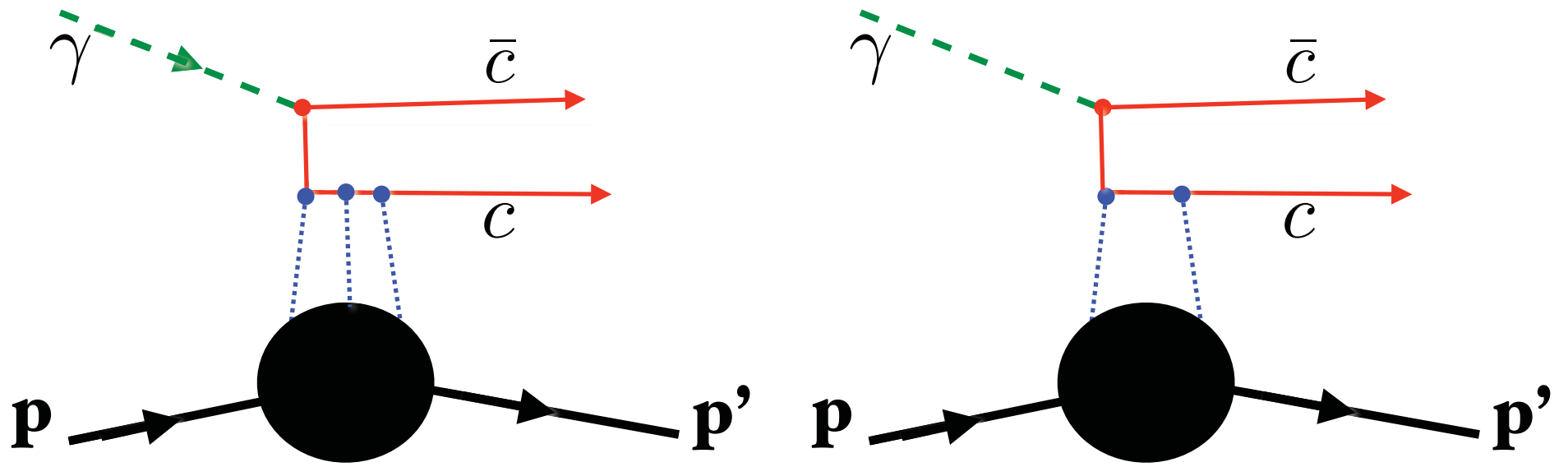


Integration over on-shell domain produces phase i

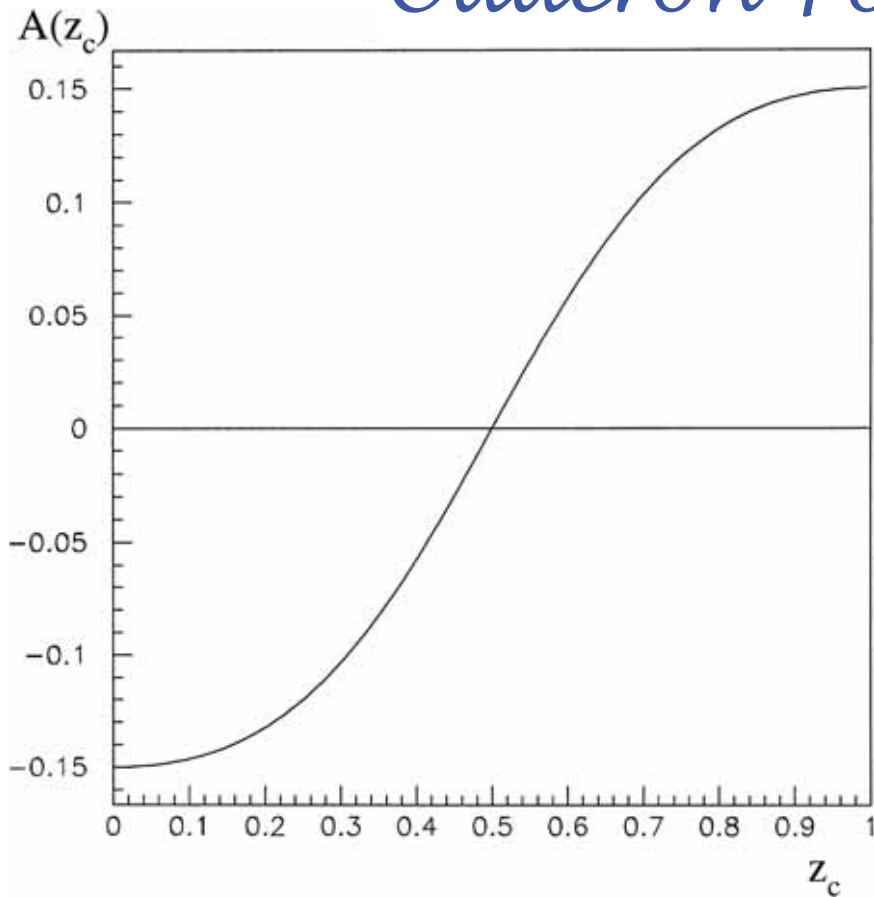
Need Imaginary Phase to Generate Pomeron

Need Imaginary Phase to Generate
T-Odd Single-Spin Asymmetry

Physics of FSI not in Wavefunction of Target



Odderon-Pomeron Interference!



$$\frac{d\sigma}{dz_c}(\gamma p \rightarrow c\bar{c}p')$$

$$\mathcal{A}(t \approx 0, M_X^2, z_c) \approx 0.45 \left(\frac{s_{\gamma p}}{M_X^2} \right)^{-0.25} \frac{2z_c - 1}{z_c^2 + (1 - z_c)^2}$$

Measure charm momentum asymmetry in photon fragmentation region

Only one charm quark needs to be measured

Merino, Rathsman, sjb

Single-spin asymmetries

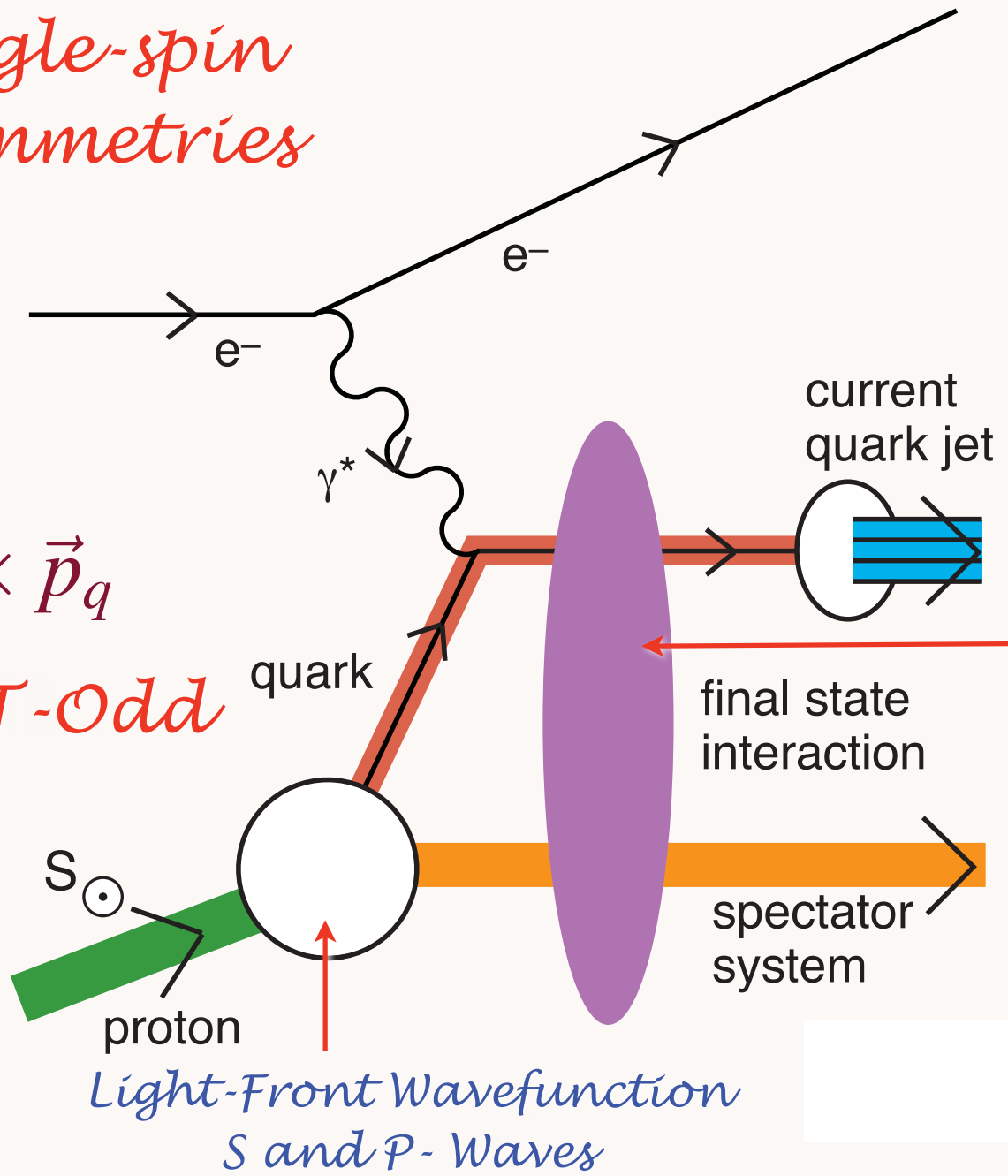
Leading Twist Sivers Effect

Hwang,
Schmidt, sjb

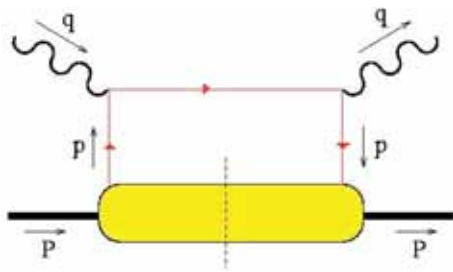
Collins, Burkardt
Ji, Yuan

*QCD S- and P-Coulomb Phases
--Wilson Line*

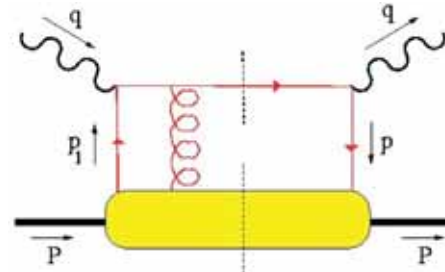
$i \vec{S}_p \cdot \vec{q} \times \vec{p}_q$
Pseudo-T-Odd



*Light-Front Wavefunction
S and P-Waves*



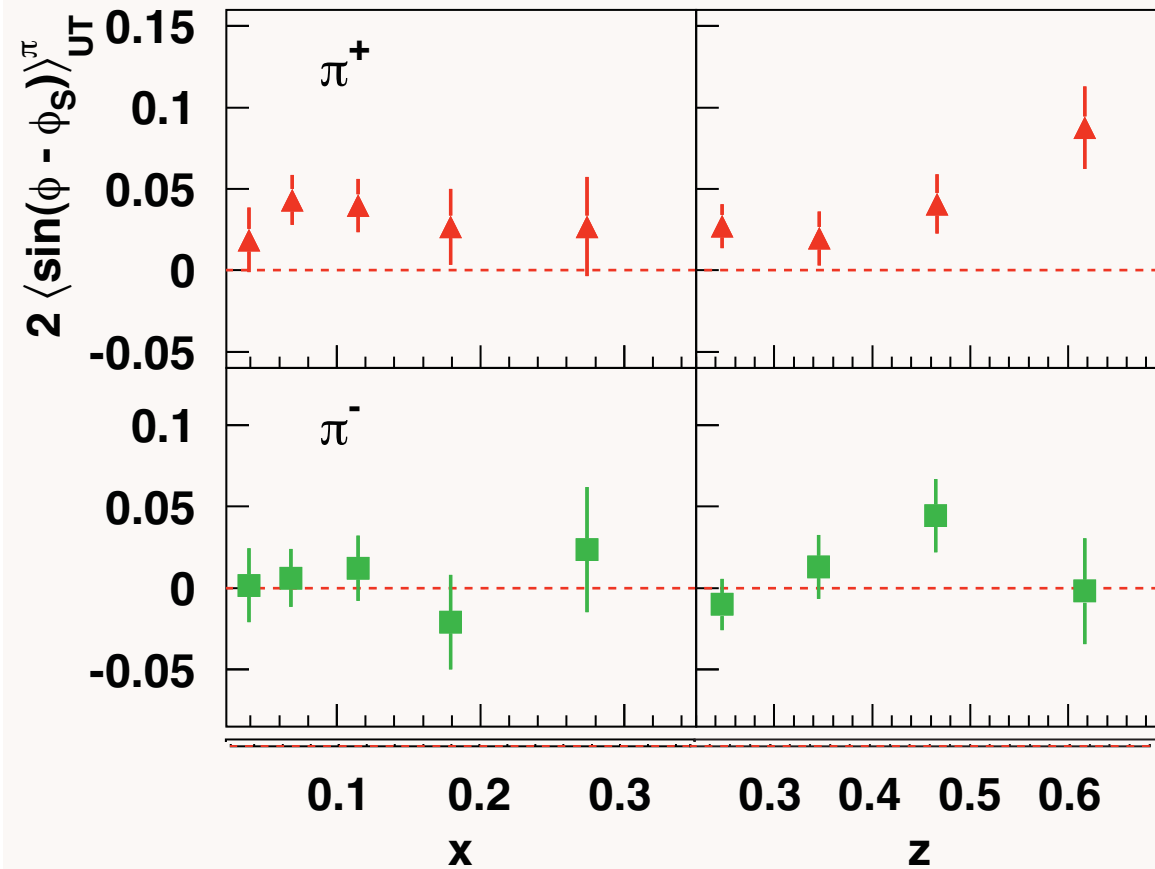
can interfere with



and produce a T-odd effect!
(also need $L_z \neq 0$)

HERMES coll., A. Airapetian et al., Phys. Rev. Lett. 94 (2005) 012002.

Sivers asymmetry from HERMES



- First evidence for non-zero Sivers function!
- \Rightarrow presence of non-zero quark orbital angular momentum!
- **Positive** for π^+ ...
Consistent with zero for π^- ...

Gamberg: Hermes data compatible with BHS model

Schmidt, Lu: Hermes charge pattern follow quark contributions to anomalous moment

Imperial College,
September 16, 2008

Novel QCD Physics

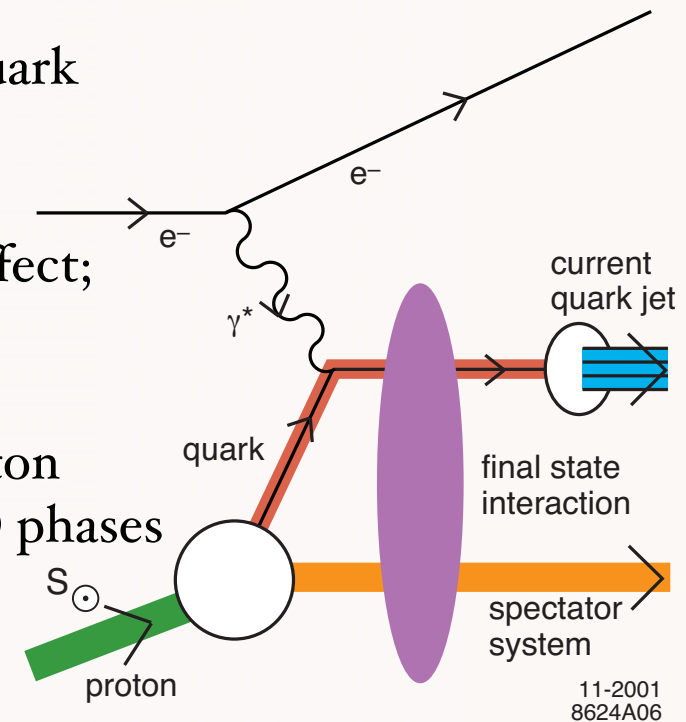
12

Stan Brodsky, SLAC

Final-State Interactions Produce Pseudo-T-Odd (Sivers Effect)

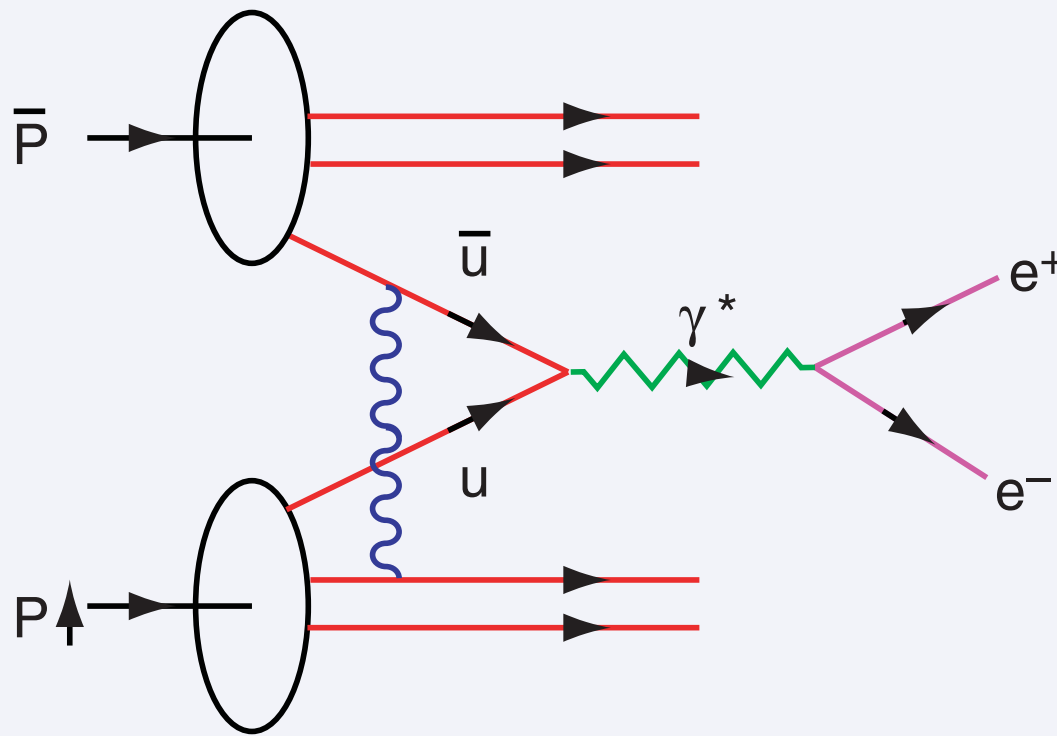
- Leading-Twist Bjorken Scaling!
- Requires nonzero orbital angular momentum of quark
- Arises from the interference of Final-State QCD Coulomb phases in S- and P- waves; Wilson line effect; gauge independent
- Relate to the quark contribution to the target proton anomalous magnetic moment and final-state QCD phases
- QCD phase at soft scale!
- New window to QCD coupling and running gluon mass in the IR
- QED S and P Coulomb phases infinite -- difference of phases finite!

$$\mathbf{i} \vec{S} \cdot \vec{p}_{jet} \times \vec{q}$$



11-2001
8624A06

Predict Opposite Sign SSA in DY !



Collins;
Hwang, Schmidt.
sjb

Single Spin Asymmetry In the Drell Yan Process

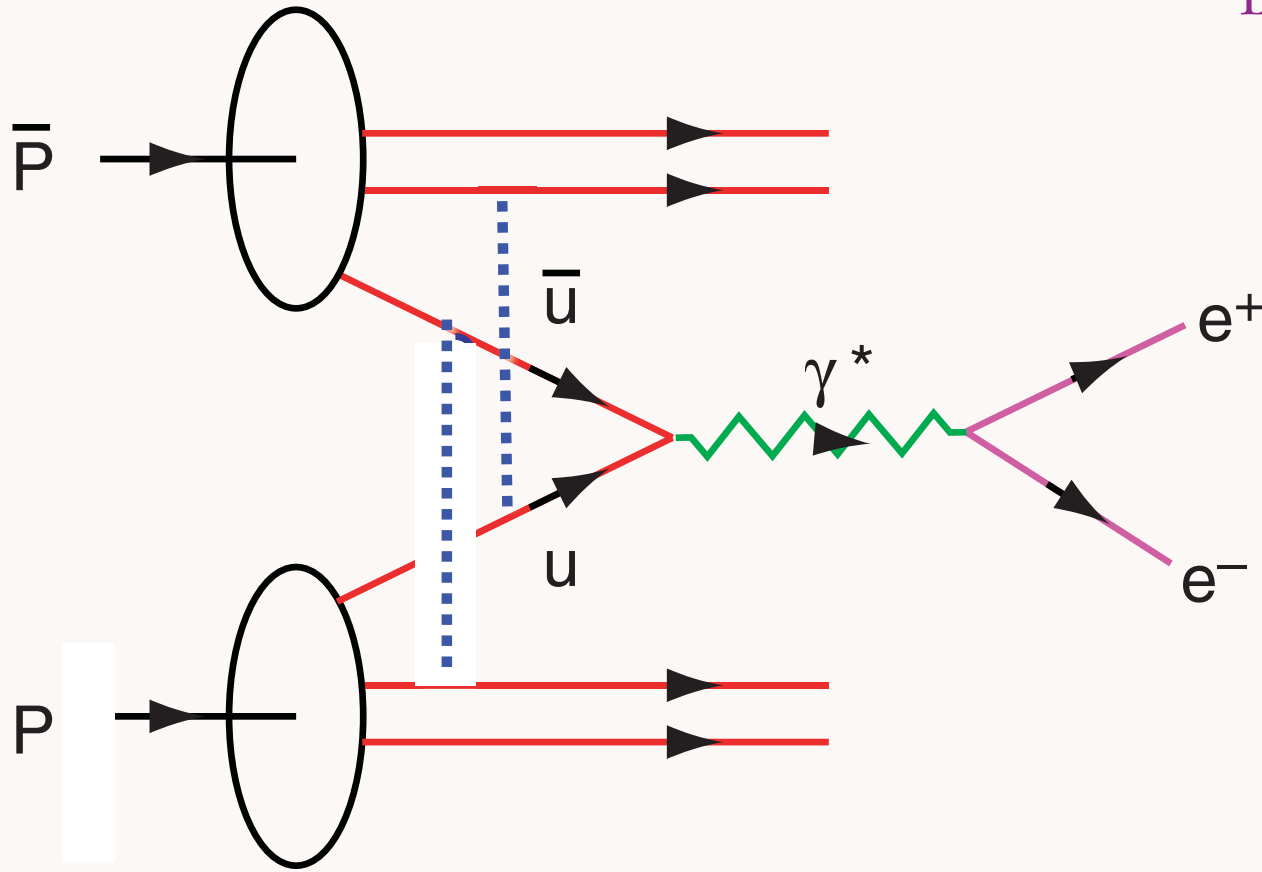
$$\vec{S}_p \cdot \vec{p} \times \vec{q}_{\gamma^*}$$

Quarks Interact in the Initial State

Interference of Coulomb Phases for S and P states

Produce Single Spin Asymmetry [Siver's Effect] Proportional to the Proton Anomalous Moment and α_s .

Opposite Sign to DIS! No Factorization



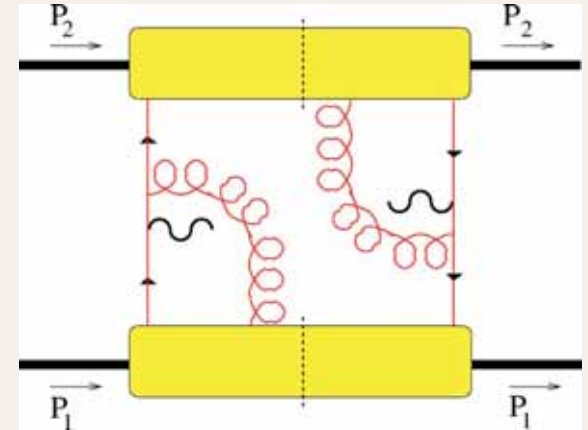
$DY \cos 2\phi$ correlation at leading twist from double ISI

Anomalous effect from Double ISI in Massive Lepton Production

Boer, Hwang, sjb

$\cos 2\phi$ correlation

- Leading Twist, valence quark dominated
- Violates Lam-Tung Relation!
- Not obtained from standard PQCD subprocess analysis
- Normalized to the square of the single spin asymmetry in semi-inclusive DIS
- No polarization required
- Challenge to standard picture of PQCD Factorization



Double Initial-State Interactions

generate anomalous $\cos 2\phi$

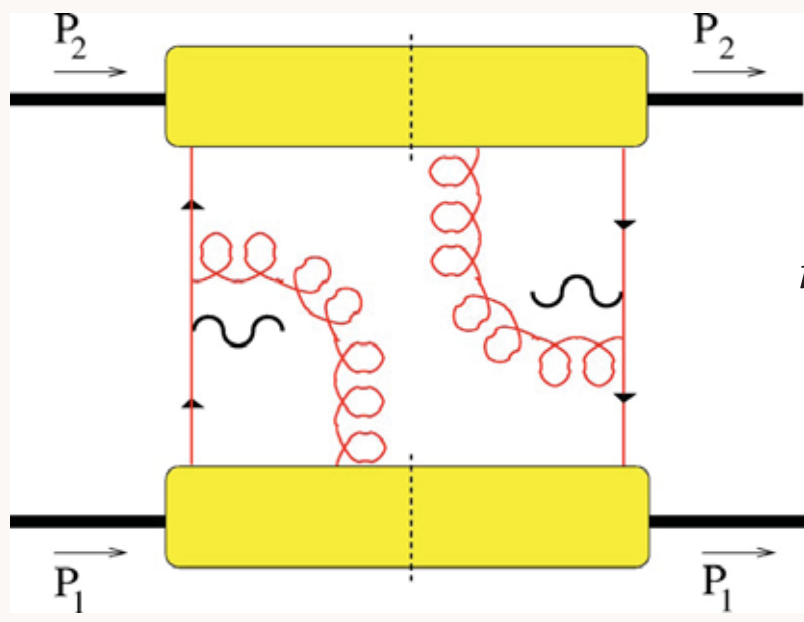
Boer, Hwang, sjb

Drell-Yan planar correlations

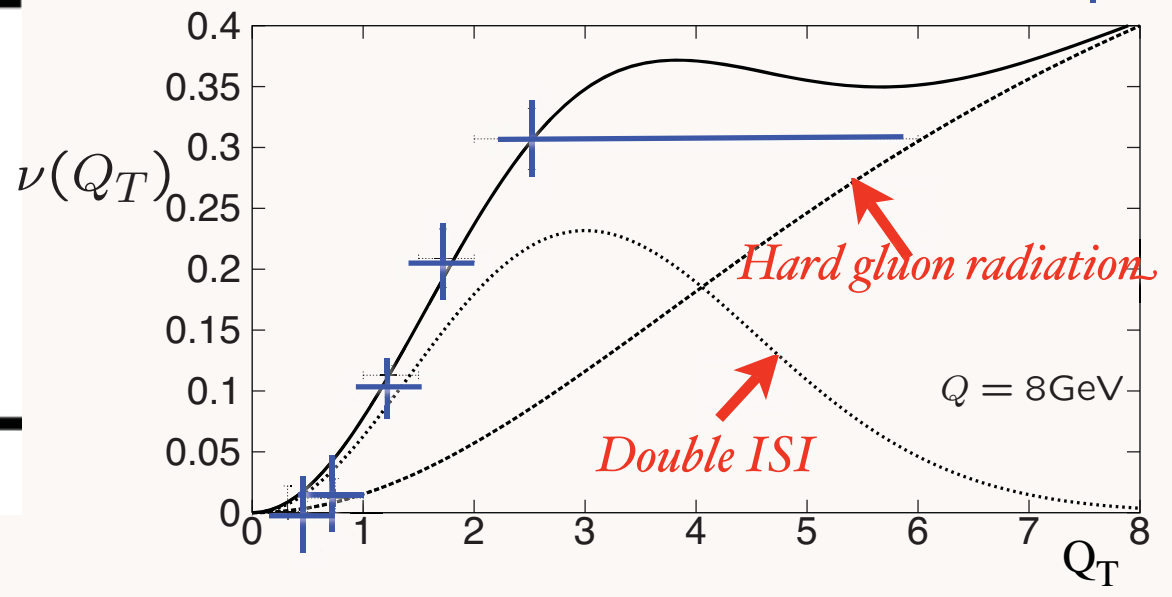
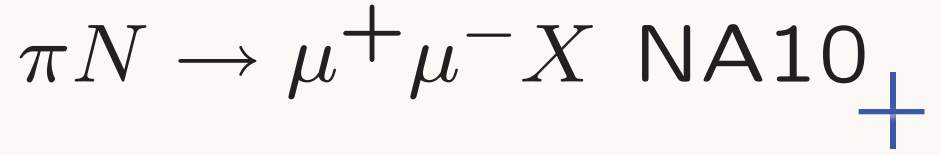
$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} \propto \left(1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi \right)$$

PQCD Factorization (Lam Tung): $1 - \lambda - 2\nu = 0$

$$\frac{\nu}{2} \propto h_1^\perp(\pi) h_1^\perp(N)$$

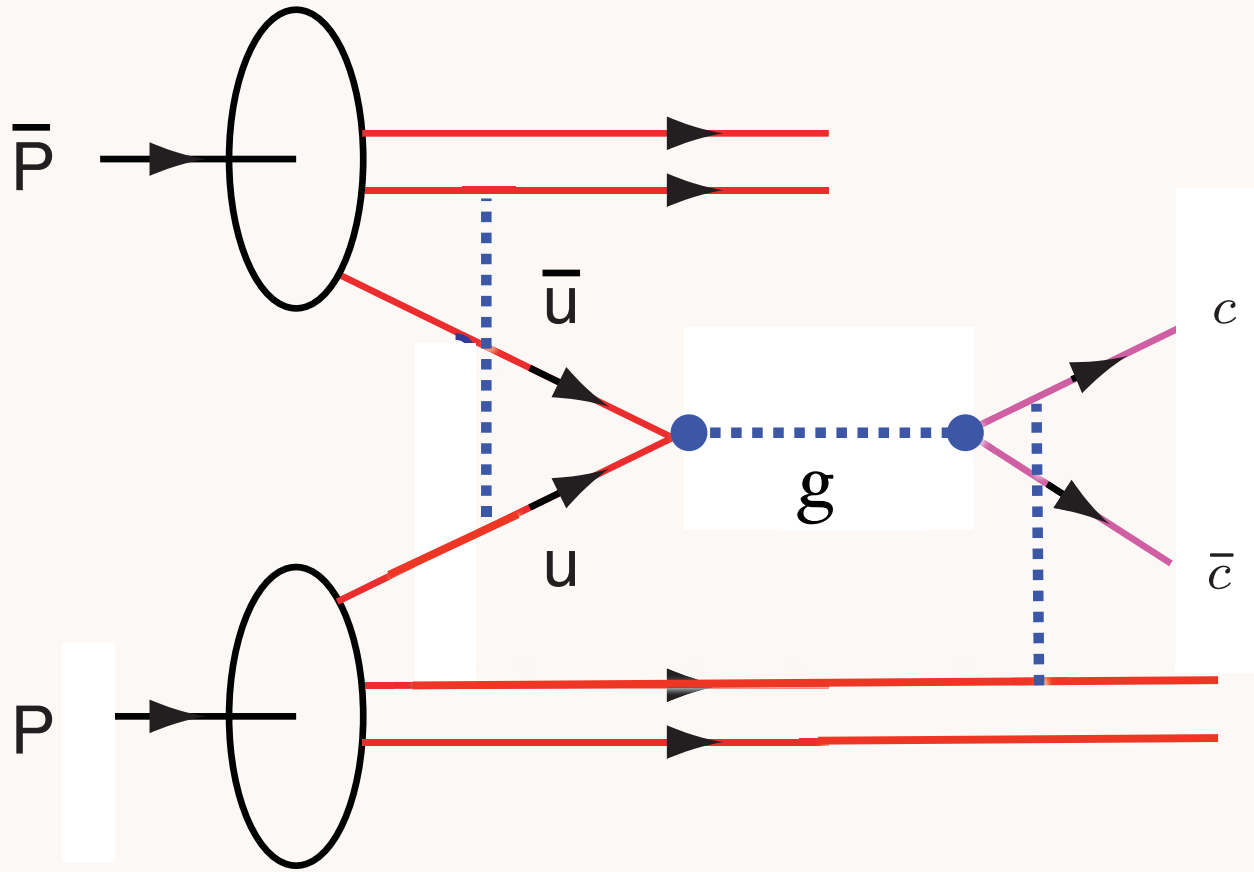


Violates Lam-Tung relation!



Model: Boer,

Stan Brodsky, SLAC

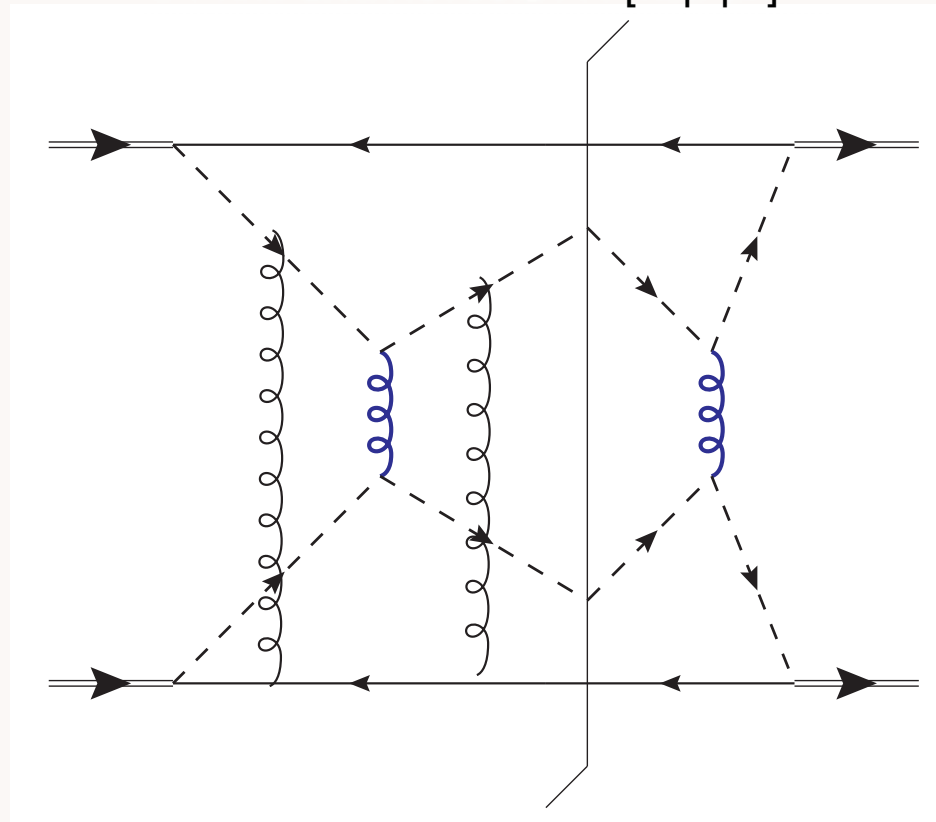


Problem for factorization when both ISI and FSI occur

Factorization is violated in production of high-transverse-momentum particles in hadron-hadron collisions

John Collins, [Jian-Wei Qiu](#) . ANL-HEP-PR-07-25, May 2007.

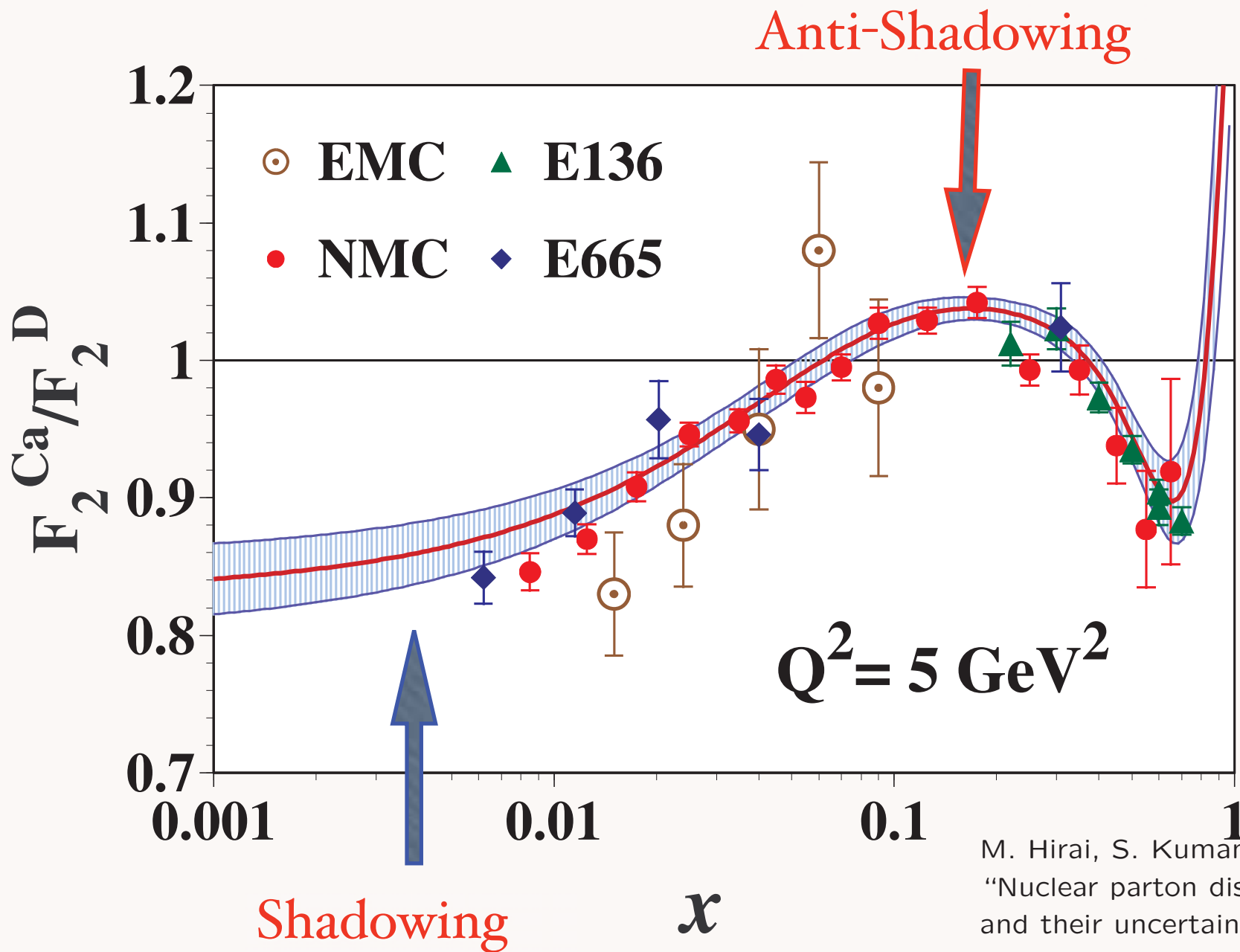
e-Print: [arXiv:0705.2141](#) [hep-ph]



The exchange of two extra gluons, as in this graph, will tend to give non-factorization in unpolarized cross sections.

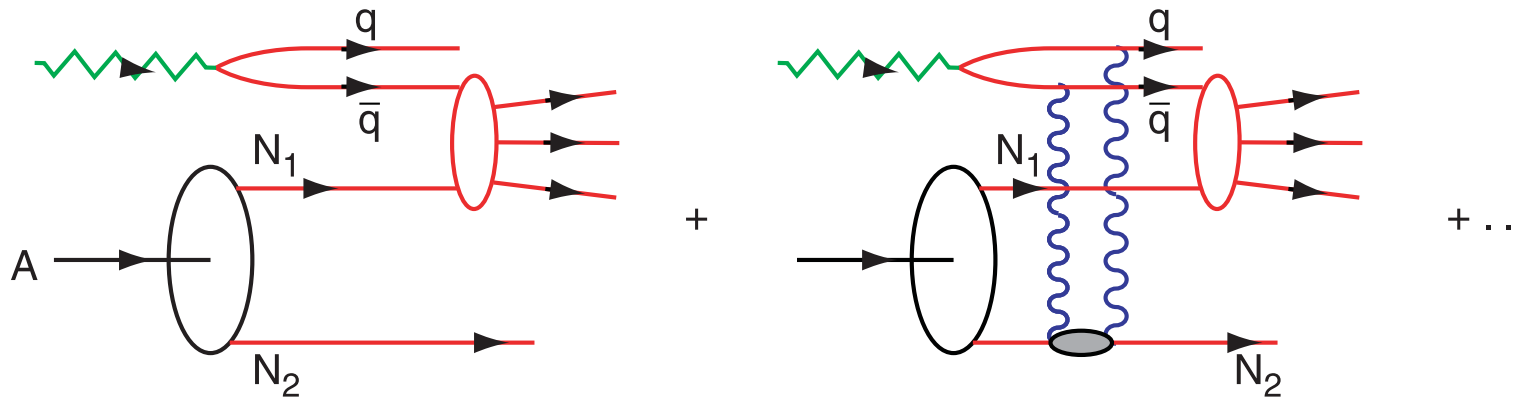
Novel Aspects of QCD in ep scattering

- **Initial and final-state interactions are **not** power suppressed DIS; Wilson line correction to handbag diagram in DVCS**
- **Leading-twist Bjorken-scaling single-spin asymmetry:**
- **Leading-twist Bjorken-scaling Diffractive DIS**
- **Diffractive Electroproduction; Color Transparency**
- **DIS at high energy reflects interactions of color-dipole of virtual photon with proton and nucleus: shadowing, saturation:**
- **Breakdown of parton model concepts: Structure functions are **not** probability distributions**
- **Nuclear LFWFS are universal, but the measured nuclear parton distributions are **not** universal -- **antishadowing is flavor-dependent****



M. Hirai, S. Kumano and T. H. Nagai,
 "Nuclear parton distribution functions
 and their uncertainties,"
 Phys. Rev. C **70**, 044905 (2004)
 [arXiv:hep-ph/0404093].

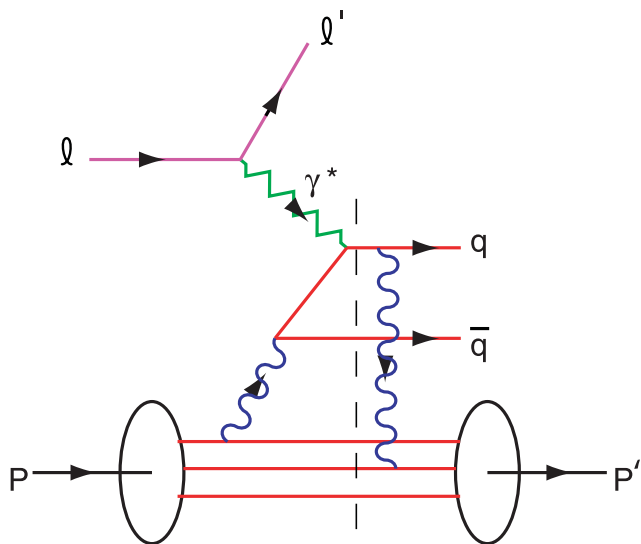
Nuclear Shadowing in QCD



Shadowing depends on understanding leading twist-diffraction in DIS

Nuclear Shadowing not included in nuclear LFWF !

Dynamical effect due to virtual photon interacting in nucleus



Shadowing depends on leading-twist DDIS

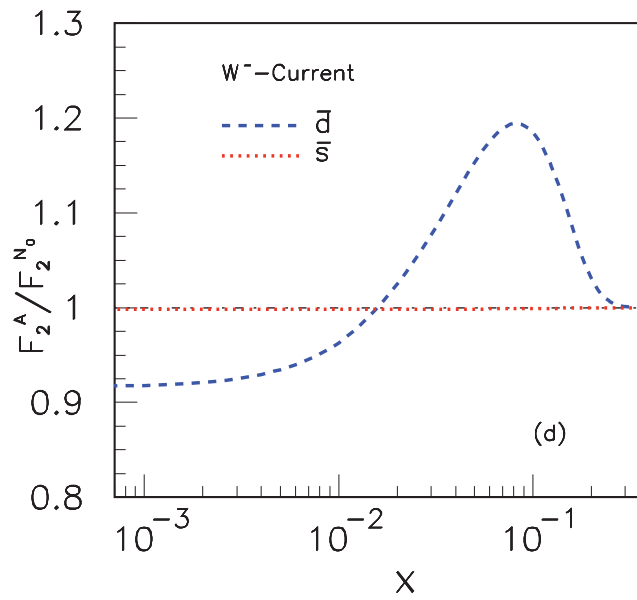
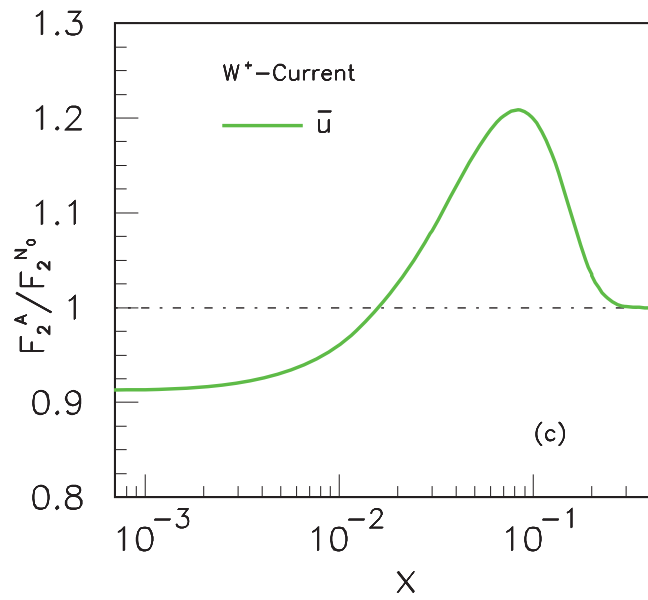
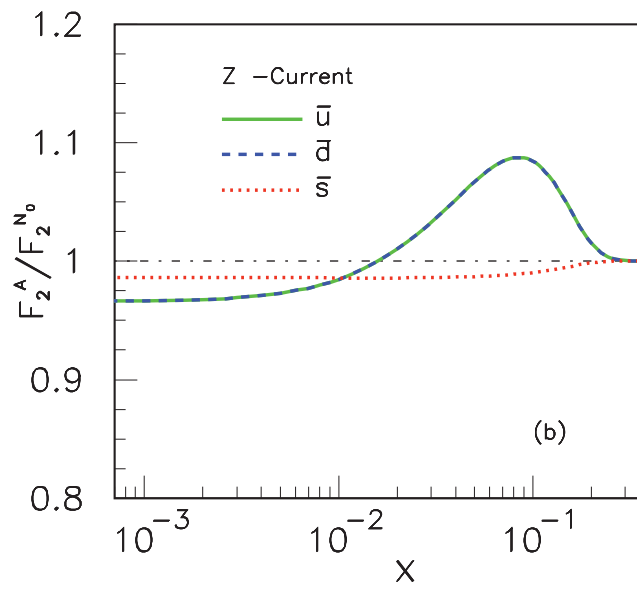
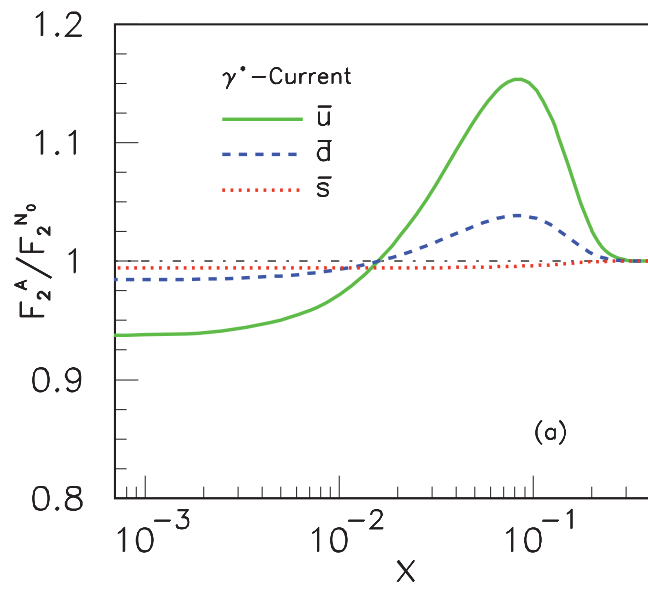
Integration over on-shell domain produces phase i

Need Imaginary Phase to Generate Pomeron

Need Imaginary Phase to Generate T-Odd Single-Spin Asymmetry

Physics of FSI not in Wavefunction of Target

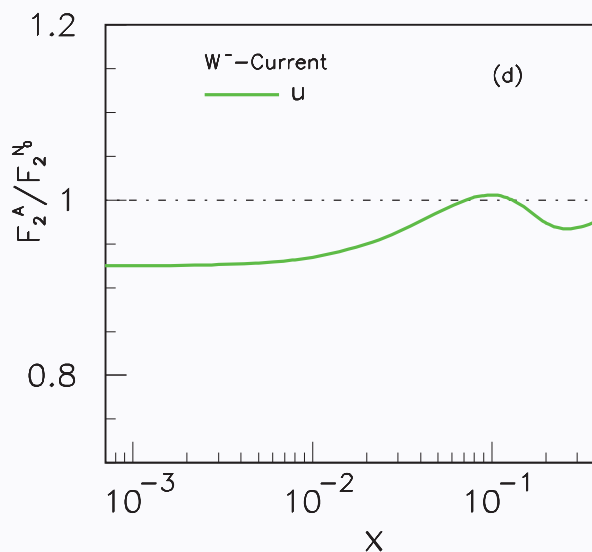
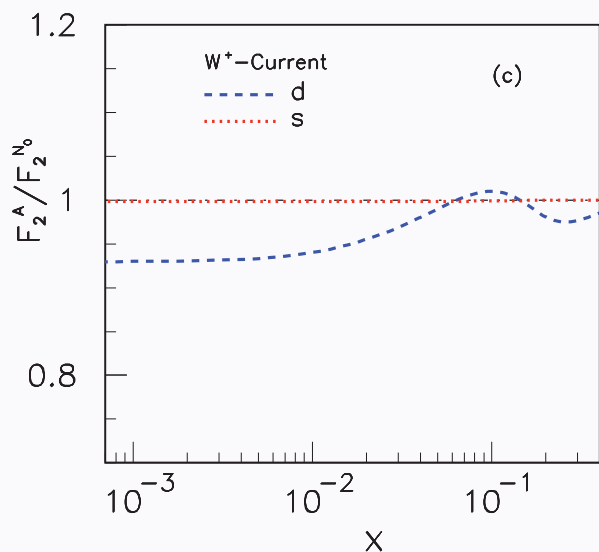
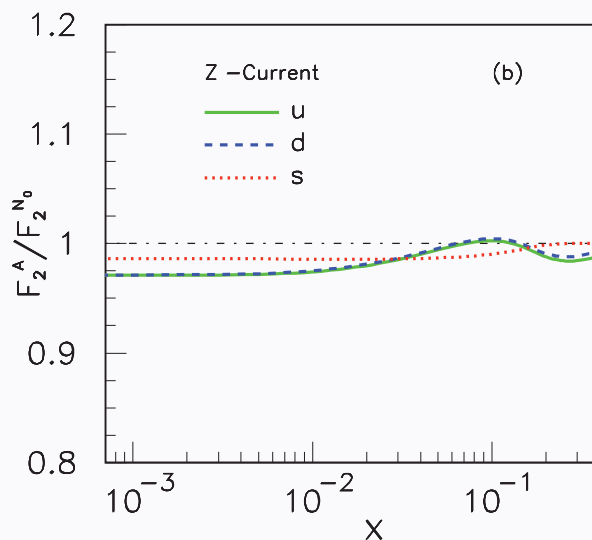
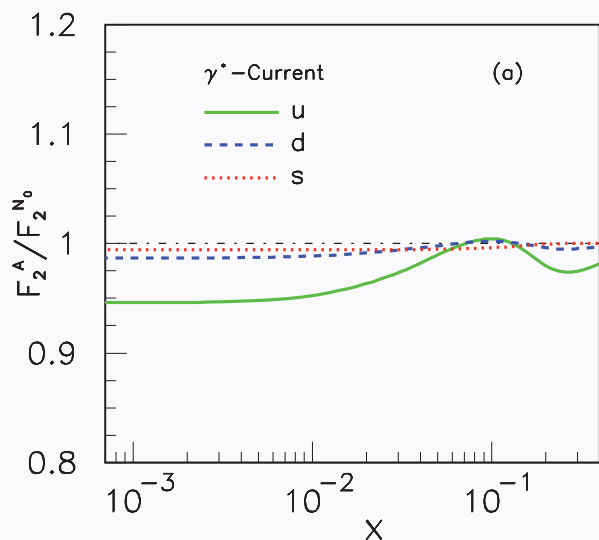
Antishadowing (Reggeon exchange) is not universal!



Schmidt, Yang; sjb

Nuclear Antishadowing not universal !

Shadowing and Antishadowing of DIS Structure Functions



S. J. Brodsky, I. Schmidt and J. J. Yang,
 “Nuclear Antishadowing in
 Neutrino Deep Inelastic Scattering,”
 Phys. Rev. D 70, 116003 (2004)
 [arXiv:hep-ph/0409279].

**Modifies
 NuTeV extraction of
 $\sin^2 \theta_W$**

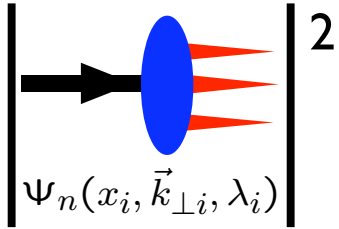
**Test in flavor-tagged
 lepton-nucleus collisions**

Physics of Rescattering

- Diffractive DIS
- Non-Unitary Correction to DIS: Structure functions are not probability distributions
- Nuclear Shadowing, Antishadowing- Not in Target WF
- Single Spin Asymmetries -- opposite sign in DY and DIS
- DY angular distribution at leading twist from double ISI-- not given by PQCD factorization -- breakdown of factorization!
- Wilson Line Effects not 1 even in LCG
- Must correct hard subprocesses for initial and final-state soft gluon attachments
- Corrections to Handbag Approximation in DVCS

Hoyer, Marchal, Peigne, Sannino, sjb

Static vs. Dynamic Structure Functions

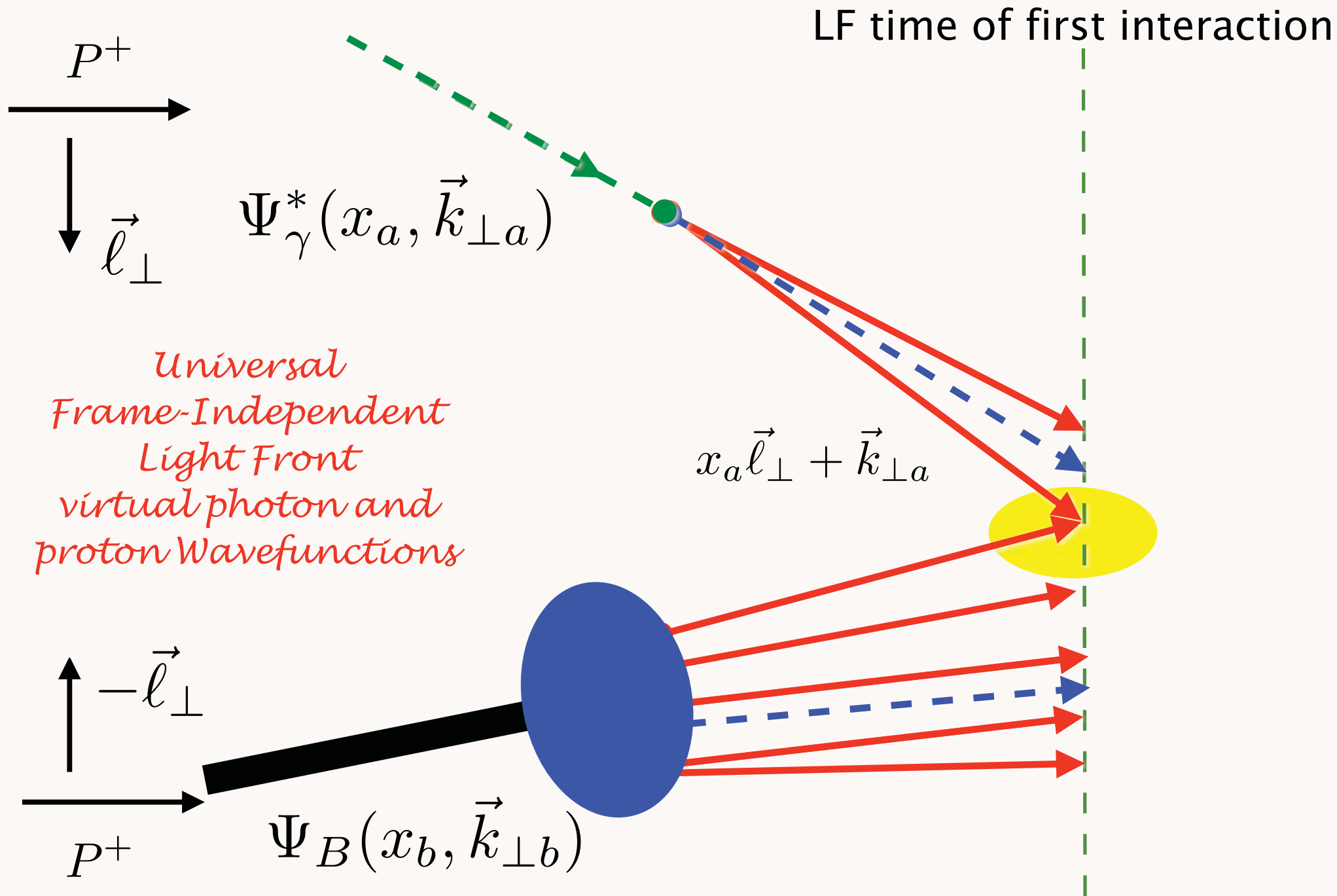


Static

- Square of Target LFWFs
- No Wilson Line
- Probability Distributions
- Process-Independent
- T-even Observables
- No Shadowing, Anti-Shadowing
- Sum Rules: Momentum and J^z
- DGLAP Evolution; mod. at large x
- No Diffractive DIS

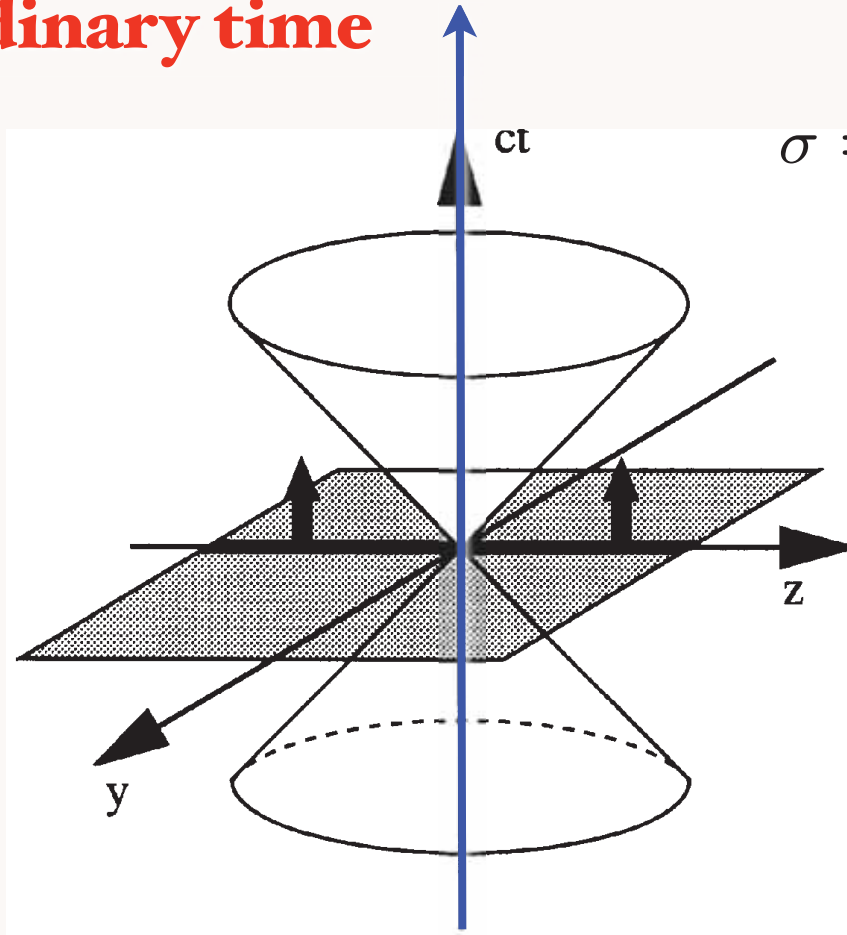
Dynamic

- Modified by Rescattering: ISI & FSI
- Contains Wilson Line, Phases
- No Probabilistic Interpretation
- Process-Dependent - From Collision
- T-Odd (Sivers, Boer-Mulders, etc.)
- Shadowing, Anti-Shadowing, Saturation
- Not Proven
- DGLAP Evolution
- Hard Pomeron and Odderon: DDIS



Dirac's Amazing Idea: The Front Form

**Evolve in
ordinary time**

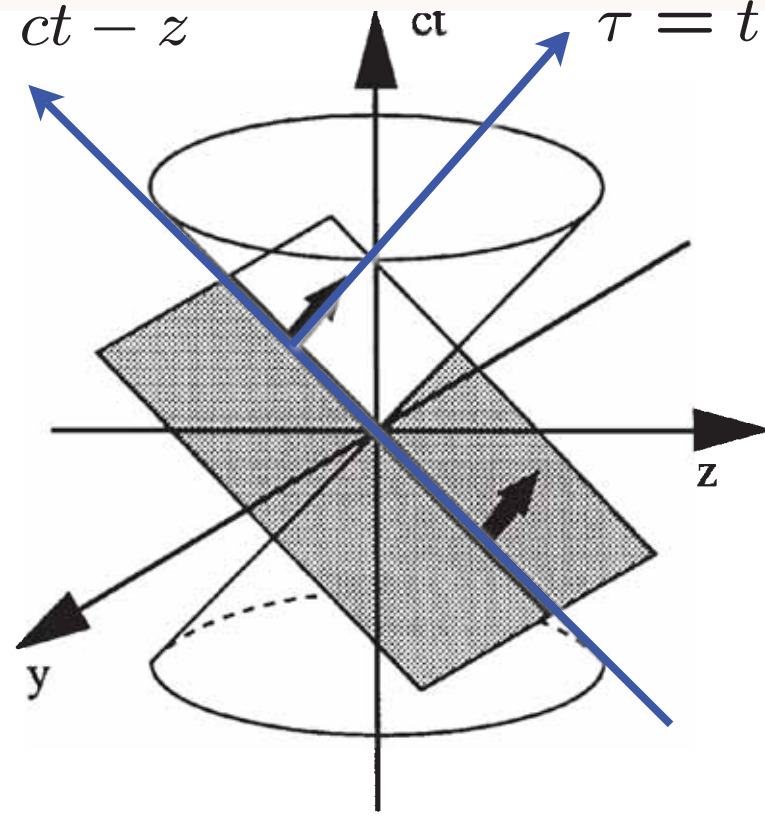


Instant Form

**Evolve in
light-front time!**

$$\sigma = ct - z$$

$$\tau = t + z/c$$



Front Form

*Each element of
flash photograph
illuminated
at same LF time*

$$\tau = t + z/c$$

Evolve in LF time

$$P^- = i \frac{d}{d\tau}$$

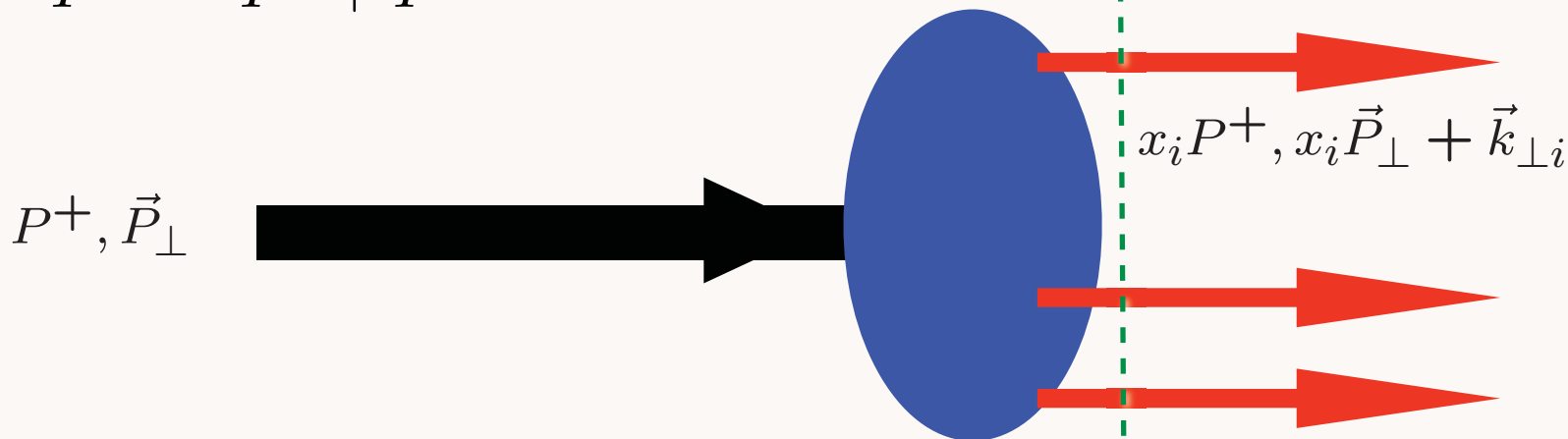
Eigenstate -- independent of τ



Light-Front Wavefunctions: rigorous representation of composite systems in quantum field theory

$$x = \frac{k^+}{P^+} = \frac{k^0 + k^3}{P^0 + P^3}$$

Fixed $\tau = t + z/c$



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

$$\sum_i^n x_i = 1$$

$$\sum_i^n \vec{k}_{\perp i} = \vec{0}_{\perp}$$

Invariant under boosts! Independent of P^μ

Light-Front Wavefunctions

Dirac's Front Form: Fixed $\tau = t + z/c$

$$\Psi(x, k_{\perp}) \quad x_i = \frac{k_i^+}{P^+}$$

Invariant under boosts. Independent of p^{μ}

$$H_{LF}^{QCD} |\psi\rangle = M^2 |\psi\rangle$$

Direct connection to QCD Lagrangian

*Remarkable new insights from AdS/CFT,
the duality between conformal field theory
and Anti-de Sitter Space*

Light-Front QCD

Heisenberg Matrix Formulation

Physical gauge: $A^+ = 0$

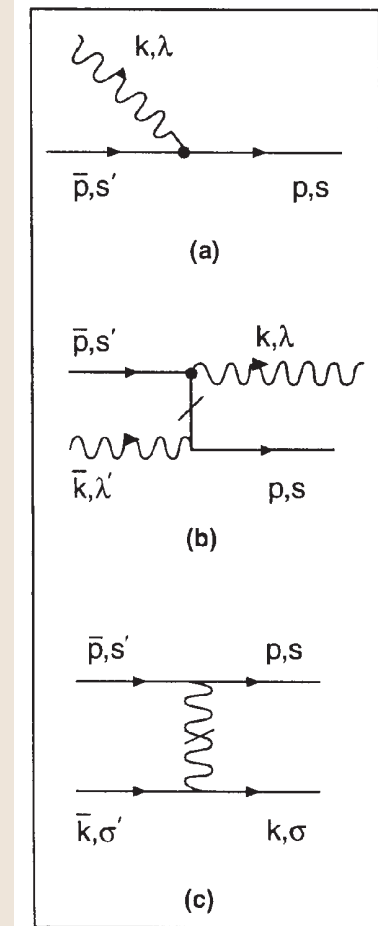
$$L^{QCD} \rightarrow H_{LF}^{QCD}$$

$$H_{LF}^{QCD} = \sum_i \left[\frac{m^2 + k_{\perp}^2}{x} \right]_i + H_{LF}^{int}$$

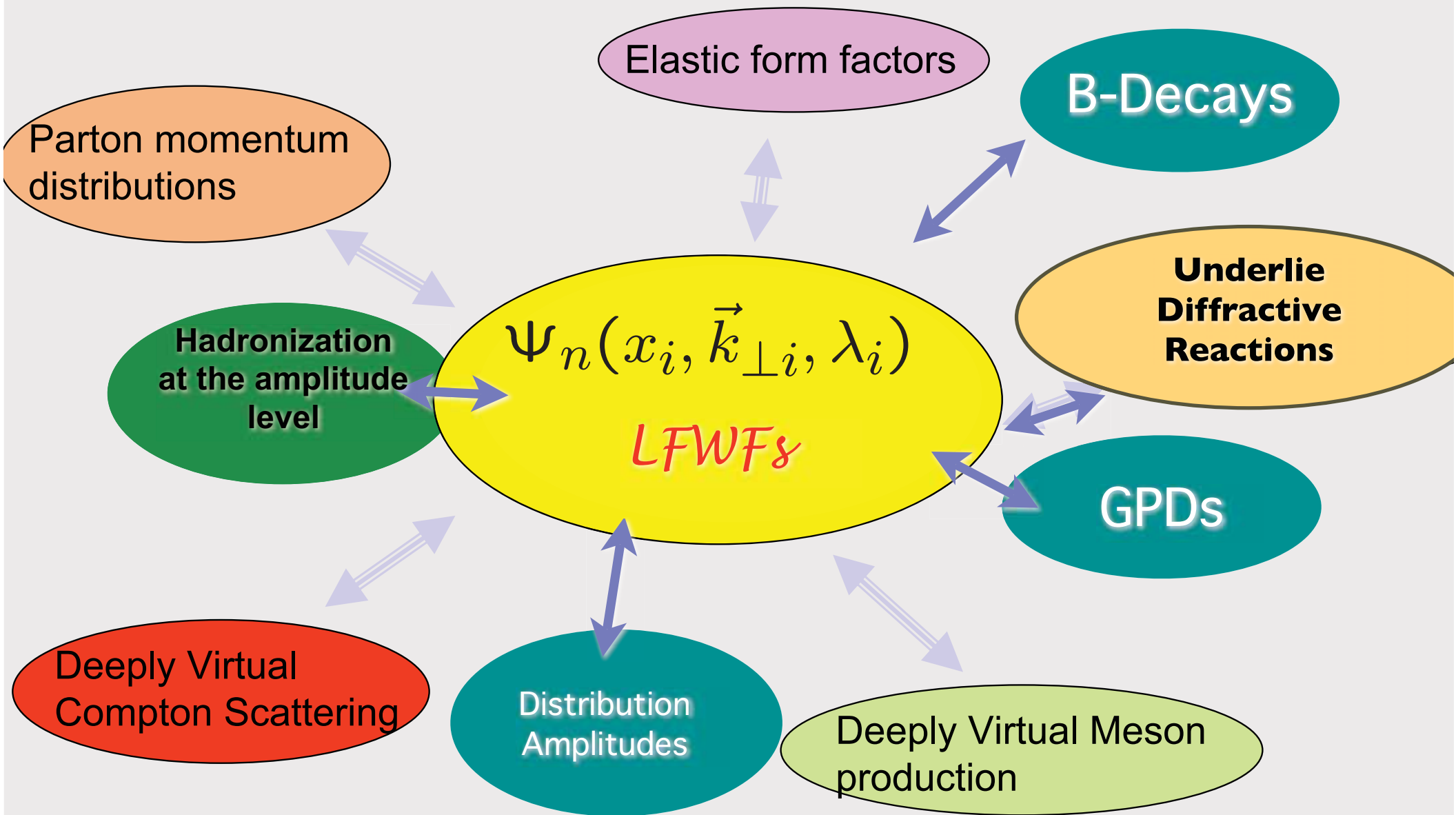
H_{LF}^{int} : Matrix in Fock Space

$$H_{LF}^{QCD} |\Psi_h\rangle = \mathcal{M}_h^2 |\Psi_h\rangle$$

Eigenvalues and Eigensolutions give Hadron Spectrum and Light-Front wavefunctions



A Unified Description of Hadron Structure



$$|p, S_z\rangle = \sum_{n=3} \Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i) |n; \vec{k}_{\perp i}, \lambda_i\rangle$$

sum over states with $n=3, 4, \dots$ constituents

The Light Front Fock State Wavefunctions

$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$

are boost invariant; they are independent of the hadron's energy and momentum P^μ .

The light-cone momentum fraction

$$x_i = \frac{k_i^+}{p^+} = \frac{k_i^0 + k_i^z}{P^0 + P^z}$$

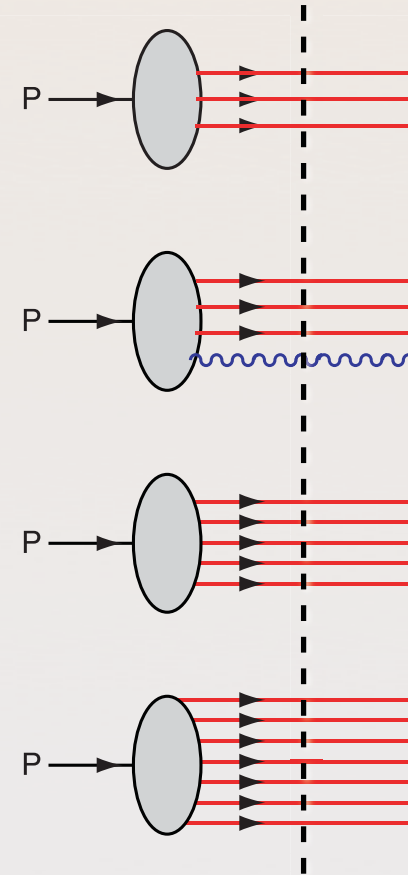
are boost invariant.

$$\sum_i^n k_i^+ = P^+, \quad \sum_i^n x_i = 1, \quad \sum_i^n \vec{k}_{i\perp} = \vec{0}^\perp.$$

Intrinsic heavy quarks,

$$\bar{s}(x) \neq s(x)$$

$$\bar{u}(x) \neq \bar{d}(x)$$



Soft gluons in the infinite momentum wave function and the BFKL pomeron.

[Alfred H. Mueller](#) ([SLAC](#) & [Columbia U.](#)) . SLAC-PUB-10047, CU-TP-609, Aug 1993. 12pp.

Published in **Nucl.Phys.B415:373-385,1994.**

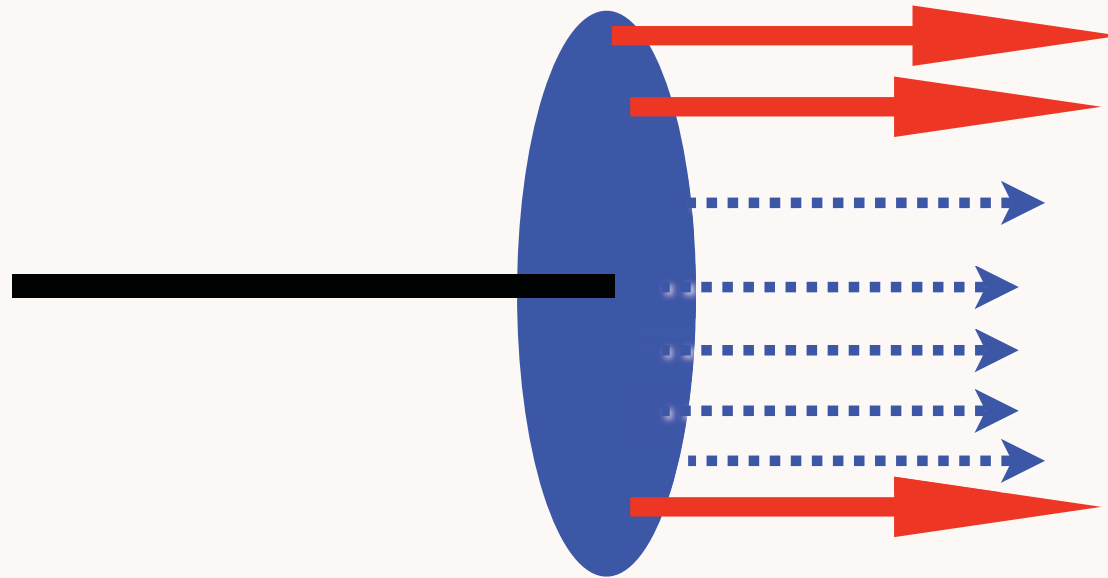
Light cone wave functions at small x.

[F. Antonuccio](#) ([Heidelberg, Max Planck Inst.](#) & [Heidelberg U.](#)) , [S.J. Brodsky](#) ([SLAC](#)) , [S. Dalley](#) ([CERN](#)) .

Phys.Lett.B412:104-110,1997.

e-Print: [hep-ph/9705413](#)

Mueller: BFKL derived from multi-gluon Fock State



Antonuccio, Dalley, sjb: Ladder Relations

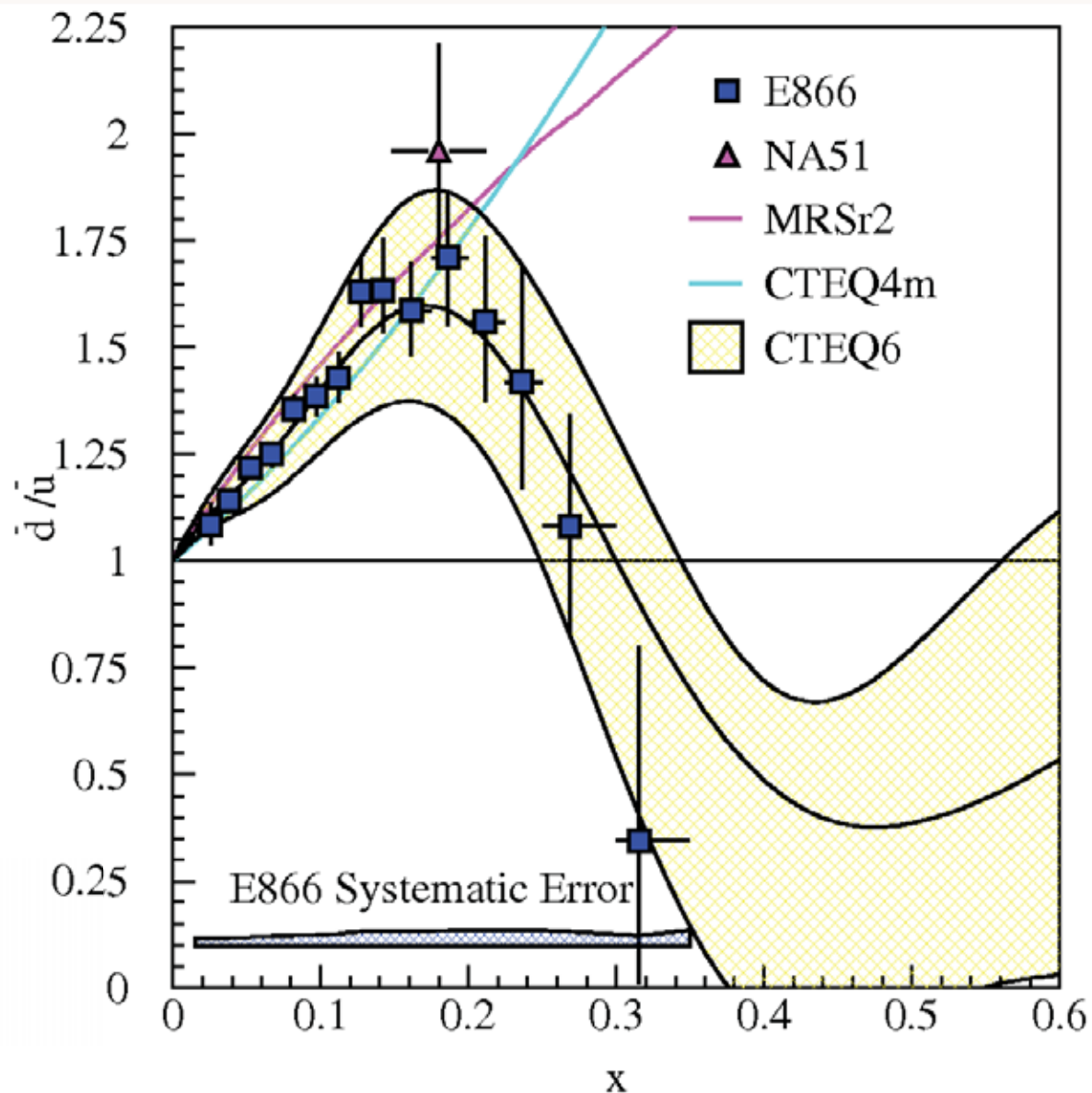
- E866/NuSea (Drell-Yan)

$$\bar{d}(x) \neq \bar{u}(x)$$

$$s(x) \neq \bar{s}(x)$$

*Intrinsic glue, sea,
heavy quarks*

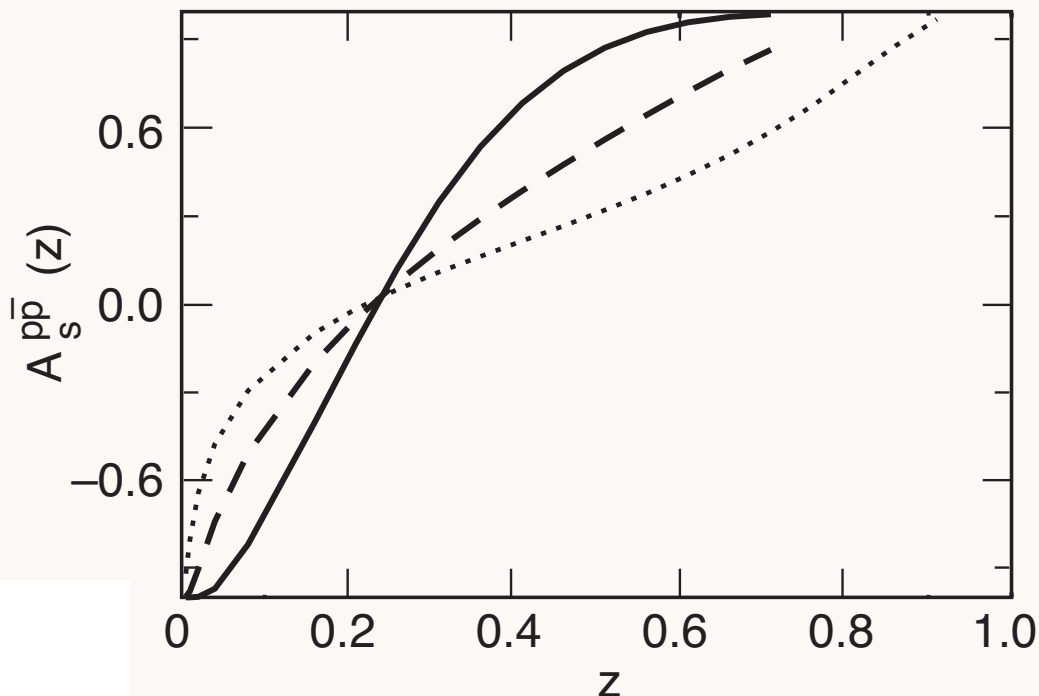
$\bar{d}(x)/\bar{u}(x)$ for $0.015 \leq x \leq 0.35$



Compare protons versus anti-proton in \bar{s} current quark fragmentation

$$D_{s \rightarrow p}(z) \neq D_{s \rightarrow \bar{p}}(z)$$

Tag s quark via high x_F Λ production in proton fragmentation region.



B.Q. Ma and sjb

$$A_s^{p\bar{p}}(z) = \frac{D_{s \rightarrow p}(z) - D_{s \rightarrow \bar{p}}(z)}{D_{s \rightarrow p}(z) + D_{s \rightarrow \bar{p}}(z)}$$

Consequence of $s_p(x) \neq \bar{s}_p(x)$

$$|uuds\bar{s}\rangle \simeq |K^+\Lambda\rangle$$

Angular Momentum on the Light-Front

$$J^z = \sum_{i=1}^n s_i^z + \sum_{j=1}^{n-1} l_j^z.$$

Conserved
LF Fock state by Fock State

$$l_j^z = -i \left(k_j^1 \frac{\partial}{\partial k_j^2} - k_j^2 \frac{\partial}{\partial k_j^1} \right) \quad n-1 \text{ orbital angular momenta}$$

Nonzero Anomalous Moment --> Nonzero orbital angular momentum

$$\frac{F_2(q^2)}{2M} = \sum_a \int [dx][d^2\mathbf{k}_\perp] \sum_j e_j \frac{1}{2} \times$$

$$\left[-\frac{1}{q^L} \psi_a^{\uparrow*}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \psi_a^\downarrow(x_i, \mathbf{k}_{\perp i}, \lambda_i) + \frac{1}{q^R} \psi_a^{\downarrow*}(x_i, \mathbf{k}'_{\perp i}, \lambda_i) \psi_a^\uparrow(x_i, \mathbf{k}_{\perp i}, \lambda_i) \right]$$

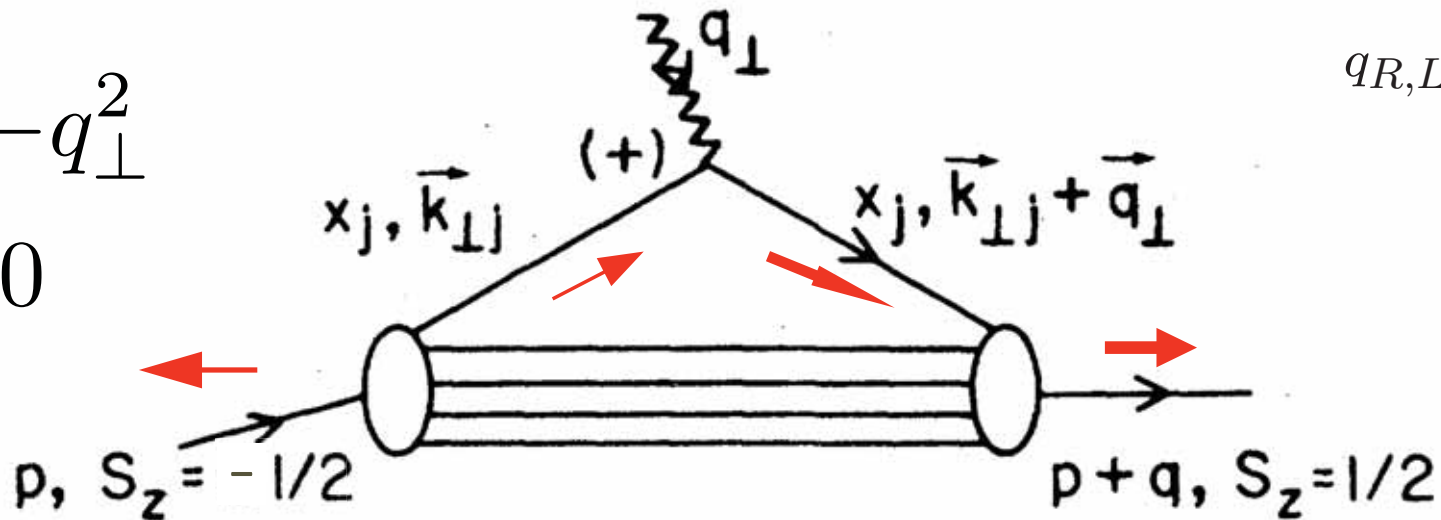
$$\mathbf{k}'_{\perp i} = \mathbf{k}_{\perp i} - x_i \mathbf{q}_\perp$$

$$\mathbf{k}'_{\perp j} = \mathbf{k}_{\perp j} + (1 - x_j) \mathbf{q}_\perp$$

$$q^2 = -q_\perp^2$$

$$q^+ = 0$$

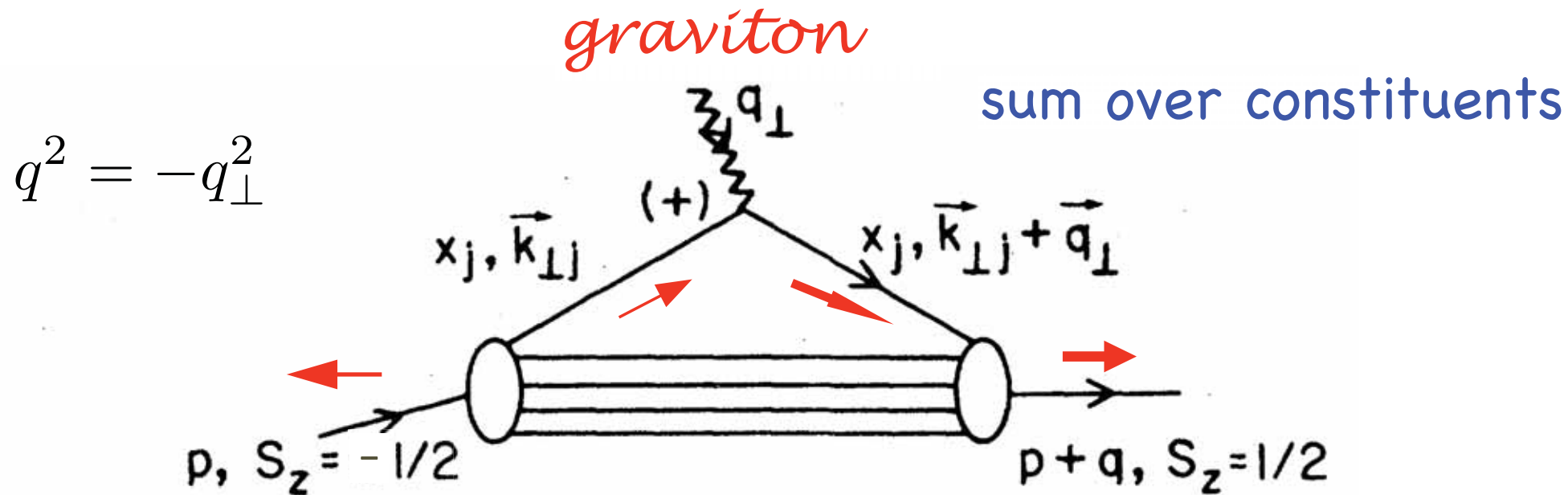
$$q_{R,L} = q^x \pm iq^y$$



Must have $\Delta l_z = \pm 1$ to have nonzero $F_2(q^2)$

Anomalous gravitomagnetic moment $B(0)$

Okun, Kobzarev, Teryaev: $B(0)$ Must vanish because of Equivalence Theorem



Hwang, Schmidt, sjb;
Holstein et al

$$B(0) = 0$$

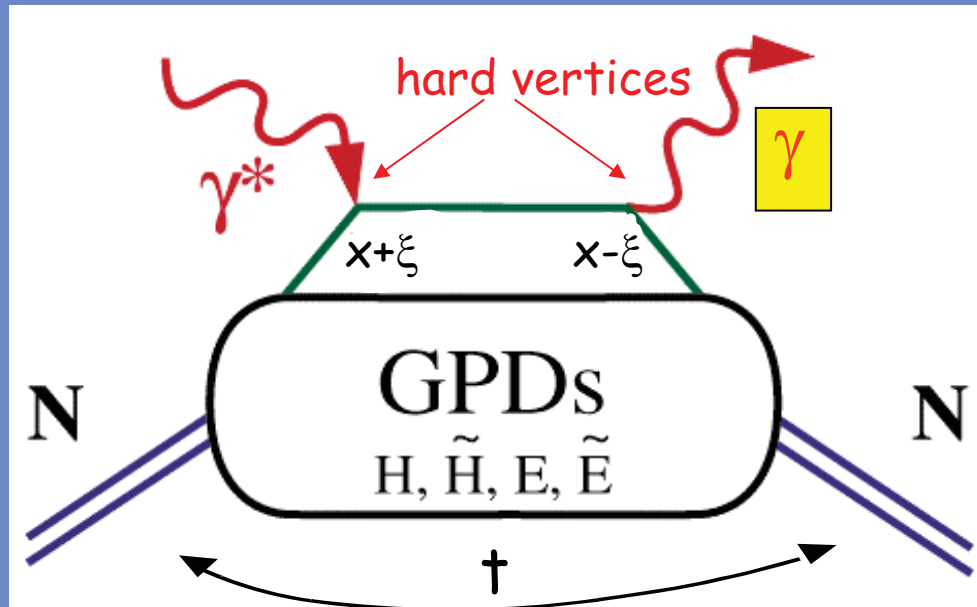
Each Fock State

Some Applications of Light-Front Wavefunctions

- Exact formulae for form factors, quark and gluon distributions; vanishing anomalous gravitational moment; edm connection to anm
- Deeply Virtual Compton Scattering, generalized parton distributions, angular momentum sum rules
- Exclusive weak decay amplitudes
- Single spin asymmetries: Role of ISI and FSI
- Factorization theorems, DGLAP, BFKL, ERBL Evolution
- Quark interchange amplitude
- Relation of spin, momentum, and other distributions to physics of the hadron itself.

GPDs & Deeply Virtual Exclusive Processes - New Insight into Nucleon Structure

Deeply Virtual Compton Scattering (DVCS)



x - quark momentum fraction

ξ - longitudinal momentum transfer

$\sqrt{-t}$ - Fourier conjugate to transverse impact parameter

$H(x, \xi, t), E(x, \xi, t), \dots$ "Generalized Parton Distributions"

Quark angular momentum (Ji sum rule)

$$J^q = \frac{1}{2} - J^G = \frac{1}{2} \int_{-1}^1 x dx [H^q(x, \xi, 0) + E^q(x, \xi, 0)]$$

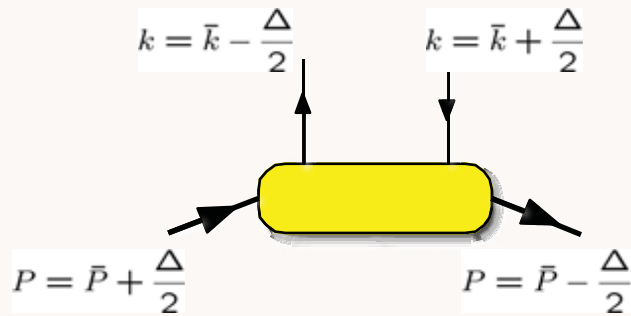
X. Ji, Phys.Rev.Lett.78,610(1997)

Light-Front Wave Function Overlap Representation

DVCS/GPD

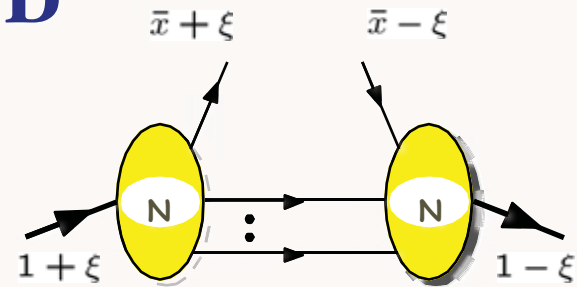
Diehl, Hwang, sjb, NPB596, 2001

See also: Diehl, Feldmann, Jakob, Kroll



$\xi < \bar{x} < 1$

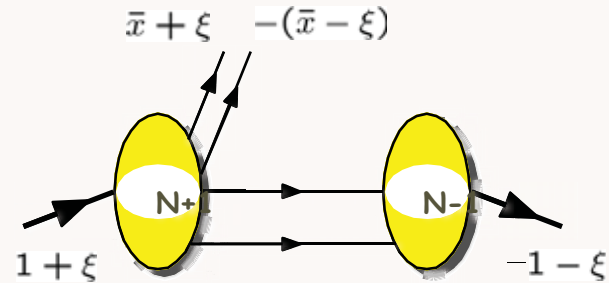
$$\sum_N$$



DGLAP region

$-\xi < \bar{x} < \xi$

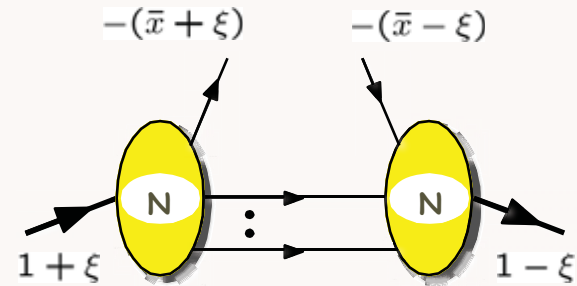
$$\sum_N$$



ERBL region

$-1 < \bar{x} < -\xi$

$$\sum_N$$



DGLAP region

$N=3$ VALENCE QUARK \Rightarrow Light-cone Constituent quark model

$N=5$ VALENCE QUARK + QUARK SEA \Rightarrow Meson-Cloud model

Link to DIS and Elastic Form Factors

DIS at $\xi=t=0$

$$H^q(x,0,0) = q(x), \quad -\bar{q}(-x)$$

$$\tilde{H}^q(x,0,0) = \Delta q(x), \quad \Delta \bar{q}(-x)$$

Form factors (sum rules)

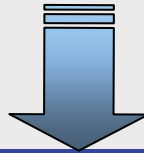
$$\int_{-1}^1 dx \sum_q [H^q(x, \xi, t)] = F_1(t) \text{ Dirac f.f.}$$

$$\int_{-1}^1 dx \sum_q [E^q(x, \xi, t)] = F_2(t) \text{ Pauli f.f.}$$

$$\int_{-1}^1 dx \tilde{H}^q(x, \xi, t) = G_{A,q}(t), \quad \int_{-1}^1 dx \tilde{E}^q(x, \xi, t) = G_{P,q}(t)$$



$$H^q, E^q, \tilde{H}^q, \tilde{E}^q(x, \xi, t)$$



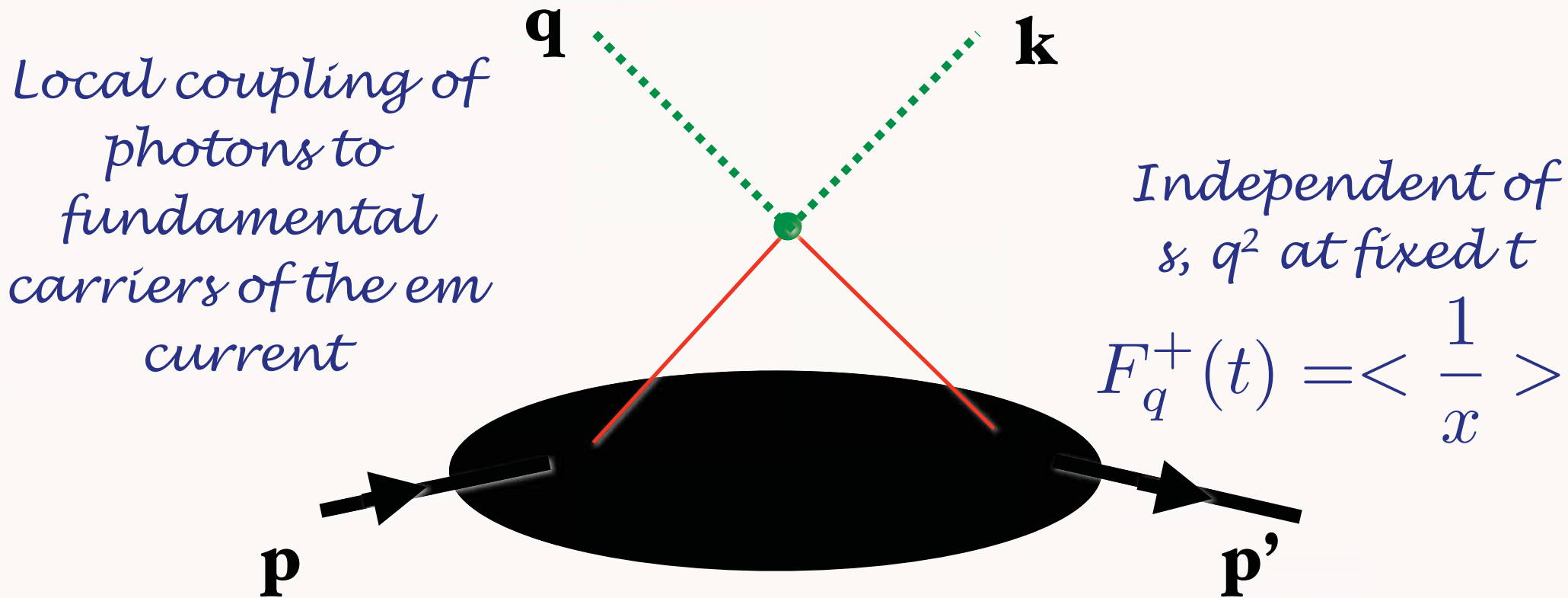
Verified using
LFWFs
Diehl, Hwang, sjb

Quark angular momentum (Ji's sum rule)

$$J^q = \frac{1}{2} - J^G = \frac{1}{2} \int_{-1}^1 x dx [H^q(x, \xi, 0) + E^q(x, \xi, 0)]$$

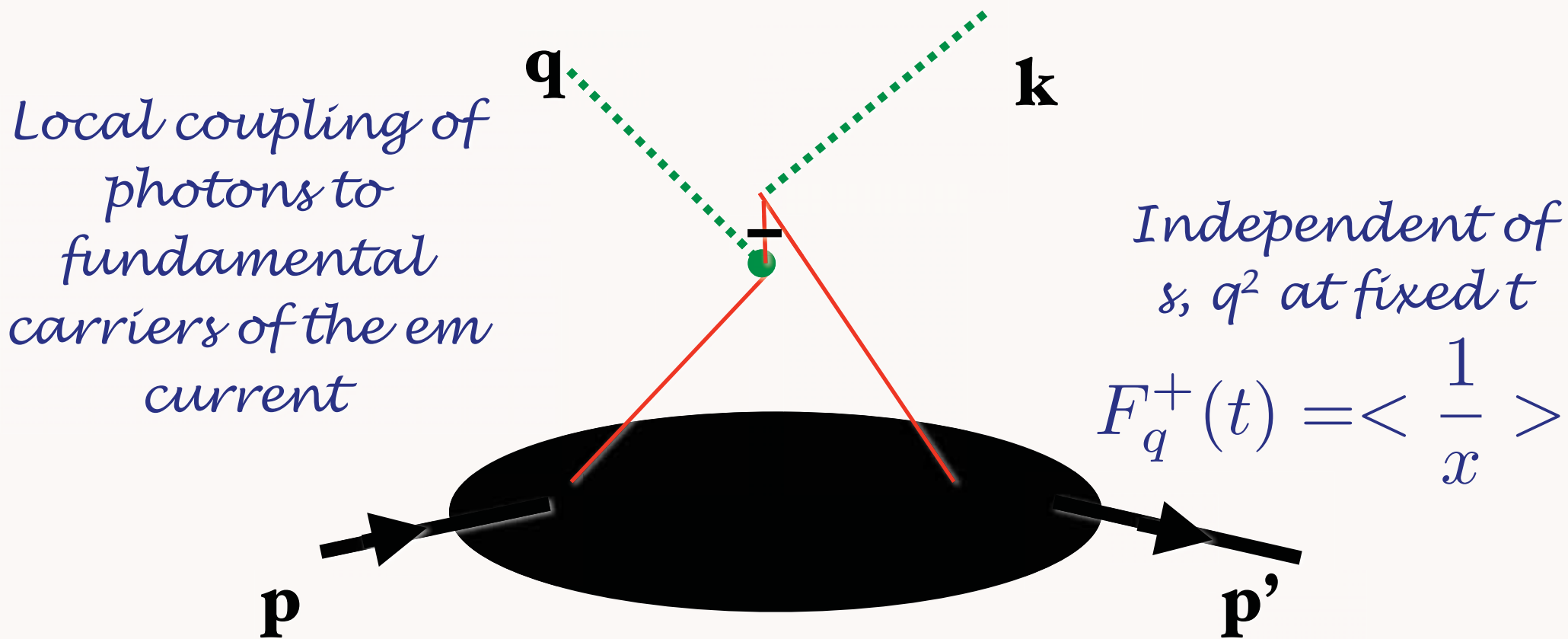
X. Ji, Phys.Rev.Lett.78,610(1997)

'Seagull' contribution to real and virtual Compton scattering

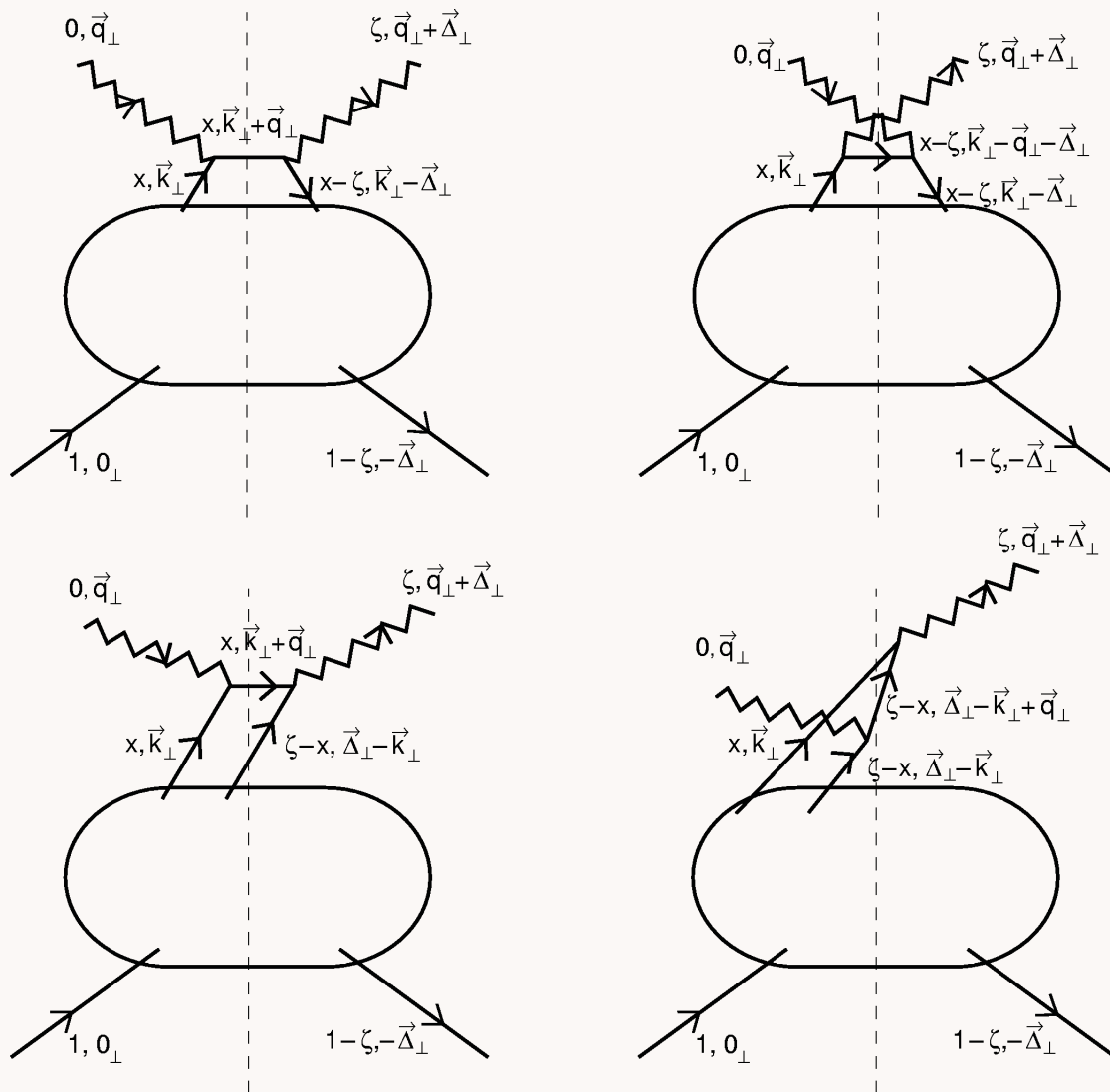


$$M = -2 \sum_{q/p} e_q^2 F_q^+(t) \vec{\epsilon} \cdot \vec{\epsilon}'$$

Instantaneous fermion exchange contribution to real and virtual Compton scattering



$$M = -2 \sum_{q/p} e_q^2 F_q^+(t) \vec{\epsilon} \cdot \vec{\epsilon}'$$



$$A_{J=0} \sim e_q^2 s^0 F(t)$$

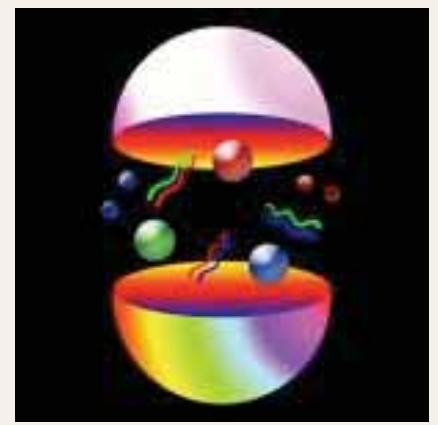
*Local J=0
fixed pole
contribution*

Close, Gunion, sjb;
Szczepaniak, Llanes-
Estrada, sjb

Light-cone wavefunction representation of deeply
virtual Compton scattering [☆]

Stanley J. Brodsky ^a, Markus Diehl ^{a,1}, Dae Sung Hwang ^b

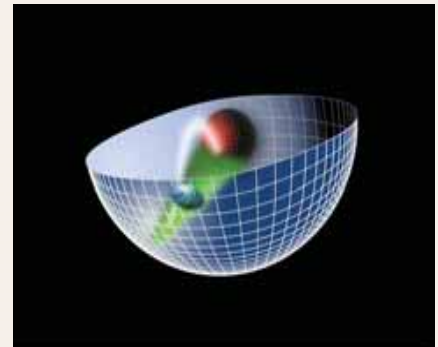
- Quarks and Gluons:
Fundamental constituents of hadrons and nuclei



- *Quantum Chromodynamics (QCD)*

- New Insights from higher space-time dimensions: *AdS/QCD*

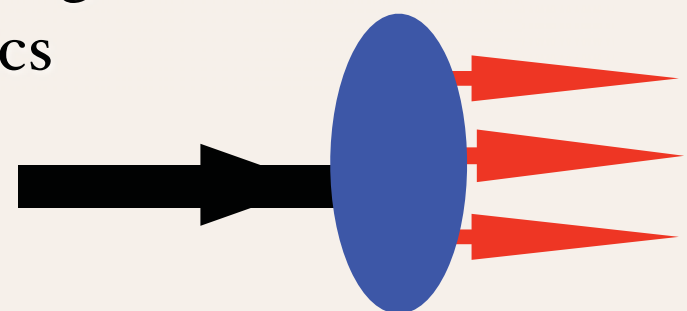
- *Light-Front Holography*



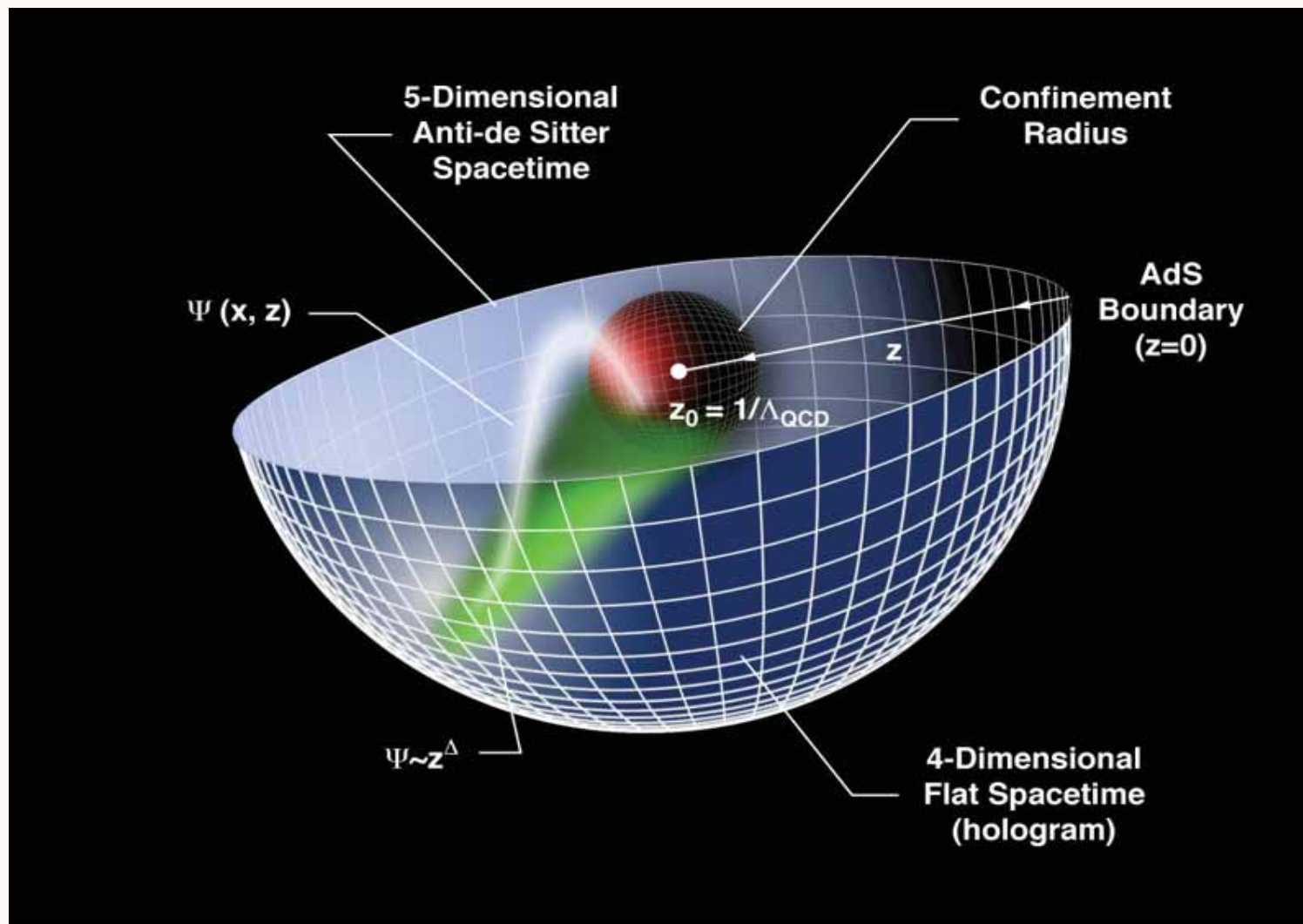
- *Hadronization at the Amplitude Level*

- *Light Front Wavefunctions:* analogous to the Schrodinger wavefunctions of atomic physics

$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$



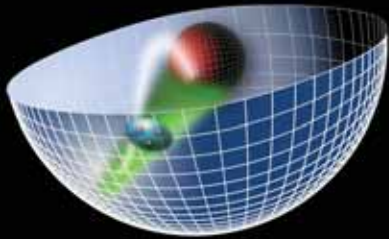
Applications of AdS/CFT to QCD



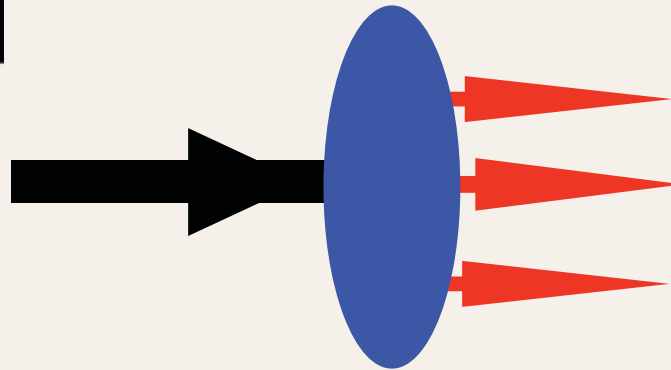
Changes in physical length scale mapped to evolution in the 5th dimension z

in collaboration with Guy de Teramond

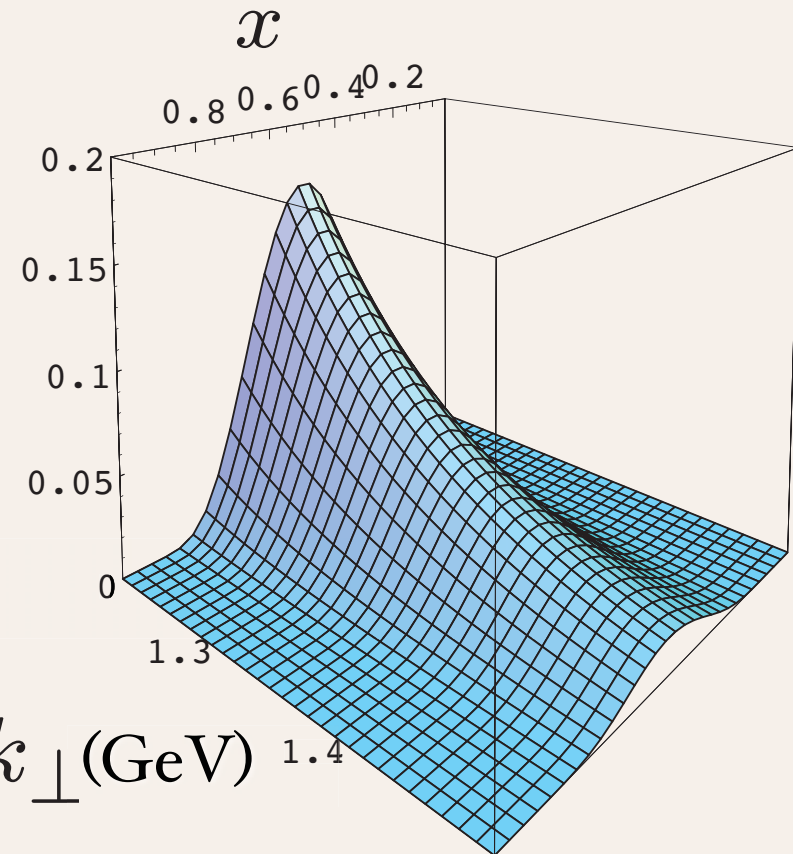
$$\phi(z)$$



- *Light-Front Holography*



$$\Psi_n(x_i, \vec{k}_{\perp i}, \lambda_i)$$



- *Light Front Wavefunctions:*

Schrödinger Wavefunctions
of Hadron Physics

$LF(3+1)$

AdS_5

$$\psi(x, \vec{b}_\perp)$$

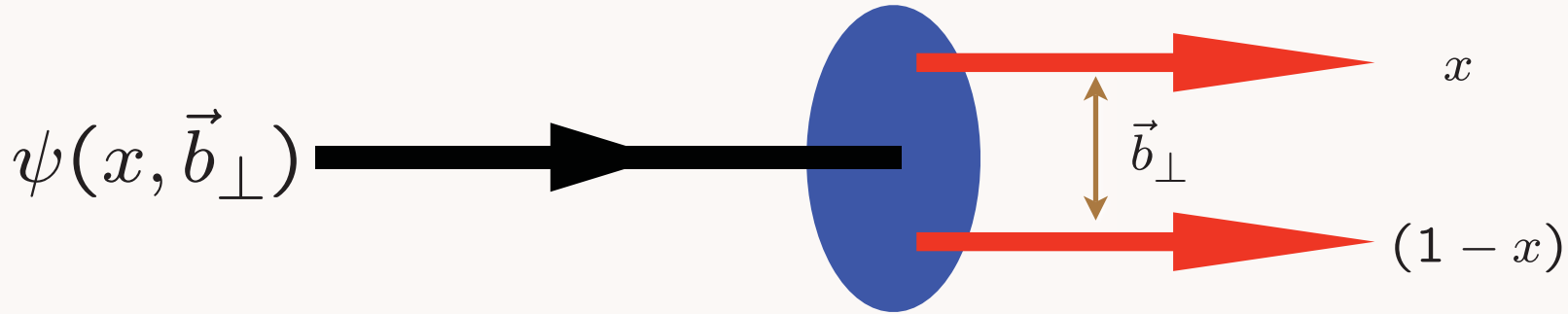


$$\phi(z)$$

$$\zeta = \sqrt{x(1-x)} \vec{b}_\perp^2$$



$$z$$



$$\psi(x, \zeta) = \sqrt{x(1-x)} \zeta^{-1/2} \phi(\zeta)$$

Light-Front Holography: Unique mapping derived from equality of LF and AdS formula for current matrix elements

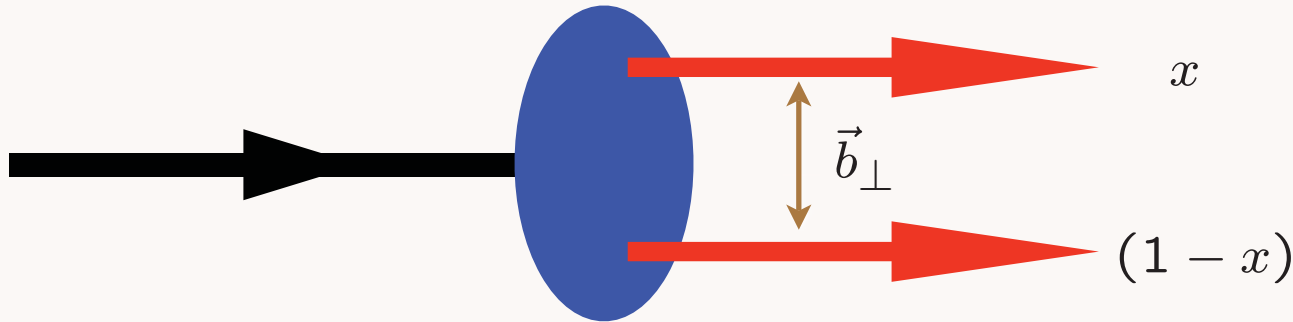
Light-Front Holography: Map AdS/CFT to 3+1 LF Theory

Relativistic LF radial equation

Frame Independent

$$\left[-\frac{d^2}{d\zeta^2} + \frac{1 - 4L^2}{4\zeta^2} + U(\zeta) \right] \phi(\zeta) = \mathcal{M}^2 \phi(\zeta)$$

$$\zeta^2 = x(1-x)\mathbf{b}_\perp^2.$$



G. de Teramond, sjb

$$U(\zeta) = \kappa^4 \zeta^2$$

*soft wall
confining potential:*

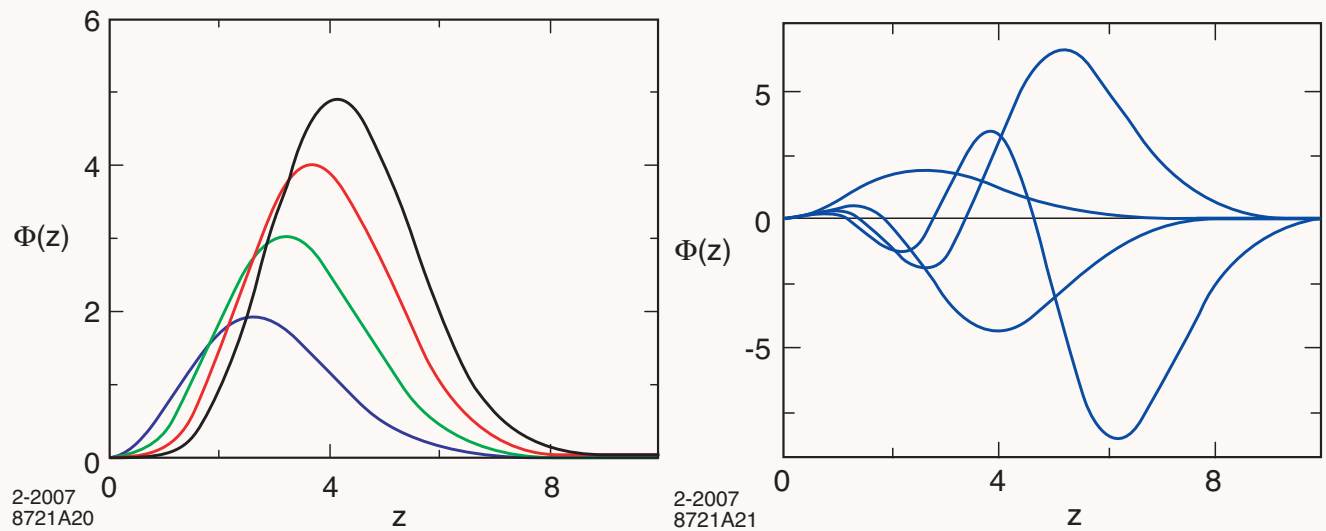
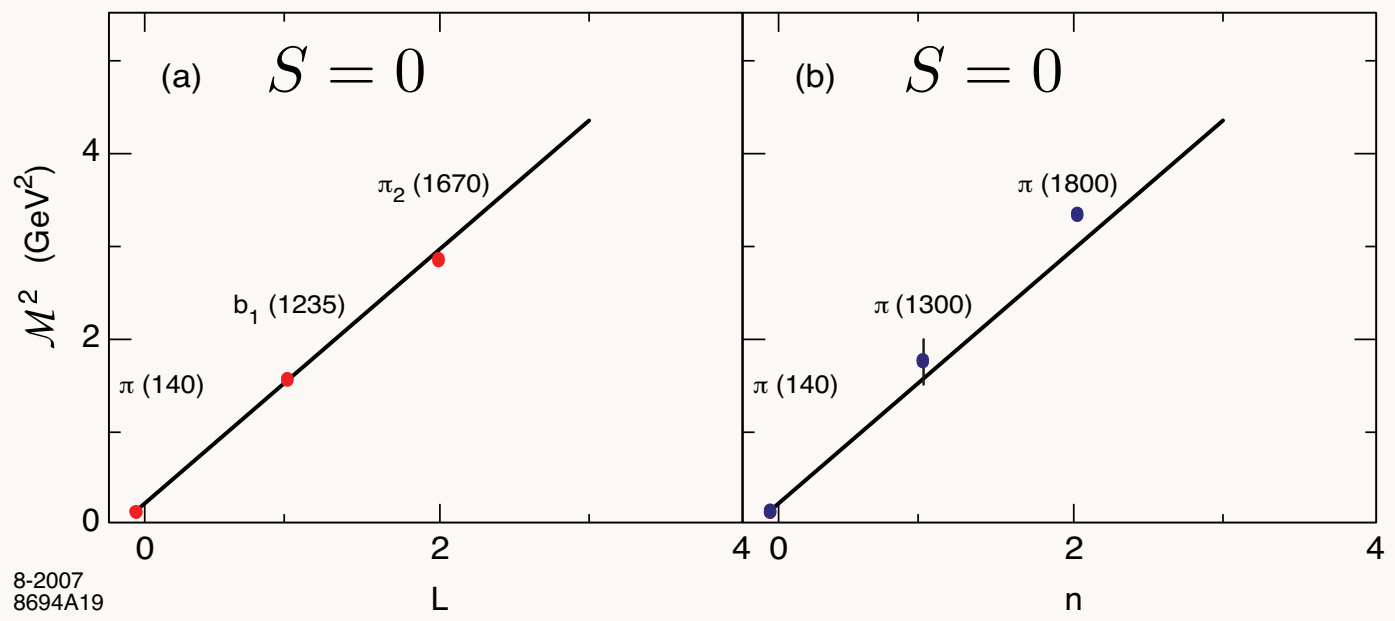


Fig: Orbital and radial AdS modes in the soft wall model for $\kappa = 0.6$ GeV .

Soft Wall Model



Light meson orbital (a) and radial (b) spectrum for $\kappa = 0.6$ GeV.

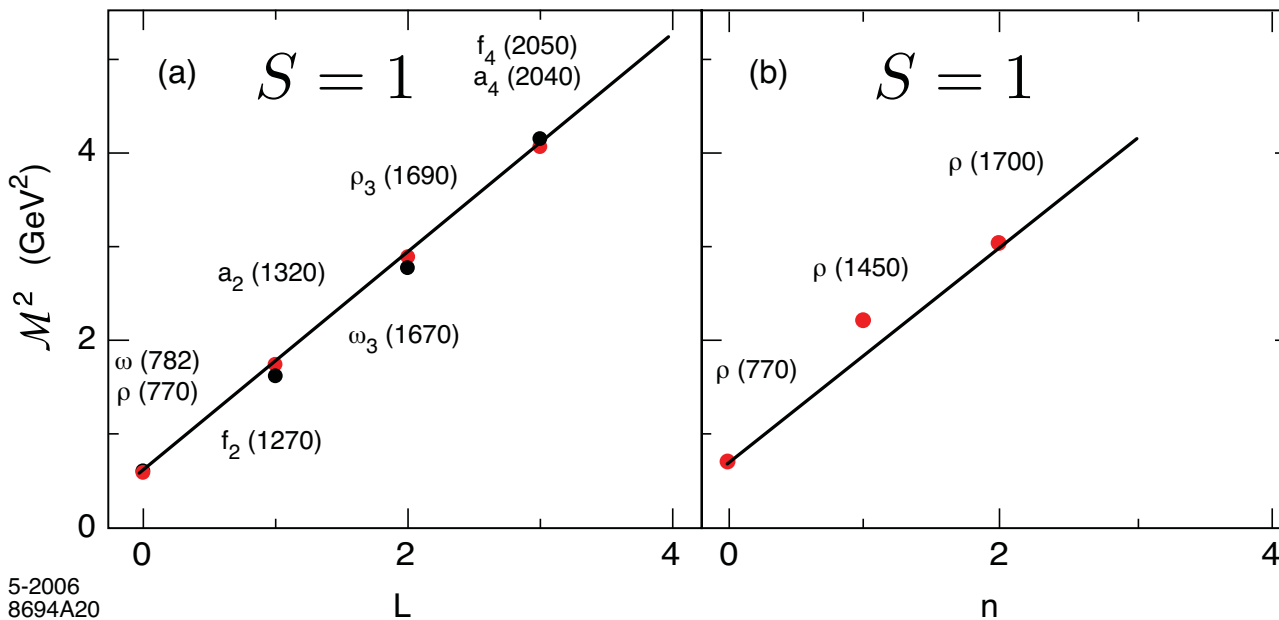
- Effective LF Schrödinger wave equation

$$\left[-\frac{d^2}{dz^2} - \frac{1 - 4L^2}{4z^2} + \kappa^4 z^2 + 2\kappa^2(L + S - 1) \right] \phi_S(z) = \mathcal{M}^2 \phi_S(z)$$

with eigenvalues $\mathcal{M}^2 = 2\kappa^2(2n + 2L + S)$.

Same slope in n and L

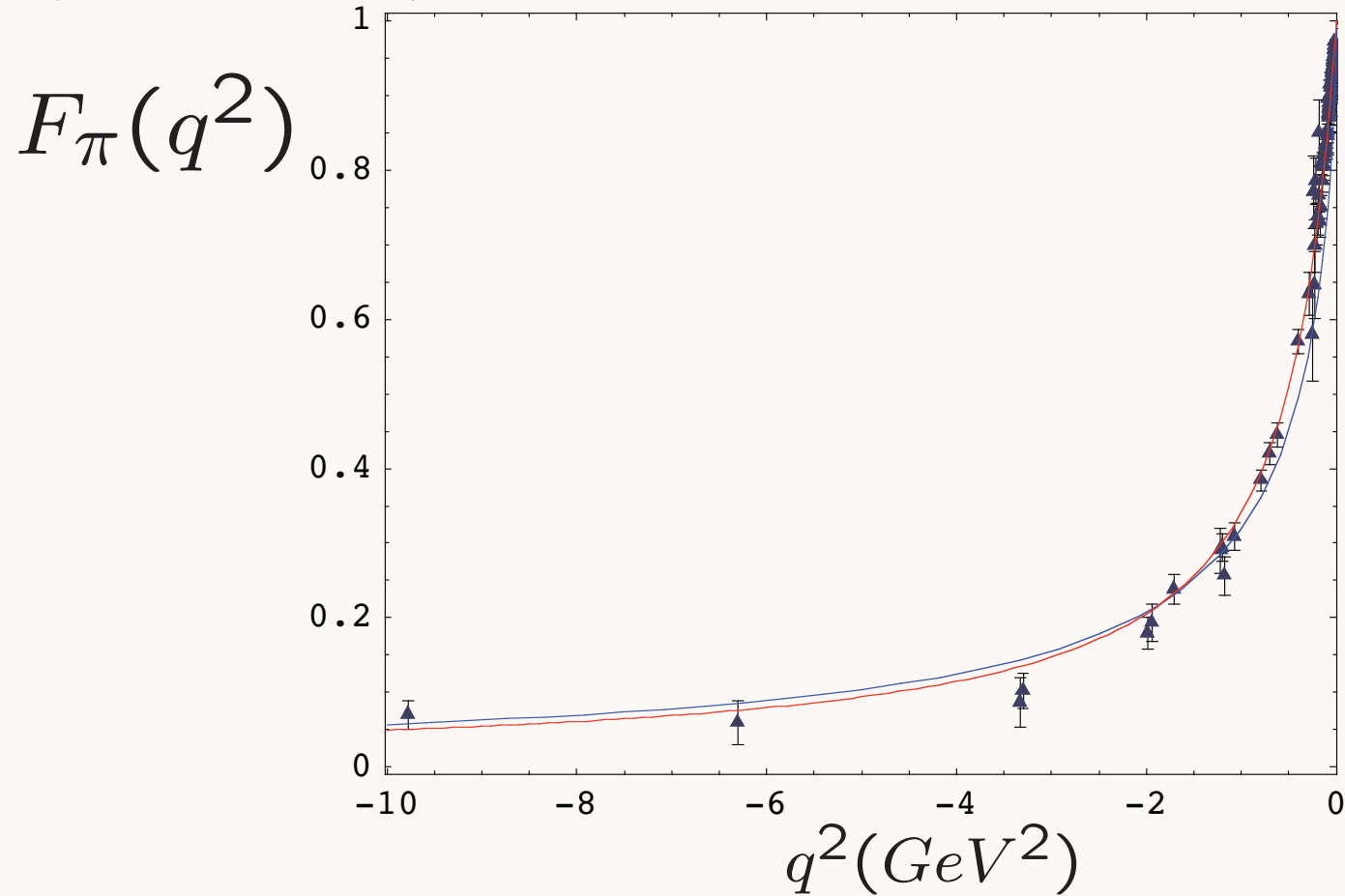
- Compare with Nambu string result (rotating flux tube): $M_n^2(L) = 2\pi\sigma(n + L + 1/2)$.



Vector mesons orbital (a) and radial (b) spectrum for $\kappa = 0.54$ GeV.

- Glueballs in the bottom-up approach: (HW) Boschi-Filho, Braga and Carrion (2005); (SW) Colangelo, De Facio, Jugeau and Nicotri(2007).

Spacelike pion form factor from AdS/CFT



Data Compilation
Baldini, Kloe and Volmer

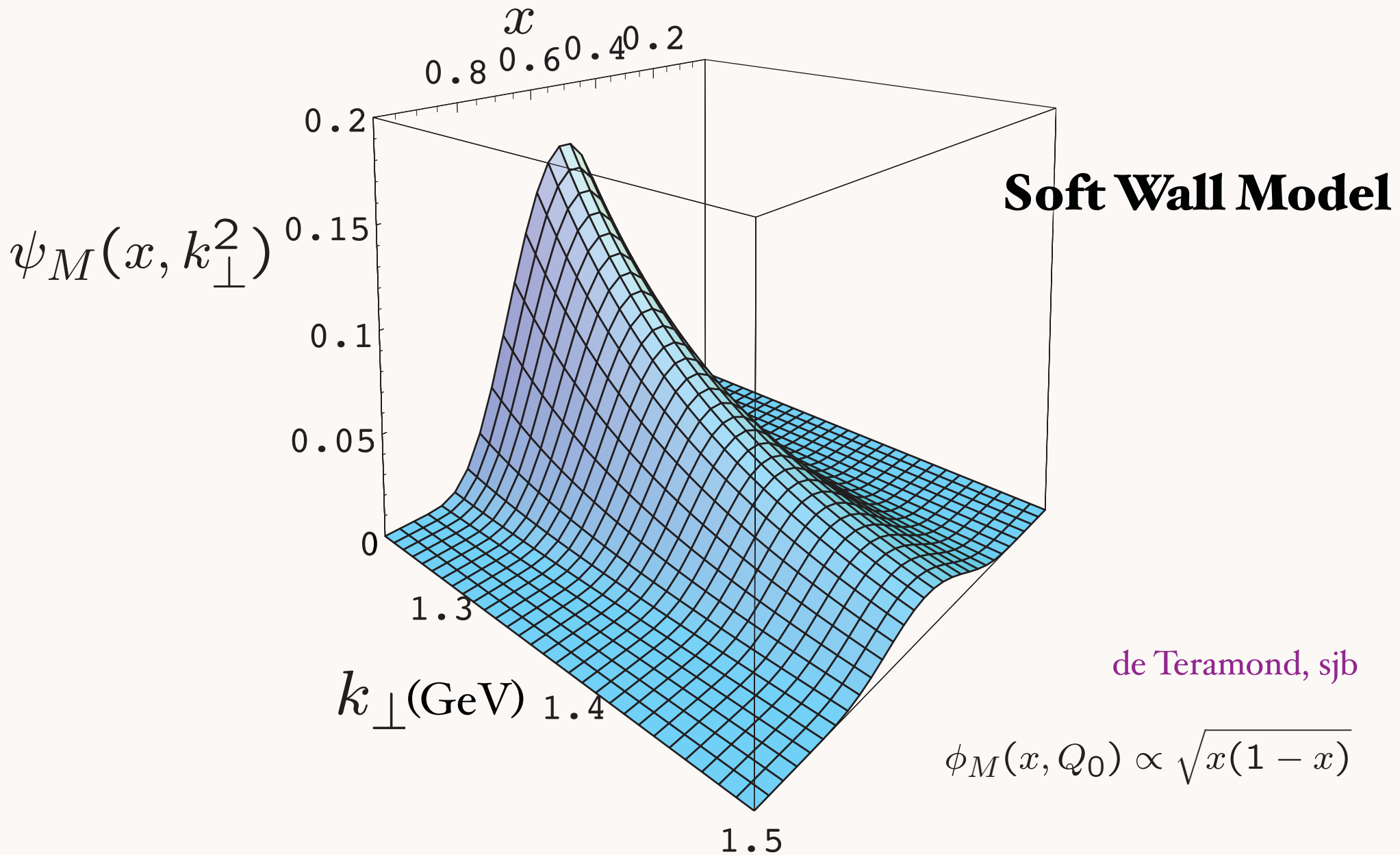
— Soft Wall: Harmonic Oscillator Confinement

— Hard Wall: Truncated Space Confinement

One parameter - set by pion decay constant

de Teramond, sjb
See also: Radyushkin

Prediction from AdS/CFT: Meson LFWF



Increases PQCD prediction for $F_{\pi}(Q^2)$ by 16/9

Second Moment of Pion Distribution Amplitude

$$\langle \xi^2 \rangle = \int_{-1}^1 d\xi \xi^2 \phi(\xi)$$

$$\xi = 1 - 2x$$

$$\langle \xi^2 \rangle_{\pi} = 1/5 = 0.20$$

$$\phi_{asympt} \propto x(1-x)$$

$$\langle \xi^2 \rangle_{\pi} = 1/4 = 0.25$$

$$\phi_{AdS/QCD} \propto \sqrt{x(1-x)}$$

$$\text{Lattice (I)} \quad \langle \xi^2 \rangle_{\pi} = 0.28 \pm 0.03$$

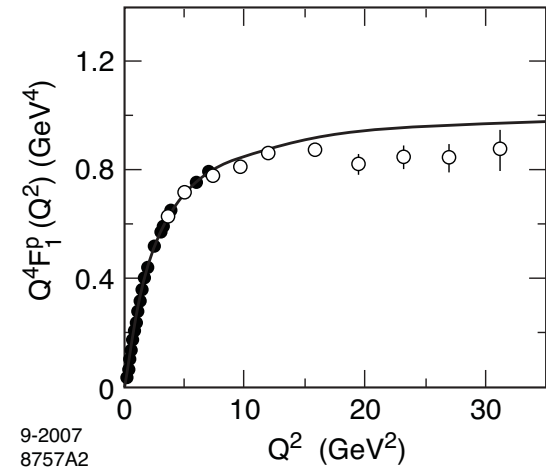
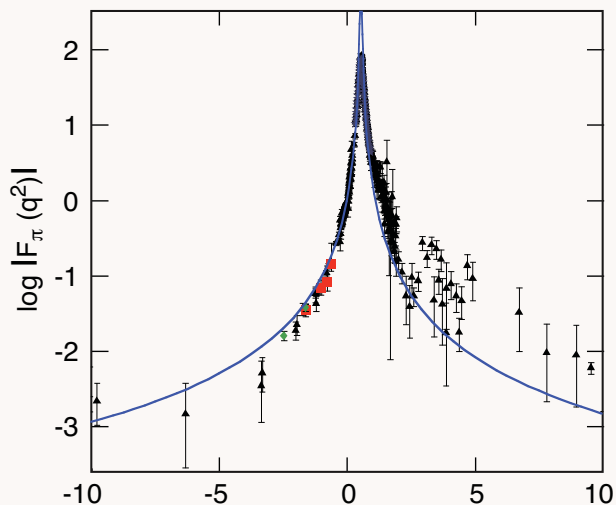
Donnellan et al.

$$\text{Lattice (II)} \quad \langle \xi^2 \rangle_{\pi} = 0.269 \pm 0.039$$

Braun et al.

Other Applications of Light-Front Holography

- Light baryon spectrum
- Light meson spectrum
- Nucleon form-factors: space-like region
- Pion form-factors: space and time-like regions
- Gravitational form factors of composite hadrons
- n -parton holographic mapping
- Heavy flavor mesons



9-2007
8757A2

hep-th/0501022
hep-ph/0602252
arXiv:0707.3859
arXiv:0802.0514
arXiv:0804.0452

Imperial College,
September 16, 2008

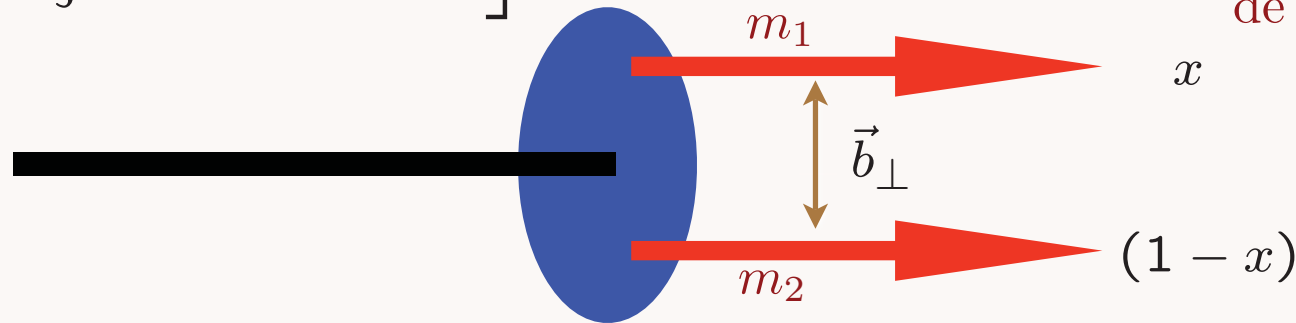
Novel QCD Physics

59

Stan Brodsky, SLAC

$$\left[-\frac{d^2}{d\zeta^2} + V(\zeta) \right] \phi(\zeta) = \mathcal{M}^2 \phi(\zeta)$$

de Teramond, sjb



$$\zeta = \sqrt{x(1-x)} \vec{b}_\perp^2$$

Holographic Variable

$$-\frac{d}{d\zeta^2} \equiv \frac{k_\perp^2}{x(1-x)}$$

LF Kinetic Energy in momentum space

Assume LFWF is a dynamical function of the quark-antiquark invariant mass squared

$$-\frac{d}{d\zeta^2} \rightarrow -\frac{d}{d\zeta^2} + \frac{m_1^2}{x} + \frac{m_2^2}{1-x} \equiv \frac{k_\perp^2 + m_1^2}{x} + \frac{k_\perp^2 + m_2^2}{1-x}$$

Result: Soft-Wall LFWF for massive constituents

$$\psi(x, \mathbf{k}_\perp) = \frac{4\pi c}{\kappa \sqrt{x(1-x)}} e^{-\frac{1}{2\kappa^2} \left(\frac{\mathbf{k}_\perp^2}{x(1-x)} + \frac{m_1^2}{x} + \frac{m_2^2}{1-x} \right)}$$

LFWF in impact space: soft-wall model with massive quarks

$$\psi(x, \mathbf{b}_\perp) = \frac{c\kappa}{\sqrt{\pi}} \sqrt{x(1-x)} e^{-\frac{1}{2}\kappa^2 x(1-x) \mathbf{b}_\perp^2 - \frac{1}{2\kappa^2} \left[\frac{m_1^2}{x} + \frac{m_2^2}{1-x} \right]}$$

$$z \rightarrow \zeta \rightarrow \chi$$

$$\chi^2 = b^2 x(1-x) + \frac{1}{\kappa^4} \left[\frac{m_1^2}{x} + \frac{m_2^2}{1-x} \right]$$

J/ψ

$\psi_{J/\psi}(x, b)$

LFWF peaks at

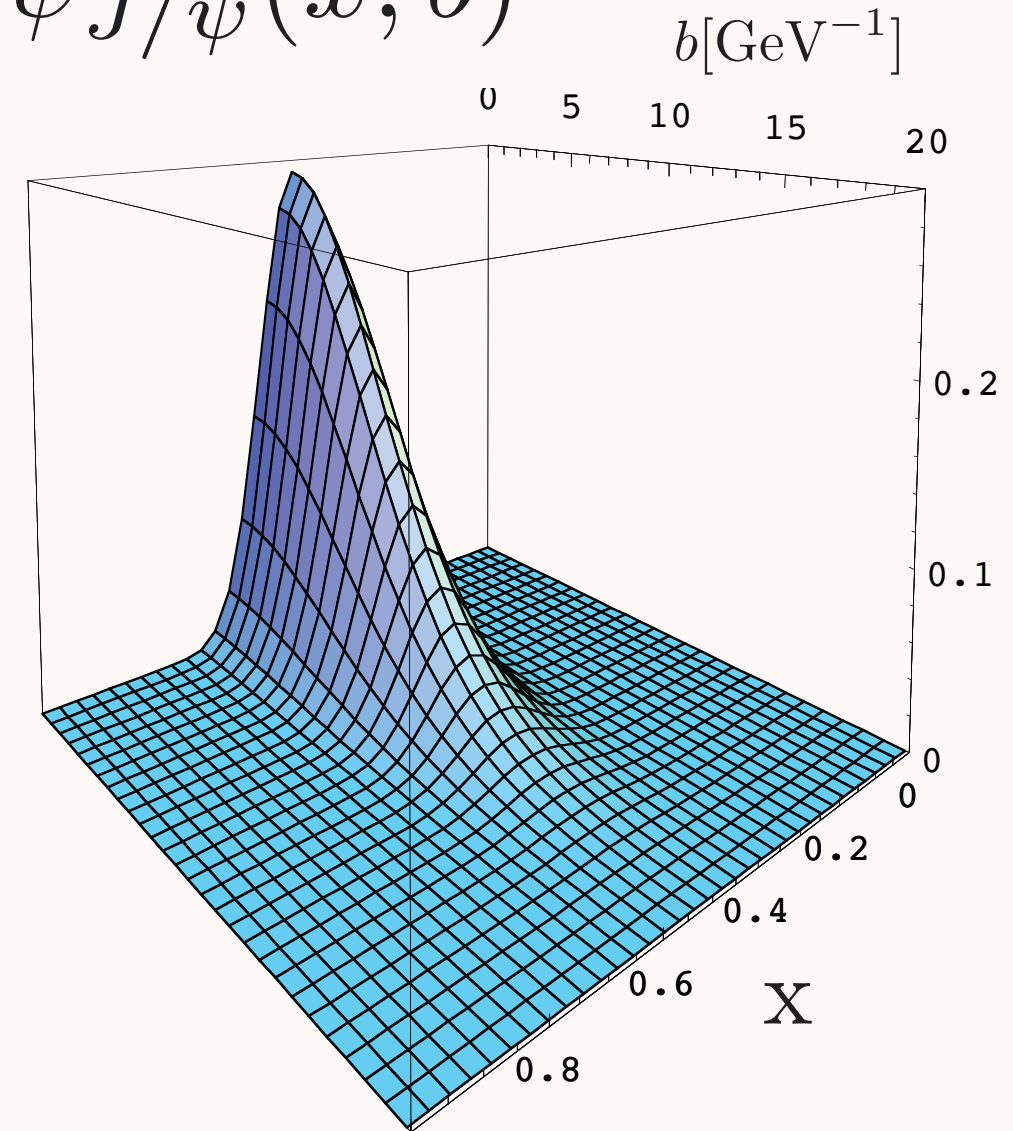
$$x_i = \frac{m_{\perp i}}{\sum_j^n m_{\perp j}}$$

where

$$m_{\perp i} = \sqrt{m^2 + k_{\perp}^2}$$

*minimum of LF
energy
denominator*

$$\kappa = 0.375 \text{ GeV}$$

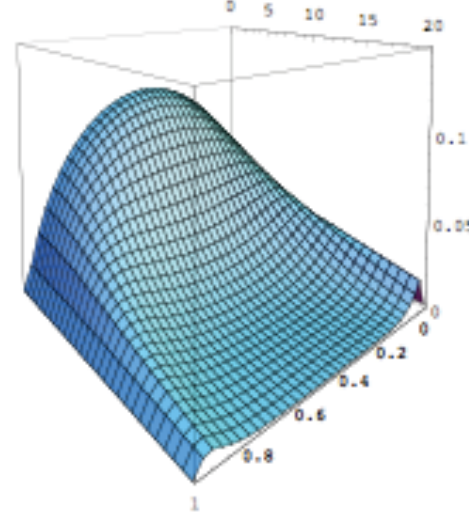
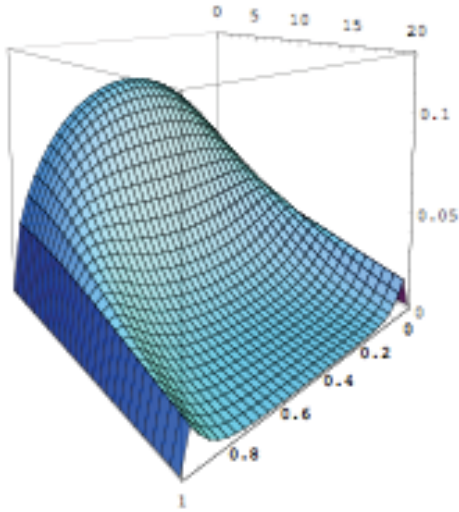


$$m_a = m_b = 1.25 \text{ GeV}$$

$$|\pi^+\rangle = |u\bar{d}\rangle$$

$$m_u = 2 \text{ MeV}$$

$$m_d = 5 \text{ MeV}$$

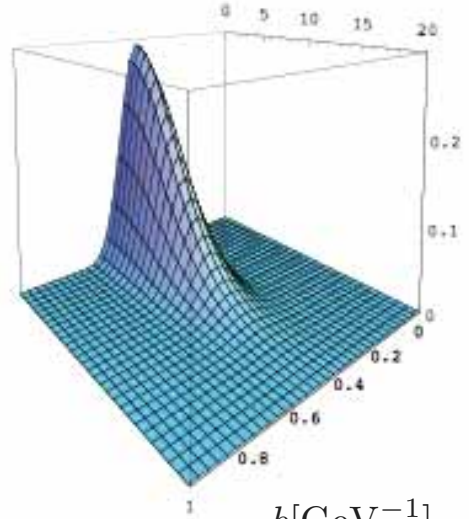
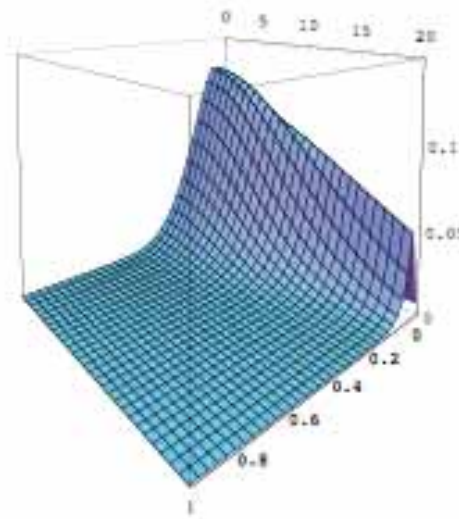


$$|K^+\rangle = |u\bar{s}\rangle$$

$$m_s = 95 \text{ MeV}$$

$$|D^+\rangle = |c\bar{d}\rangle$$

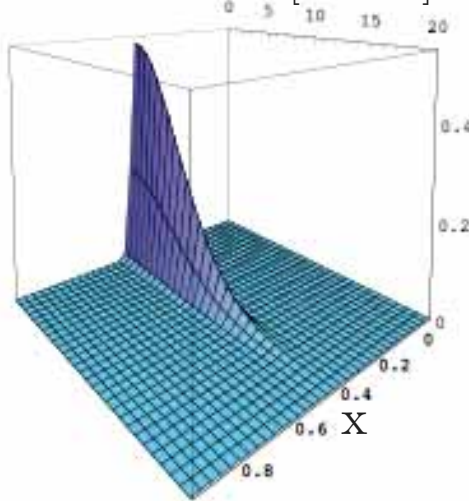
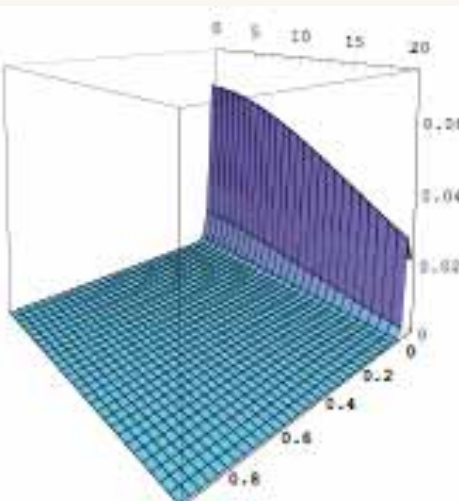
$$m_c = 1.25 \text{ GeV}$$



$$|\eta_c\rangle = |c\bar{c}\rangle$$

$$|B^+\rangle = |u\bar{b}\rangle$$

$$m_b = 4.2 \text{ GeV}$$

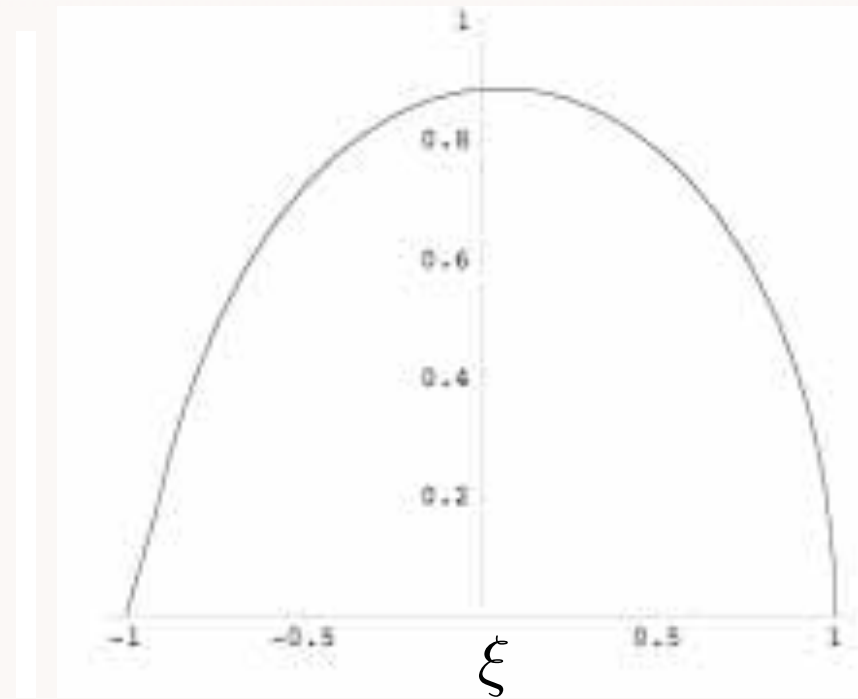


$$|\eta_b\rangle = |b\bar{b}\rangle$$

$$\kappa = 375 \text{ MeV}$$

First Moment of Kaon Distribution Amplitude

$$\langle \xi \rangle = \int_{-1}^1 d\xi \xi \phi(\xi)$$
$$\xi = 1 - 2x$$



$$\langle \xi \rangle_K = 0.04 \pm 0.002 \quad \kappa = 375 \text{ MeV}$$

Range from $m_s = 65 \pm 25 \text{ MeV}$ (PDG)

$$\langle \xi \rangle_K = 0.029 \pm 0.002$$

Donnellan et al.

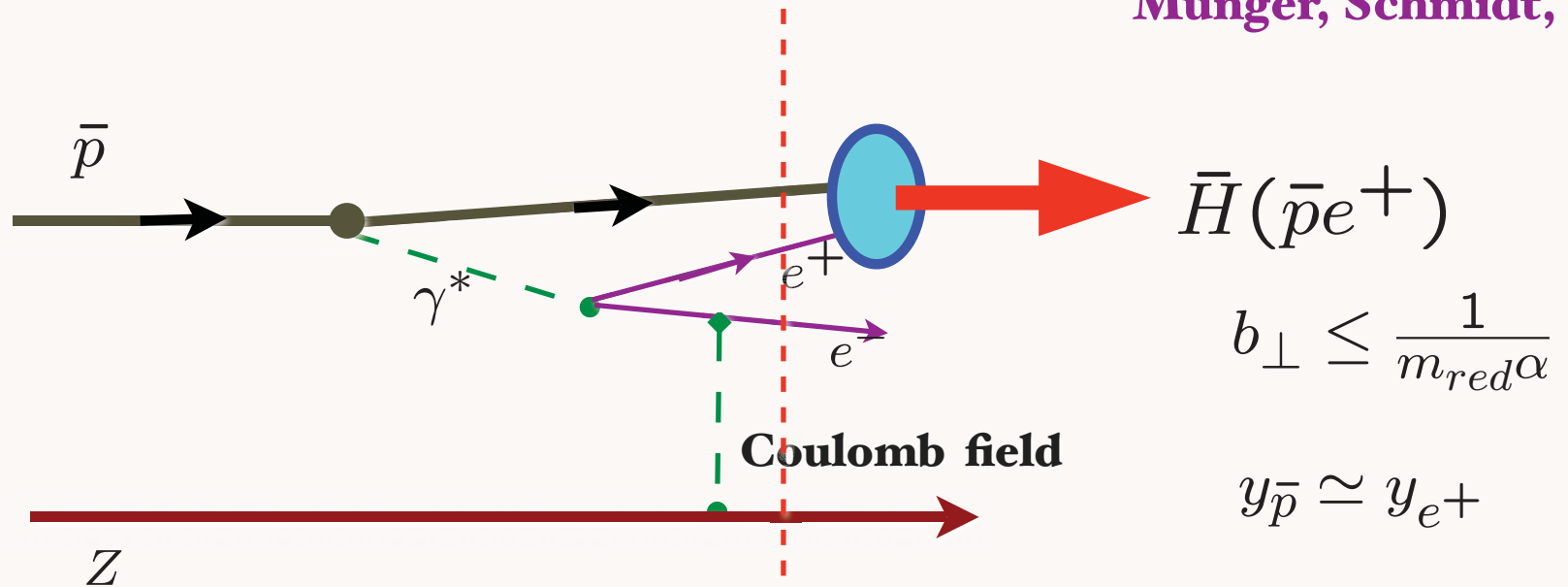
$$\langle \xi \rangle_K = 0.0272 \pm 0.0005$$

Braun et al.

Formation of Relativistic Anti-Hydrogen

Measured at CERN-LEAR and FermiLab

Munger, Schmidt, sjb

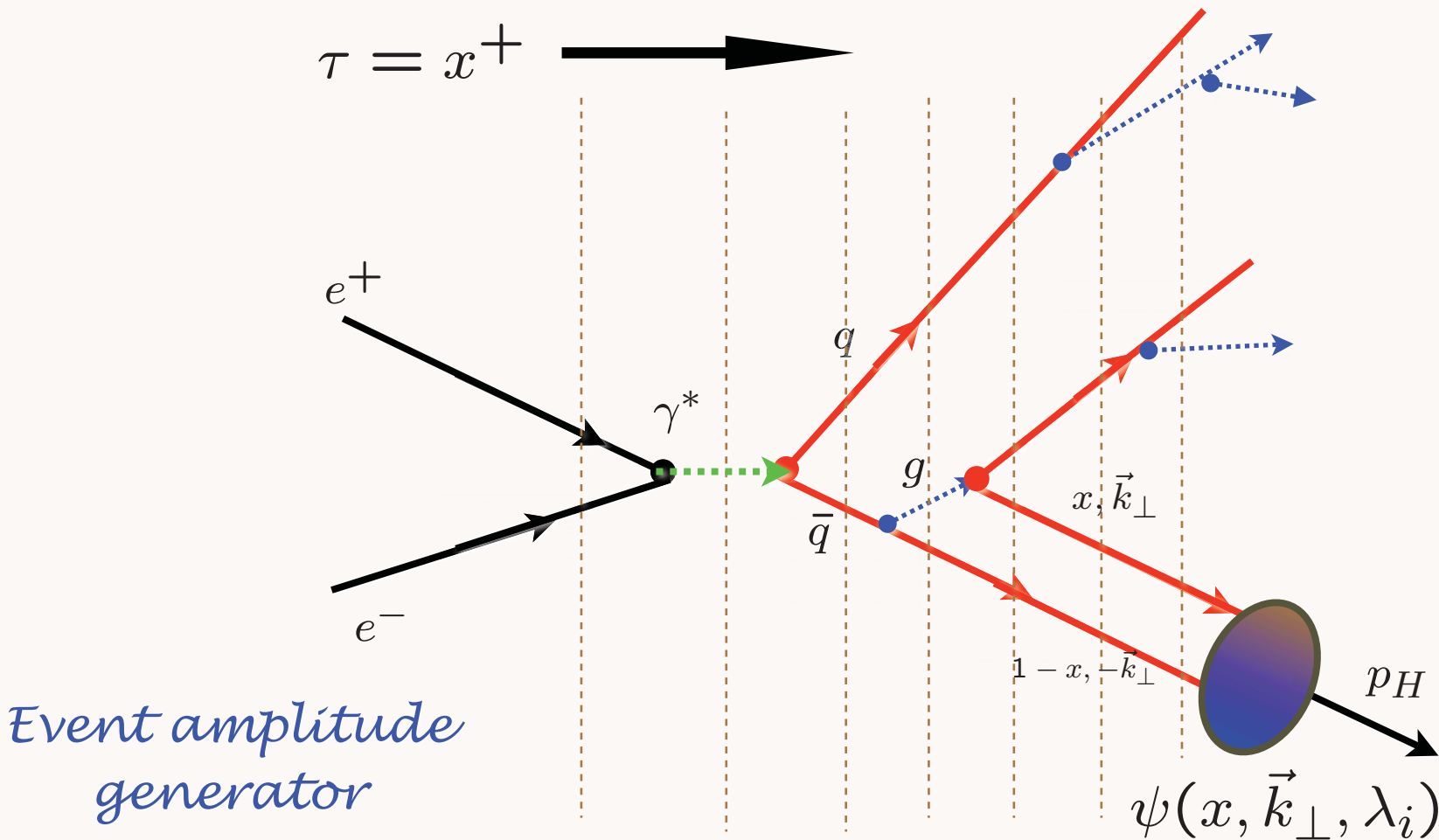


Coalescence of off-shell co-moving positron and antiproton

Wavefunction maximal at small impact separation and equal rapidity

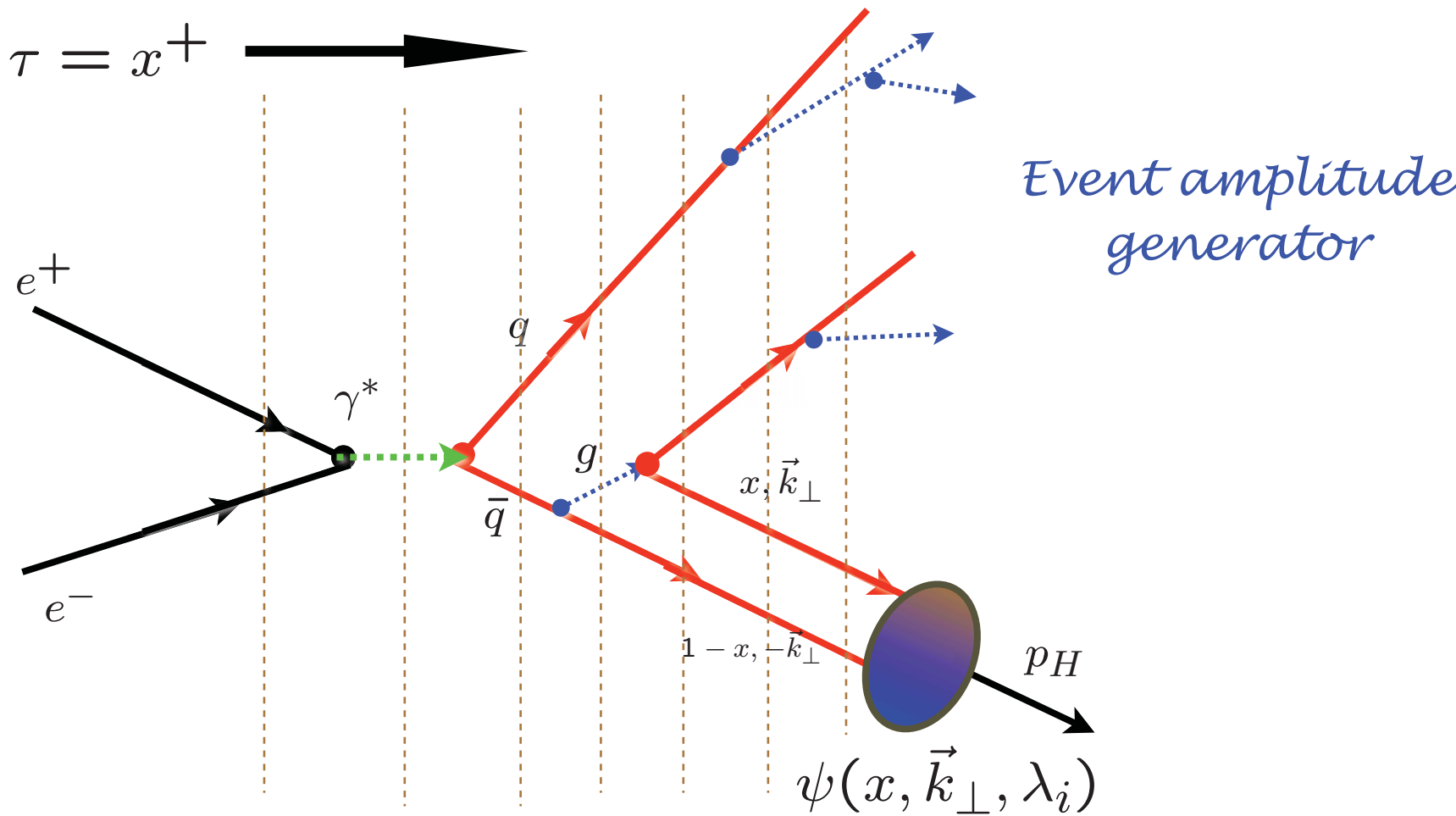
“Hadronization” at the Amplitude Level

Hadronization at the Amplitude Level



Construct helicity amplitude using Light-Front Perturbation theory; coalesce quarks via LFWFs

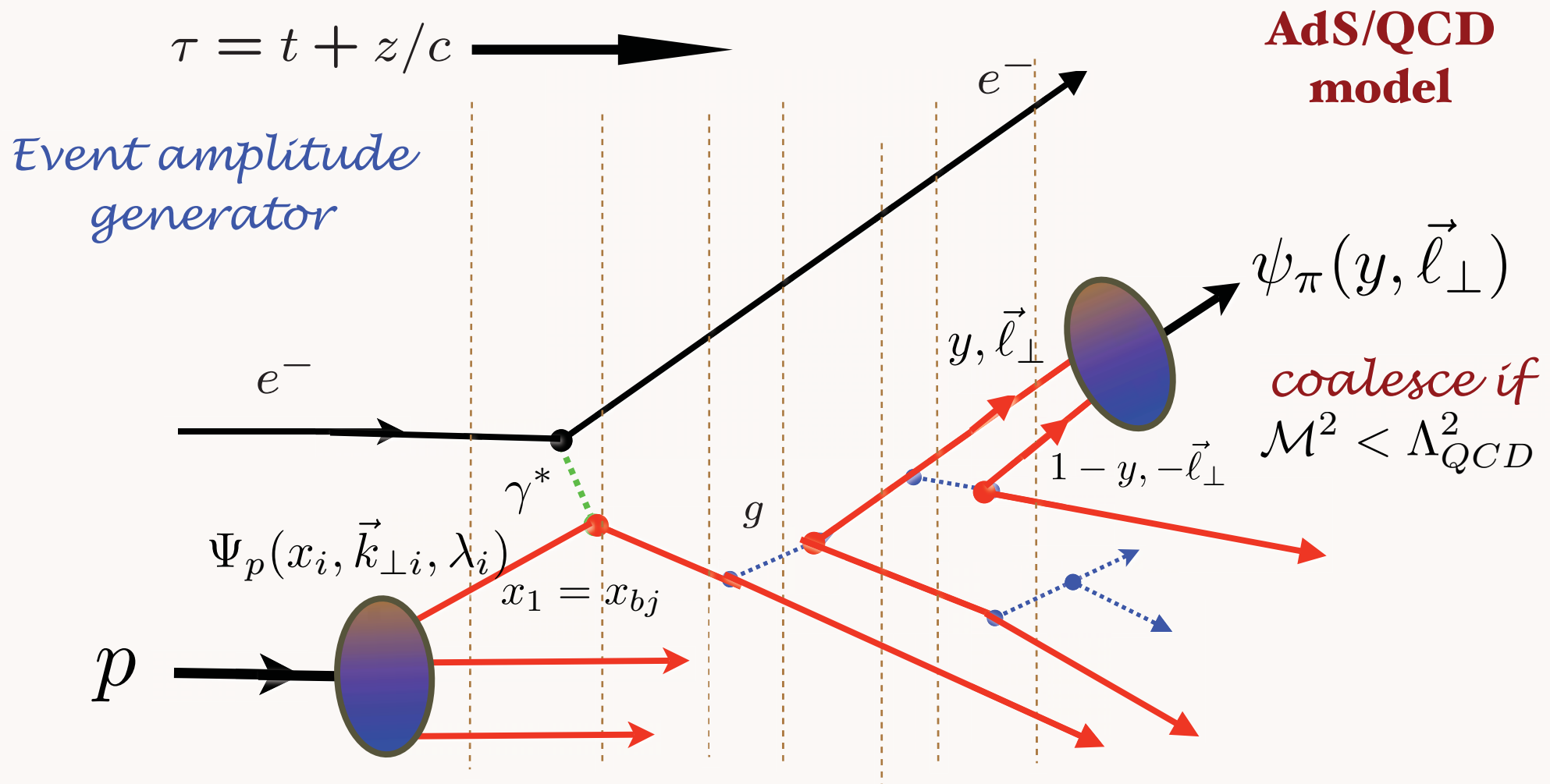
Hadronization at the Amplitude Level



AdS/QCD
Hard Wall
Confinement:

Capture if $\zeta^2 = x(1-x)b_\perp^2 > \frac{1}{\Lambda_{QCD}^2}$
 i.e.,
 $\mathcal{M}^2 = \frac{k_\perp^2}{x(1-x)} < \Lambda_{QCD}^2$

Jet Hadronization at the Amplitude Level



Construct helicity amplitude using Light-Front Perturbation theory; coalesce quarks via Light-Front Wavefunctions

Three Pictures of High Energy Lepton-Proton Collisions

Infinite momentum frame

Parton Model

Simple Virtual Photon Probes Complex Evolved Proton

Proton Rest Frame

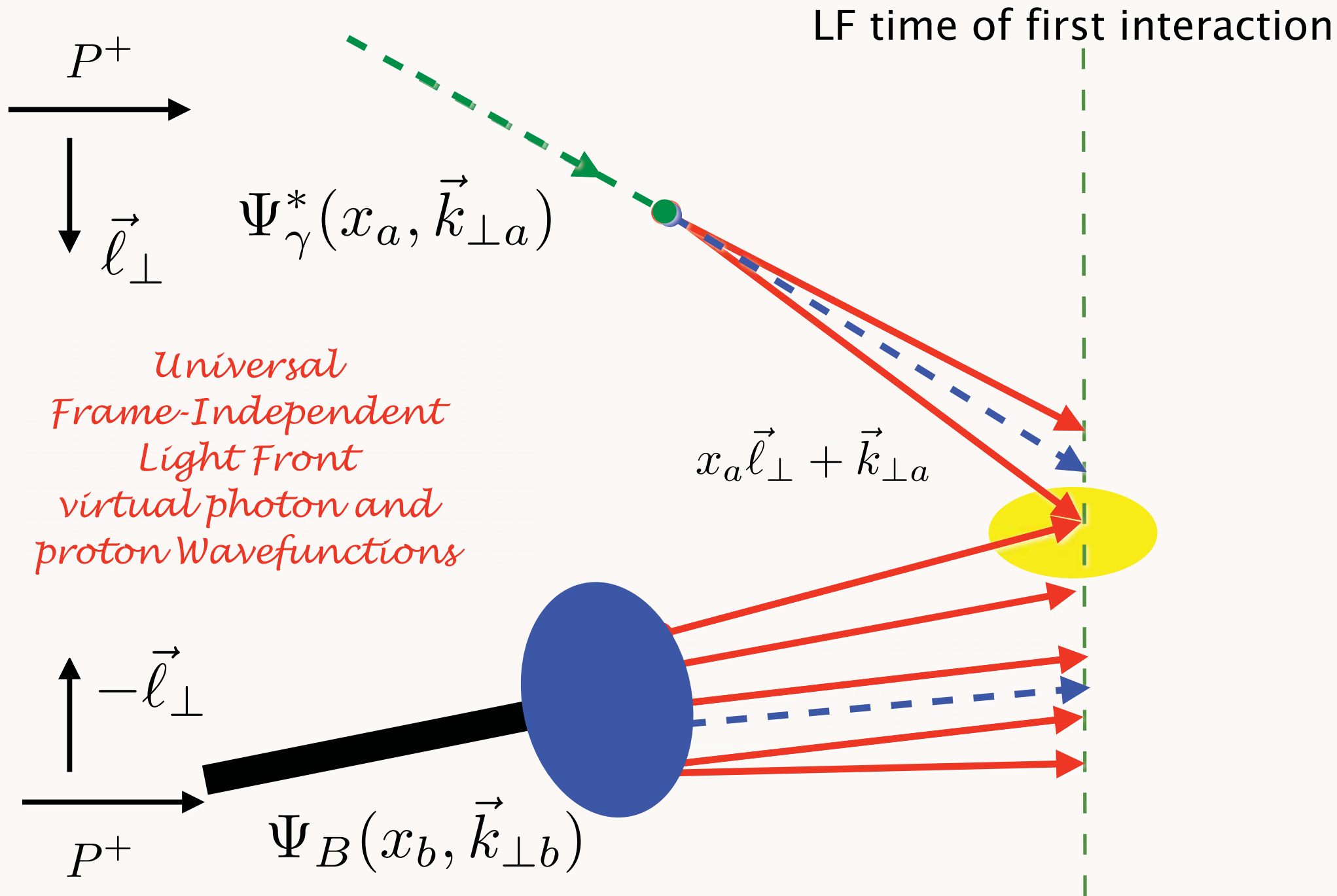
Color-Dipole Model

Color Dipole of Virtual Photon Scatters on a Static Proton

Frame-Independent

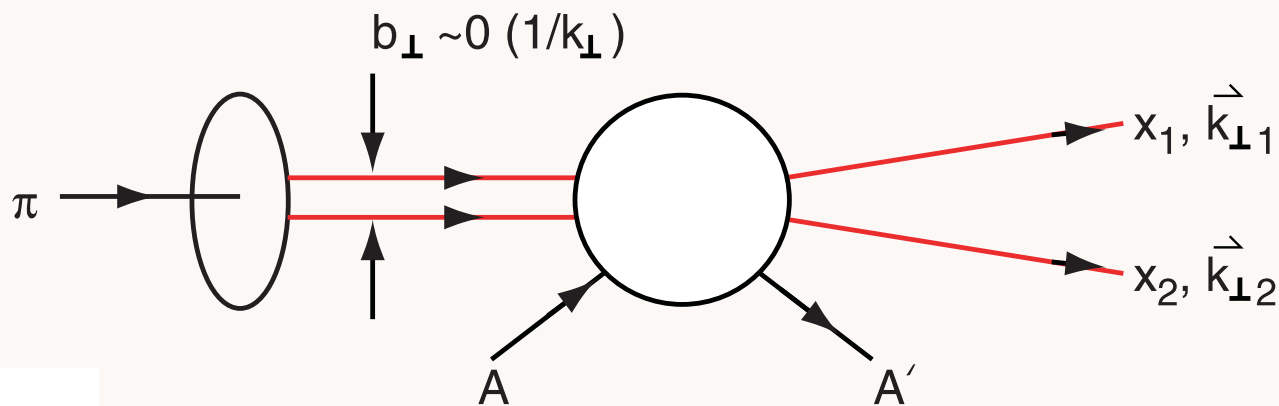
**Light-Front
Hamiltonian Theory**

Collision of Light-Front Wavefunctions
of Virtual Photon and Proton



Diffraction Dissociation of Pion into Quark Jets

E791 Ashery et al.

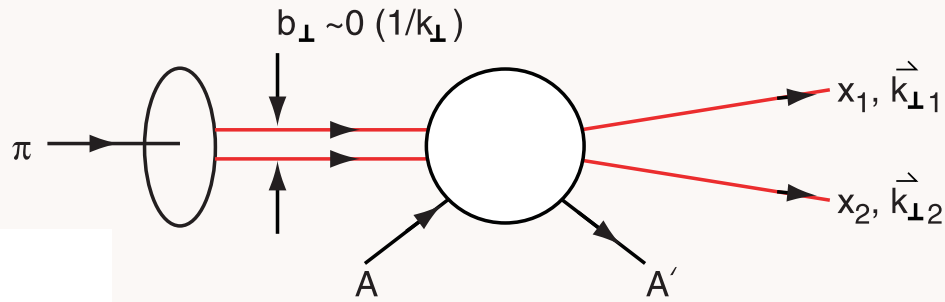


$$M \propto \frac{\partial^2}{\partial^2 k_{\perp}} \psi_{\pi}(x, k_{\perp})$$

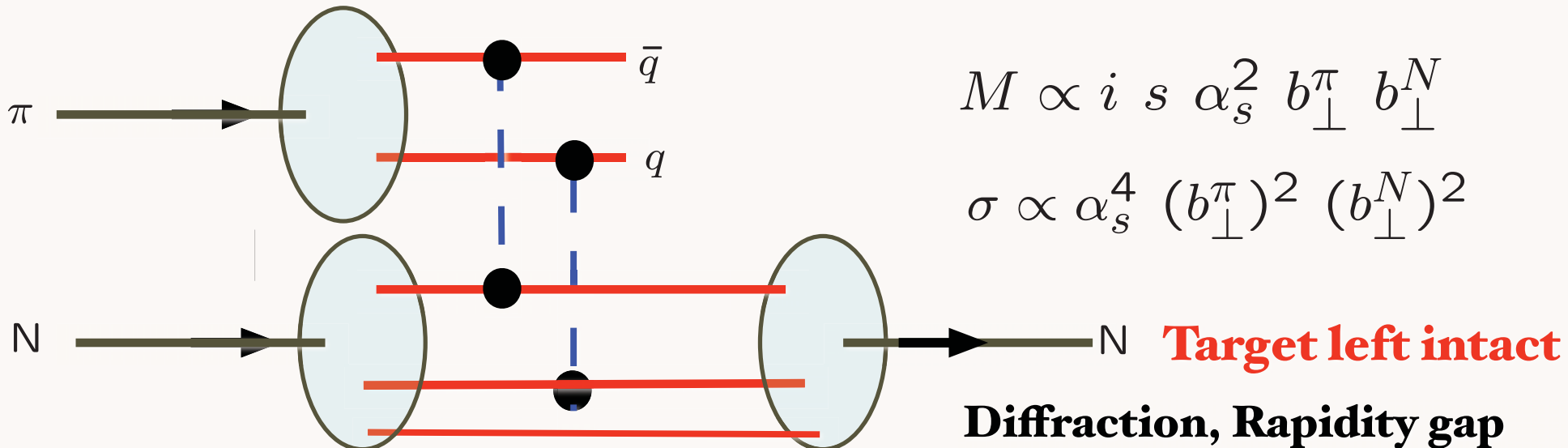
Measure Light-Front Wavefunction of Pion

Minimal momentum transfer to nucleus
Nucleus left Intact!

Key Ingredients in Ashery Experiment



Two-gluon exchange gives imaginary amplitude proportional to energy, constant diffractive cross sections



$$M \propto i s \alpha_s^2 b_{\perp}^{\pi} b_{\perp}^N$$

$$\sigma \propto \alpha_s^4 (b_{\perp}^{\pi})^2 (b_{\perp}^N)^2$$